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in the case of systems**

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## **Bohm Aharonov effects for bounded states in the case of systems**

by

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**ABSTRACT.** — We study the comparison problem for the eigenvalues of the covariant Laplacian with electric potential acting on the sections of vector bundle with structure group  $\mathcal{U}(m)$ .

**RÉSUMÉ.** — On s'intéresse dans cet article à un problème de comparaison de valeurs propres pour le Laplacien covariant, avec potentiel électrique, agissant sur les sections d'un fibré vectoriel de groupe structural  $\mathcal{U}(m)$  ( $m \in \mathbb{N}^*$ ).

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### **INTRODUCTION**

Let  $(M, g)$  be an  $n$ -dimensional connected orientable Riemannian manifold with (possibly empty) boundary  $\partial M$ ,  $(E, (\cdot, \cdot))$  be a Hermitian ( $C^\infty$ ) bundle over  $M$  with rank  $m$ . We denote by  $A^0(M, E) = C^\infty(M, E)$  the set of  $C^\infty$  sections of  $E$ . More generally we denote by  $A^p(M)$  the set of the  $C^\infty - p$  forms on  $M$  and by  $A^p(M, E)$  the set of  $E$ -valued  $C^\infty - p$  forms on  $M$ .

As usual, we put

$$A_0^p(M, E) = \{ \Theta \in A^p(M, E) : \text{supp } \Theta \subset \text{int}(M) = M \setminus \partial M \}$$

and we introduce on  $A_0^0(M, E)$ ,  $A_0^1(M, E)$  the inner products  $[ \cdot, \cdot ]_0$ ,  $[ \cdot, \cdot ]_1$  defined by:

$$[\xi, \xi']_0 = \int_M (\xi, \xi')(x) dv, \quad \text{for } \xi, \xi' \in A_0^0(M, E),$$

$$[\Theta, \Theta']_1 = \int_M \langle \Theta, \Theta' \rangle(x) dv, \quad \text{for } \Theta, \Theta' \in A_0^1(M, E).$$

where  $\langle \cdot, \cdot \rangle$ ,  $dv$  denote the natural metric in  $T^*M \otimes E$  induced by  $g$  and the Riemannian volume element, respectively. Let  $\nabla : C^\infty(M, E) \rightarrow A^1(M, E)$  be a connection on  $E$ , compatible with the Hermitian structure (cf. [13]). The dual operator

$$\nabla^* : A_0^1(M, E) \rightarrow C_0^\infty(M, E)$$

of  $\nabla|_{C_0^\infty(M, E)}$  is defined by:

$$\forall \Theta \in A_0^1(M, E), \quad [\nabla^* \Theta, \xi]_0 = [\Theta, \nabla \xi]_1 \quad \forall \xi \in C_0^\infty(M, E). \quad (0.1)$$

We consider a positive  $C^\infty$  function  $V$  on  $M$  and we introduce the two following positive formally self-adjoint elliptic operators  $H_{\nabla, V}$ ,  $H_V$  defined by:

$$\text{Dom}(H_{\nabla, V}) = C_0^\infty(M, E), \quad H_{\nabla, V} = \nabla^* \cdot \nabla + V,$$

$$\text{Dom}(H_V) = C_0^\infty(M), \quad H_V = d^* \cdot d + V.$$

In the case  $\partial M \neq \emptyset$ , the Bochner-Laplace (resp. Laplace) operator  $H_{\nabla, V}^{M, E}$  (resp.  $H_V^M$ ) is the Dirichlet realization for  $M$  in the completion  $L^2(M, E)$  [resp.  $L^2(M)$ ] of the pre-Hilbert space

$$(A_0^0(M, E), [ \cdot, \cdot ]_0) \text{ (resp. } C_0^\infty(M)$$

with the usual scalar product). If  $\partial M = \emptyset$ , we denote by  $H_{\nabla, V}^{M, E}$ ,  $H_V^M$  the unique self-adjoint extension (the closure) [7] of operators  $H_{\nabla, V}$ ,  $H_V$  in the space  $L^2(M, E)$  and  $L^2(M)$ , respectively. The problem we want to address in this work is, assuming to simplify  $H_{\nabla, V}^{M, E}$  and  $H_V^M$  with compact resolvent, is the following:

Under which conditions on  $E$  and  $\nabla$  do the operators  $H_{\nabla, V}^{M, E}$ ,  $H_V^M$  admit the same first eigenvalue or more generally the same spectrum.

We shall consider two cases:

**Case I.** —  $E = M \times \mathbb{C}^m$  and  $M$  satisfies one of the following properties:

- (P1)  $M$  is compact
- (P2)  $M$  is the closure of an open set (possibly unbounded)  $Q$  of  $\mathbb{R}^n$  with regular bounded boundary  $\partial Q$ ,
- (P3)  $M = \mathbb{R}^n$ .

