# ZhangXin Chen <br> Analysis of mixed methods using conforming and nonconforming finite element methods 

M2AN - Modélisation mathématique et analyse numérique, tome 27, n ${ }^{\mathrm{o}} 1$ (1993), p. 9-34
[http://www.numdam.org/item?id=M2AN_1993__27_1_9_0](http://www.numdam.org/item?id=M2AN_1993__27_1_9_0)
© AFCET, 1993, tous droits réservés.
L'accès aux archives de la revue «M2AN - Modélisation mathématique et analyse numérique » implique l'accord avec les conditions générales d'utilisation (http://www.numdam.org/conditions). Toute utilisation commerciale ou impression systématique est constitutive d'une infraction pénale. Toute copie ou impression de ce fichier doit contenir la présente mention de copyright.

## Numdam

Article numérisé dans le cadre du programme
Numérisation de documents anciens mathématiques
http://www.numdam.org/

# ANALYSIS OF MIXED METHODS USING CONFORMING AND NONCONFORMING FINITE ELEMENT METHODS (*) 

by Zhangxin Chen ( ${ }^{1}$ )

Communicated by J. Douglas


#### Abstract

An abstract framework under which an equivalence between mixed finite element methods and certain modified versions of conforming and nonconforming finite element methods is established for second order elliptic problems with variable coefficients. It is shown, based on the equivalence, that mixed methods can be implemented through usual conforming or nonconforming methods modified in a cost-free manner and that new error estimates for these methods can be derived. The Raviart-Thomas, Brezzi-Douglas-Marini, and Marini-Pietra mixed methods for second order elliptic problems are analyzed by means of the present techniques.


Résumé. - On établit un cadre abstrait pour établir, dans le cas de problèmes elliptiques du $2^{e}$ ordre à coefficients variables, l'équivalence entre des méthodes d'éléments finis mixtes et certaines versions modifiées de méthodes d'éléments finis conformes et non conformes. En se basant sur cette équivalence, on montre que des méthodes mixtes peuvent être mises en œuvre à partir des méthodes conformes et non conformes habituelles, à moindre coût; on montre également que des nouvelles estimations d'erreur peuvent être obtenues. On analyse par ces techniques, dans le cadre des problèmes elliptiques du