We assume, in the case when  $M$  is not compact, that the electric potential  $V$  verifies:

$$V(x) \rightarrow +\infty \quad \text{as } |x| \rightarrow +\infty. \quad (0.2)$$

**Case II.** –  $E$  is not necessarily trivial but  $M$  is compact.

It is well known ([10], [11], ...) that if  $M$  is compact, the spectra of  $H_{\nabla, V}^{M, E}$  and  $H_V^M$  are increasing sequences of positive eigenvalues tending to  $+\infty$ . When  $M$  is not compact, this follows from the condition (0.2) (See [11] for  $H_V^M$ ; and Theorem 2.3 of [6], Theorem 1.2 of [1] for the operator  $H_{\nabla, V}^{M, E}$  with  $E = M \times \mathbb{C}^m$ ). As we shall see, the comparison problem for the spectra of two such operators is naturally related to the gauge transformations. In section 2 of this work, we discuss briefly this idea and we give a characterization for the trivial connections. We present in section 2 comparison theorems for the **case I** generalizing results obtained by Helffer [5], Shigekawa [12] in the scalar case and Manabe-Shigekawa [10] in the case of systems. We study the **case II** in section 4 and we give a theorem extending results of Kuwabara [8].

I would like to thank my adviser Bernard Helffer who suggested me this study.

## 1. GAUGE TRANSFORMATIONS AND TRIVIAL CONNECTIONS

Let  $e_B = (e_B^1, \dots, e_B^m)$  be a local orthonormal frame over an open set  $B$  of  $M$ , *i.e.*,  $e^i \in C^\infty(B, E|_B)$  for  $1 \leq i \leq m$  such that  $(e_B^i(x))_{1 \leq i \leq m}$  is an orthonormal basis of a fibre  $E_x$  for each  $x \in B$ . Then,

$$\nabla e_B^i = \sum_s \omega_{si} \otimes e^s, \quad \text{where } \omega_{si} \in A^1(B) \quad \text{for } 1 \leq i, s \leq m. \quad (1.1)$$

We call the matrix 1-form  $\omega = [\omega_{is}]_{i,s}$  the connection form of  $\nabla$  with respect to the frame  $e_B$ . Because  $\nabla$  is compatible with the metric  $(\cdot, \cdot)_{E|_B}$ ,  $\omega$  takes values in the Lie algebra  $\mathcal{M}_{m,a}$  of the unitary group  $\mathcal{U}(m)$ . Let  $\xi \in A^0(M, E)$  and  $\xi_{|B} = (\xi_1, \dots, \xi_m)$  be the (local) trivialization of  $\xi$  with respect to  $e_B$  (defined by  $\xi_{|B} = \sum_i \xi_i e_B^i$ ).

If  $f_B = (f_B^1, \dots, f_B^m)$  is another orthonormal frame over  $B$  and if  $T = [t_{ij}]$  is the  $\mathcal{U}(m)$ -valued function on  $B$  such that:  $f_B^i = \sum_s t_{si} e_B^s$ , or in matrix-notations  $f_B = e_B \cdot T$ , then, the connection form  $\omega'$  of  $\nabla$  and the trivialization  $\xi'_B$  of  $\xi$  with respect to  $f_B$  are given by:

$$\omega' = T^* \omega T + T^* dT, \quad (1.2)$$

$$\xi'_B = T^* \xi_B. \quad (1.3)$$

Transformations of the form (1.2) and (1.3) are called (local) gauge transformations. If  $E$  is trivialisable and if  $e_M, f_M$  are (global) frames of  $E$  over  $M$ , then, for  $\xi \in A_0^0(M, E)$ , we have (with the notations of (2.2) and

(2.3)):

$$H_{\omega, \nu}(\xi_M) = (T \cdot H_{\omega', \nu} \cdot T^*)(\xi'_M), \tag{1.4}$$

where  $T \in C^\infty(M, \mathcal{U}(m))$  and  $H_{\omega, \nu} = (d + \omega)^* \cdot (d + \omega) + V \otimes 1$  is the representation of  $H_{\nabla, \nu}$  with respect to the frame  $e_M$ . Consequently, in the case I,  $H_{\nabla, \nu}^E$  is nothing but a Schrödinger operator  $H_{\omega, \nu}^M$  with magnetic potential  $\omega \in A^1(M, \mathcal{M}_{m, a})$ .

Properties (2.4) and (2.2) say that the operators  $H_{\omega, \nu}^M$  and  $H_{\nu}^M \otimes 1$  are unitary equivalent if there exists  $S \in C^\infty(M, \mathcal{U}(m))$  such that  $dS = \omega \cdot S$  on  $M$ . A such form  $\omega$  is called trivial.

Our problem is now to find characterizations of such forms. Let  $\omega \in A^1(M, \mathcal{M}_{m, a})$ . We call  $\omega$  flat if its curvature  $K(\omega) = d\omega + \omega \wedge \omega$  vanishes. It is easy to see that a trivial 1-form  $\omega$  is flat. Let  $\gamma: [0, 1] \rightarrow M$  be a closed curve in  $M$ ,  $\gamma^*(\omega) = A_{\gamma, \omega}(t) dt$  be the pull-back of  $\omega$  by  $\gamma$ , and consider the associated system of differential equations:

$$\psi' = \psi \cdot A_{\gamma, \omega}, \quad \Psi(0) = I_m. \tag{1.5}$$

It is well known (See for example [2]) that a system (1.5) has a unique solution  $g$  in  $C^1([0, 1], \mathcal{U}(m))$ . Let us define the holonomy class of  $\omega$  with respect to  $\gamma$  by:

$$U_\gamma(\omega) = \{ U \in \mathcal{U}(m) \text{ such that: } U \text{ and } g(1) \text{ are unitary equivalent} \}.$$

For example, we have for a closed 1-form  $\omega$  in  $A^1(M, \mathcal{M}_{1, a})$ :

$$U_\gamma(\omega) = \left\{ \exp \left( \int_\gamma \omega \right) \right\}.$$

One can verify (See [4]) that, if  $\omega$  is flat, then  $U_\gamma(\omega)$  depends only on the homotopy class of  $\gamma$  and that for  $T \in C^1(M, \mathcal{U}(m))$ ,  $\omega_T = T^* \cdot \omega \cdot T + T^* \cdot dT$ ; we have:

$$K(\omega_T) = T^* \cdot K(\omega) \cdot T = 0, \tag{1.6}$$

$$U_\gamma(\omega_T) = U_\gamma(\omega). \tag{1.7}$$

The following theorem is probably classical (See [4])

**THEOREM 1.1.** — *For  $\omega \in A^1(M, \mathcal{M}_{m, a})$ . The following conditions (i) and (ii) are equivalent:*

(a)  $\omega$  is trivial,

(ii) (a):  $\omega$  is flat, (b):  $U_\gamma(\omega) = \{ I_m \}$ , for each closed curve  $\gamma$  in  $M$ .

**COROLLARY 1.2.** — *If  $M$  is simply connected. Then,  $\omega$  is trivial if and only if it is flat.*

Let us look at the more general case of connections and consider a system  $(B_\alpha, e_\alpha)_{\alpha \in I}$  of local trivializations of  $E$ , i. e.,  $(B_\alpha)_\alpha$  is an open connected cover of  $M$  and  $e_\alpha$  is an orthonormal frame over  $B_\alpha$  for each  $\alpha \in I$ . For  $B_{\alpha\beta} = B_\alpha \cap B_\beta \neq \emptyset$ , the  $\mathcal{U}(m)$ -valued functions  $g_{\alpha\beta}$  on  $B_{\alpha\beta}$  such that

$e_\beta = e_\alpha g_{\alpha\beta}$  are called transition functions. If  $\omega_\alpha$  is the connection form of  $\nabla$  with respect to  $e_\alpha$ ,  $K(\omega_\alpha)$  is called the curvature form of  $\nabla$  with respect to  $e_\alpha$ . By (2.2), we have:

$$\omega_\beta = g_{\alpha\beta}^* \cdot \omega_\alpha \cdot g_{\alpha\beta} + g_{\alpha\beta}^* \cdot dg_{\alpha\beta}, \quad (1.8)$$

$$K(\omega_\beta) = g_{\alpha\beta}^* \cdot K(\omega_\alpha) \cdot g_{\alpha\beta} \quad \text{on } B_{\alpha\beta}. \quad (1.9)$$

The property (1.9) says that the condition,  $K(\omega_\alpha) = 0$  for each  $\alpha \in I$ , depends only on the connection  $\nabla$ . Connections which satisfy this condition are called flats. We say that  $\nabla$  is trivial if there exist a system of local trivializations  $(B_\alpha, f_\alpha)_{\alpha \in I}$  of  $E$  such that the corresponding transition functions (resp. connection forms)  $g'_{\alpha\beta}$  (resp.  $\omega'_\alpha$ ) are all identity functions (resp. zero forms). As a necessary condition,  $E$  is trivialisable and  $\nabla$  is flat. We start from these conditions and we consider the connection form  $\omega$  of  $\nabla$  with respect to a given global frame  $e_M$  of  $E$ . It is clear, by (1.6) and (1.9), that for a closed curve  $\gamma$  in  $M$ , the class  $U_\gamma(\omega)$  is independent of a choice of  $e_M$ . We define the holonomy class of  $\nabla$  with respect to  $\gamma$  by:  $U_\gamma(\nabla) = U_\gamma(\omega)$ . We can then state Theorem 1.1 as follows:

**THEOREM 1.1.** — *Suppose that  $E$  is trivialisable. Then, the following conditions are equivalent:*

- (i)  $\nabla$  is trivial,
- (ii) (a):  $\nabla$  is flat, (b):  $U_\gamma(\nabla) = \{I_m\}$ , for each closed curve  $\gamma$  in  $M$ .

**REMARK 1.3.** — Let  $\nabla$  be a flat connection on  $E$  (unnecessarily trivialisable). Using the fact that a flat connection is locally trivial, we construct in [4] a holonomy class  $U_\gamma(\nabla)$ , which coincides in the case of a trivialisable vector bundle  $E$  with the class defined above, and such that, if  $U_\gamma(\nabla) = \{I_m\}$ , then  $E$  is trivialisable and  $\nabla$  is trivial.

## 2. COMPARISON THEOREMS, CASE I

Through this section, we assume that  $E = M \times \mathbb{C}^m$  and that  $M$  satisfies one of the properties (P1), (P2), (P3) mentioned in Section 1. If  $A^0(M, E)$  is identified (in a natural way) with  $C^\infty(M, \mathbb{C}^m)$ , then  $H_{\nabla, V}$  can be regarded [by (1.4)] as a Schrödinger operator  $H_{\omega, V} = \nabla_\omega^* \cdot \nabla_\omega + V$ , where  $\nabla_\omega = d + \omega$ , with a (fixed) magnetic potential  $\omega$  in  $A^1(M, \mathcal{M}_{m, a})$  and electric potential  $V$ . Recall that if  $\partial M \neq \emptyset$ ,  $H_{\omega, V}^M$  is the Friedrichs' extension [11] associated to the positive sesquilinear form  $q_{\omega, V}$  defined on  $C_0^\infty(M, \mathbb{C}^m)$  by:

$$q_{\omega, V}(\varphi, \psi) = \int_M (\langle \nabla_\omega \varphi, \nabla_\omega \psi \rangle + (V\varphi, \psi))(x) dv,$$

for  $\varphi$  and  $\psi$  in  $C_0^\infty(M, \mathbb{C}^m)$ .

Let  $\lambda_\omega^M$  (resp.  $\lambda_0^M$ ) be the first eigenvalue of  $H_{\omega, \nu}^M$  (resp.  $H_\nu^M$ ). As we know by the Kato's inequality (given in [6] for the case of systems), we have:

$$\lambda_0^M \leq \lambda_\omega^M. \tag{2.1}$$

Let  $u_0$  be the first eigenfunction of  $H_\nu^M$  attached to  $\lambda_0^M$ . We know that  $u_0$  can be chosen such that  $u_0 > 0$  on  $\text{int}(M)$  and  $\|u_0\|_0 = 1$ . Using elementary computations and the fact that  $\omega$  is skew Hermitian, we get the following lemma (due essentially to Lavine-O'Carroll [9]):

LEMMA 2.1. — For  $\varphi \in C_0^\infty(M, \mathbb{C}^m)$ ,

$$\|\nabla_\omega \varphi - du_0 \cdot \varphi / u_0\|_1^2 = q_{\omega, \nu}(\varphi, \varphi) - \lambda_0^M \|\varphi\|_0^2.$$

The first consequence is of course that we get, as in [5], another proof of (2.1). Suppose now that  $\lambda_\omega^M = \lambda_0^M$  and consider a normalized eigenfunction  $u_\omega$  of  $H_{\omega, \nu}^M$  attached to  $\lambda_\omega^M$ . We deduce from Lemma 2.1 and using a minimizing sequence tending to  $u_\omega$  in  $L^2(M, \mathbb{C}^m)$  that:

$$[\nabla_\omega u_\omega - du_0 \cdot u_\omega / u_0, \alpha]_1 = 0, \text{ for each } \alpha \in A_0^1(M, \mathbb{C}^m). \tag{2.2}$$

Consequently,

$$u_0 \cdot \nabla_\omega (u_\omega / u_0) = \nabla_\omega u_\omega - du_0 \cdot u_\omega / u_0 = 0,$$

on  $\text{int}(M)$  (since  $u_\omega$  and  $u_0$  are  $C^\infty$  on  $M$ ).

That is to say,

$$\nabla_\omega (u_\omega / u_0) = 0, \text{ on } \text{int}(M). \tag{2.3}$$

Now, let  $\lambda_{\omega, 1}^M, \lambda_{\omega, 2}^M, \dots, \lambda_{\omega, k}^M$  ( $k \leq m$ ) be the  $k$ -first eigenvalues of  $H_{\omega, \nu}^M$ . Then, we have

PROPOSITION 2.2. — If  $\lambda_{\omega, 1}^M = \lambda_{\omega, 2}^M = \dots = \lambda_{\omega, k}^M = \lambda_0^M$ . Then, there exists  $\varphi_1, \varphi_2, \dots, \varphi_k$  in  $C^\infty(M, \mathbb{C}^m)$  such that, for each  $x \in M$ ,  $(\varphi_q(x))_q$  form an orthonormal system of  $\mathbb{C}^m$  with:

$$\nabla_\omega \varphi_q = 0, \text{ for each } 1 \leq q \leq k. \tag{2.4}$$

*Proof.* — Let  $(u_{\omega, q})_{1 \leq q \leq k}$  be a system of  $k$  normalized eigenfunctions of  $H_{\omega, \nu}^M$  attached to  $\lambda_\omega^M$ , and define  $\varphi_q = u_{\omega, q} / u_0$  on  $\text{int}(M)$ . It is clear that  $(\varphi_q)_{1 \leq q \leq k}$  satisfies (2.4) on  $\text{int}(M)$ . On the other hand, using maximum principle (Lemma 3.4 in [3] applied to  $\Delta - V$  and  $-u_0$ ), we get that:

$$\partial u_0 / \partial N = \nabla u_0 \cdot N < 0 \text{ on } \partial M,$$

where  $N: \partial M \rightarrow \mathbb{R}^n$  is the outward normal vector field to  $\partial M$  (note that  $\partial M$  is a regular bounded set). Then, let us define  $\varphi_q(x_0)$ , for  $x_0 \in \partial M$  and  $1 \leq q \leq k$ , by:

$$\begin{aligned} \varphi_q(x_0) &= \lim_{t \rightarrow 0, t > 0} \{ u_{\omega, q}(x_0 - t N(x_0)) / u_0(x_0 - t N(x_0)) \} \\ &= (\partial u_{\omega, q} / \partial N)(x_0) / (\partial u_0 / \partial N)(x_0). \end{aligned}$$

In order to show that  $\varphi_q$  verifies (2.4) on  $\partial M$ , it is sufficient to consider the case  $M = \bar{Q}$ . Let  $\mathcal{V}$  be a neighbourhood of  $\partial Q$  and  $\Phi$  in  $C^\infty(\mathcal{V})$  such that:

$$\partial Q = \{x \in \mathcal{V} : \Phi(x) = 0\} \quad \text{and} \quad (\nabla \Phi)(x) \neq 0 \quad \text{for } x \in \mathcal{V}.$$

Then, the field  $\bar{N}$  defined on  $\mathcal{V}$  by:  $N(x) = (\nabla \Phi) / |(\nabla \Phi)|$ , is  $C^\infty$  on  $\mathcal{V}$  and extend  $N$  on  $\bar{Q}$ . Let  $\vec{A} = (A_1, A_2, \dots, A_n) \in C^\infty(\bar{Q}, \mathcal{M}_{m,a})$  such that:  $\omega = \sum_j A_j dx_j$  on  $\bar{Q}$ ,  $1 \leq q \leq k$ , and  $x_0 \in \partial Q$ . By a simple computation, we see that, on a suitable neighbourhood of  $x_0$ , we have:

$$\varphi_q = (\nabla u_{\omega, q} \cdot \nabla \Phi) / (\nabla u_0 \cdot \nabla \Phi) + ((\vec{A} \cdot \nabla \Phi) / (\nabla u_0 \cdot \nabla \Phi)) u_{\omega, q}.$$

In particular,

$$\varphi_q \in C^\infty(\bar{Q}, \mathbb{C}^m)$$

$$\text{and} \quad (\nabla \varphi_q + \vec{A} \varphi_q)(x_0) = 0, \quad \text{for } 1 \leq q \leq k \quad \text{and} \quad x_0 \in \partial Q.$$

Now, we show the second part of this proposition. Let us remark that as a consequence of the Cauchy uniqueness theorem for linear systems of differential equations, we have:

LEMMA 2.3. — *If  $x_0 \in M$ ,  $\alpha \in A^1(M, \mathcal{M}_{m,a})$  and  $\psi \in C^1(M, \mathbb{C}^m)$  such that:  $\nabla_\alpha \psi = 0$ ,  $\psi(x_0) = 0$ . Then,  $\psi = 0$  on  $M$ .*

By this lemma, we obtain easily that, for  $x \in M$ , the  $(\varphi_q(x))_q$  are linearly independent in  $\mathbb{C}^m$ . Let us verify that, for  $x \in M$  and  $1 \leq p, q \leq k$ ,  $(\varphi_p(x), \varphi_q(x)) = \delta_q^p$  (where  $\delta_q^p$  is the Kronecker delta).

By differentiation of the application  $S_q^p = (\varphi_p, \varphi_q)$  [which is in  $C^\infty(M, \mathbb{C}^m)$ ] and using the fact that  $\omega$  is skew Hermitian, we obtain:

$$\begin{aligned} dS_q^p &= \langle d\varphi_p, \varphi_q \rangle_0 + \langle \varphi_p, d\varphi_q \rangle_0 \\ &= \langle -\omega \cdot \varphi_p, \varphi_q \rangle_0 + \langle \varphi_p, -\omega \cdot \varphi_q \rangle_0 = 0. \end{aligned}$$

Here it is understood that the inner products on the right are defined by the requirement that:  $\langle \Theta, \varphi \rangle_0 = \sum_s \bar{\varphi}_s \theta_s \in A^1(M)$ , for

$$\Theta = (\theta_s)_s \in A^1(M, \mathbb{C}^m) \quad \text{and} \quad \varphi = (\varphi_s)_s \in A^0(M, \mathbb{C}^m).$$

Then,  $S_q^p$  is equal to a constant  $c_q^p$  on  $M$  (note here that  $M$  is connected) and finally

$$\delta_q^p = \int_M (u_{\omega, p}, u_{\omega, q})(x) dv = \int_M |u_0|^2(x) (\varphi_p, \varphi_q)(x) dv = c_q^p.$$

Let us translate this result on the curvature of  $\omega$ . By differentiation of (2.4), we obtain:

$$K(\omega) \varphi_q = 0, \quad \text{for } 1 \leq q \leq k. \quad (2.5)$$



Let us define the kernel of  $K(\omega)$  as the subset of the trivial bundle  $M \times \mathbb{C}^m$ :

$$\ker K(\omega) = \{ (x, v) \in M \times \mathbb{C}^m : K(\omega)(x) [\partial_j(x), \partial_l(x)], \\ v = 0, \quad \text{for } 1 \leq j, l \leq n \},$$

where  $\{\partial_j(x)\}_j$  is the natural basis of  $T_x M$ . Note here that  $\ker K(\omega)$  defined in this way is independent of a choice of a basis in  $T_x M$ . Moreover, it is invariant under global gauge transformations. Suppose that  $\lambda_{\omega, 1}^M = \lambda_{\omega, 2}^M = \dots = \lambda_{\omega, k}^M = \lambda_0^M$ , and consider  $k$ -functions  $(\varphi_q)_q$  satisfying the above proposition. Let  $\mathcal{K}$  be the trivial subbundle of  $M \times \mathbb{C}^m$  generated by  $(\varphi_q)_q$ , and  $\mathcal{K}^\perp$  the orthogonal fiber subbundle to  $\mathcal{K}$ . Condition (2.5) says that  $\ker K(\omega)$  contains  $\mathcal{K}$ . More precisely, we have:

LEMMA 2.4. — Assume that  $\lambda_{\omega, 1}^M = \lambda_{\omega, 2}^M = \dots = \lambda_{\omega, k}^M = \lambda_0^M$ . Then, the following equivalent conditions are satisfied:

- (i):  $\nabla_\omega$  restricted to  $A^0(M, \mathcal{K})$  takes values in  $A^1(M, \mathcal{K})$ ,
- (ii):  $\nabla_\omega$  restricted to  $A^0(M, \mathcal{K}^\perp)$  takes values in  $A^1(M, \mathcal{K}^\perp)$ .

In other words, the restriction of  $\nabla_\omega$  to  $A^0(M, \mathcal{K})$  define a connection  $\nabla_{\omega, \mathcal{K}}$  on  $\mathcal{K}$ .

*Proof.* — The equivalence between (i) and (ii) results from the following relation:

$$\langle \nabla_\omega f, \psi \rangle_0 = - \langle f, \nabla_\omega \psi \rangle_0, \quad \text{for } f \in A^0(M, \mathcal{K}) \text{ and } \psi \in A^0(M, \mathcal{K}^\perp).$$

Let us prove (i). Consider  $f = \Sigma_q(f, \varphi_q) \cdot \varphi_q \in A^0(M, \mathcal{K})$  and using (2.4), we obtain:

$$\begin{aligned} \nabla_\omega f &= \Sigma_q [d(f, \varphi_q) \cdot \varphi_q + (f, \varphi_q) \cdot d\varphi_q + (f, \varphi_q) \cdot \omega \cdot \varphi_q] \\ &= \Sigma_q d(f, \varphi_q) \cdot \varphi_q \in A^1(M, \mathcal{K}). \end{aligned}$$

Let us give the main theorem of this section.

THEOREM 2.5. — The following three conditions are equivalent:

- (i)  $\lambda_{\omega, 1}^M = \lambda_{\omega, 2}^M = \dots = \lambda_{\omega, k}^M = \lambda_0^M$
- (ii)  $\ker K(\omega)$  contains a trivial subbundle  $\mathcal{K}$  of  $M \times \mathbb{C}^m$  of rang  $k$ , such that:

- (a):  $\nabla_{\omega|_A} A^0(M, \mathcal{K}) : A^0(M, \mathcal{K}) \rightarrow A^1(M, \mathcal{K})$ ,
- (b):  $\nabla_{\omega, \mathcal{K}}$  is flat,
- (c):  $U_\gamma(\nabla_{\omega, \mathcal{K}}) = \{I_k\}$ , for each closed curve  $\gamma$  in  $M$ .
- (iii)  $k \cdot \text{Sp}(H_{\nabla}^M) \subset \text{Sp}(H_{\omega, \nabla}^M)$ , where

$$k \cdot \text{Sp}(H_{\nabla}^M) = \text{Sp}(H_{\nabla}^M) \cup \text{Sp}(H_{\nabla}^M) \cup \dots \cup \text{Sp}(H_{\nabla}^M), (k \text{ times}).$$

*Proof.* — The assertion (i)  $\Rightarrow$  (ii) is an easy consequence of Proposition 2.3, Lemma 2.4 and Theorem 1.1'. Let us prove (ii)  $\Rightarrow$  (iii), which is the non trivial part of the statements. Consider a frame  $\mathcal{E} = (e_q)_q$ ,

$e_q \in C^\infty(M, \mathbb{C}^m)$  for  $1 \leq q \leq k$ , of  $\mathcal{X}$  over  $M$ . Using (a), we can write:

$$\nabla_\omega e_q = \sum_s \langle (d + \omega) e_q, e_s \rangle_0 \cdot e_s.$$

This means that the 1-form  $\omega_{\mathcal{X}} = \langle \langle (d + \omega) e_i, e_l \rangle_0 \rangle_{1 \leq i, l \leq k}$  is the connection form of  $\nabla_{\omega, \mathcal{X}}$  with respect to  $\mathcal{E}$ . Now, conditions (b) and (c) say that  $\nabla_{\omega, \mathcal{X}}$  is trivial:

$$\exists W \in C^\infty(M, \mathcal{U}(k)) : dW = W \cdot \omega_{\mathcal{X}}. \quad (2.6)$$

Using elementary computations, we see that if  $(\eta_s)_s$  [resp.  $(\delta_l)_l$ ] is the canonical basis of  $\mathbb{C}^k$  (resp.  $\mathbb{C}^m$ ) and  $E = \langle [e_i, \delta_l] \rangle_{1 \leq i \leq k, 1 \leq l \leq m}$ , then

$$(dE + \omega \cdot E - E \cdot W^* \cdot dW) \eta_s \in A^1(M, \mathcal{X}) \cap A^1(M, \mathcal{X}^\perp), \quad \text{for } 1 \leq s \leq k.$$

Consequently,

$$dE + \omega \cdot E - E \cdot W^* \cdot dW = 0 \quad \text{in } A^1(M, \mathcal{M}_{m \times k}), \quad (2.7)$$

where  $\mathcal{M}_{m \times k}$  is the set of  $m \times k$ -matrix.

Let  $\lambda \in \text{Sp}(H_V^M)$ ,  $u$  an associated eigenfunction of  $H_V^M$ , and set:

$$u_q = u \cdot E \cdot W^* \cdot \eta_q \in C^\infty(M, \mathbb{C}^m), \quad \text{for } 1 \leq q \leq k.$$

Then,  $u_q$ 's are independent in  $L^2(M, \mathbb{C}^m)$  and we have for  $1 \leq q \leq k$ :

$$H_{\omega, V}^M(u_q) = E W^* \cdot (H_V^M \otimes 1) u \cdot \eta_q = \lambda u_q,$$

using (2.7). This means that  $\lambda$  is also an eigenvalue of  $H_{\omega, V}^M$  with multiplicity greater or equal to  $k$ .

As a consequence, we have:

**THEOREM 2.6.** — *The following three conditions are equivalent:*

- (i)  $\lambda_{\omega, 1}^M = \lambda_{\omega, 2}^M = \dots = \lambda_{\omega, m}^M = \lambda_0^M$ ,
- (ii)  $H_{\omega, V}^M$  and  $H_V^M \otimes 1$  are unitary equivalent,
- (iii) (a):  $K(\omega) = 0$ , (b):  $U_\gamma(\omega) = \{I_m\}$ , for each closed curve  $\gamma$  in  $M$ .

### 3. COMPARISON THEOREMS, CASE II

We look here at the **case II** and we fix a finite system of local trivializations  $(B_\alpha, e_\alpha)_{\alpha \in I}$  of  $E$ , with  $B_\alpha$  connected for each  $\alpha \in I$ . Let  $\omega_\alpha$  be the connection form of  $\nabla$  with respect to  $(B_\alpha, e_\alpha)$  and  $u_0$  (resp.  $\lambda_0^M$ ) the first eigenfunction (resp. eigenvalue of  $H_V^M$  as in Lemma 2.1).

Let us first remark that, using a partition of unity subordinate to the covering  $\{B_\alpha\}_\alpha$ , we can formulate (see [4] for the detail of the proof) this lemma in this case as follow:

$$\text{LEMMA 3.1.} \quad - \quad \|\nabla \xi - du_0 \otimes \xi / u_0\|_1^2 = [H_{\nabla, V}^M(\xi), \xi]_0 - \lambda_0^M \|\xi\|_0^2,$$

for  $\xi \in A_0^0(M, E)$ .

As a consequence of this lemma and the min-max principle [11], we have:

$$\lambda_0^M \leq \lambda_{\nabla}^{M, E}, \quad \text{where } \lambda_{\nabla}^{M, E} \text{ is the first eigenvalue of } H_{\nabla, \nabla}^{M, E}.$$

In order to formulate Proposition 2.1 in this case, we can get using local trivializations the following lemma:

LEMMA 3.2. — *If  $\xi \in A^0(M, E)$ ,  $x_0 \in M$  such that  $\nabla \xi = 0$  and  $\xi(x_0) = 0$ . Then,  $\xi = 0$ .*

Now, let us denote by  $\lambda_{\nabla, 1}^{M, E}, \lambda_{\nabla, 2}^{M, E}, \dots, \lambda_{\nabla, k}^{M, E}$  the  $k$ -first eigenvalues of  $H_{\nabla, \nabla}^{M, E}$ , and recall that  $\nabla$  is supposed compatible with the Hermitian structure of  $E$ . Namely,

$$d\langle \xi, \zeta \rangle = \langle \nabla \xi, \zeta \rangle_0 + \langle \xi, \nabla \zeta \rangle_0, \quad \text{for } \xi, \zeta \in A^0(M, E). \quad (3.1)$$

Then, using (3.1) and Lemma 3.2, we can obtain in the same way as in Proposition 3.2 the:

PROPOSITION 3.3. — *If  $\lambda_{\nabla, 1}^{M, E} = \lambda_{\nabla, 2}^{M, E} = \dots = \lambda_{\nabla, k}^{M, E} = \lambda_0^M$ , then, there exists  $k$ -sections  $(\xi_s)$  of  $E$  over  $M$  such that  $\{\xi_s(x)\}_s$  is an orthonormal system of  $E_x$  for each  $x \in M$ , and that:*

$$\nabla \xi_s = 0 \quad \text{in } A^1(M, E), \quad \text{for } 1 \leq s \leq k. \quad (3.2)$$

COROLLARY 3.4. — Under conditions:  $\lambda_{\nabla, 1}^{M, E} = \lambda_{\nabla, 2}^{M, E} = \dots = \lambda_{\nabla, k}^{M, E} = \lambda_0^M$ , we have:

(i)  $E = \mathcal{X} \oplus \mathcal{X}^\perp$  (Whitney sum), where  $\mathcal{X}$  is a trivialisable subbundle of  $E$  with rank  $k$ ,

(ii)  $\nabla = \nabla_{\mathcal{X}} \oplus \nabla_{\mathcal{X}^\perp}$ , where  $\nabla_{\mathcal{X}}$  is a flat connection on  $\mathcal{X}$  such that:  $U_\gamma(\nabla_{\mathcal{X}}) = \{I_k\}$ , for each closed curve  $\gamma$  in  $M$ .

Let us give the main theorem of this section.

THEOREM 3.5. — The three following conditions are equivalent:

- (i)  $\lambda_{\nabla, 1}^{M, E} = \lambda_{\nabla, 2}^{M, E} = \dots = \lambda_{\nabla, m}^{M, E} = \lambda_0^M$ ,
- (ii)  $\text{Sp}(H_{\nabla, \nabla}^{M, E}) = m \cdot \text{Sp}(H_{\nabla, \nabla}^M)$ ,
- (iii) (a):  $E$  is trivialisable, (b): the curvature of  $\nabla$  vanishes, (c):  $U_\gamma(\nabla) = \{I_m\}$ , for each closed curve  $\gamma$  in  $M$ .

*Proof.* — The implication (ii)  $\Rightarrow$  (i) is trivial.

The assertion (i)  $\Rightarrow$  (iii) follows directly from Corollary 3.4.

Let us prove (iii)  $\Rightarrow$  (ii). We start from (iii) (a) and we consider a family  $\{r_\alpha\}_\alpha$  of applications (*i. e.*, a trivialization of  $E$ ) such that:

$$r_\alpha \in C^\infty(B_\alpha, \mathcal{U}(m)), r_\alpha = g_{\alpha\beta} \cdot r_\beta \quad \text{on } B_{\alpha\beta}, \quad \text{for } \alpha, \beta \in I. \quad (3.3)$$

Let  $(\xi_\alpha)_\alpha$  be the local trivializations of a section  $\xi$  in the system  $(B_\alpha, e_\alpha)$ . By (3.3), we have

$$r_\alpha^* \xi_\alpha = r_\beta^* \xi_\beta \quad \text{on } B_{\alpha\beta}, \quad \text{for } \alpha, \beta \in I. \quad (3.4)$$

Then for each  $\xi \in A^1(M, E)$ , define  $F_\zeta \in C^\infty(M, \mathbb{C}^m)$  by:  $F_{\xi|_{B_\alpha}} = r_\alpha^* \cdot \xi_\alpha$  for  $\alpha \in I$ . It is easy to see that the application  $T$  defined by:  $T(\xi) = F_\xi$  is one to one.

Moreover,

$$\text{supp } \xi = \text{supp } F_\xi, \quad (3.5)$$

$$[\xi, \xi']_0 = \int_M (F_\xi, F_{\xi'}) dv \equiv [F_\xi, F_{\xi'}], \quad \text{for } \xi, \xi' \in A^0(M, E). \quad (3.6)$$

On the other hand, if  $\omega \in A^1(M, \mathcal{M}_{m,d})$  is the connection form of  $\nabla$  [which is trivial by the conditions (b), (c)] with respect to the frame defined by  $(r_\alpha)_\alpha$ , i. e.,

$$\omega|_{B_\alpha} = r_\alpha^* \cdot \omega_\alpha \cdot r_\alpha + r_\alpha^* dr_\alpha;$$

and if  $\mathcal{H}_{\omega, \nabla}^M$  is the Schrödinger operator with magnetic potential  $\omega$ . Then, by a direct computation (and using the min-max principle for the hereunder (C.3) property) we obtain the following properties:

$$(C.1): d\xi_\alpha + \omega_\alpha \xi_\alpha = r_\alpha (dF_\xi + \omega \cdot F_\xi)|_{B_\alpha} \quad \text{for } \alpha \in I,$$

$$(C.2): [H_{\nabla, \nabla}(\xi), \xi']_0 = [\mathcal{H}_{\omega, \nabla}^M(F_\xi), F_{\xi'}], \quad \text{for } \xi, \xi' \in A_0^0(M, E),$$

$$(C.3): \text{Sp}(H_{\nabla, \nabla}^M, E) = \text{Sp}(\mathcal{H}_{\omega, \nabla}^M).$$

Now, the condition (ii) results from (C.3) and Theorem 2.6, respectively.

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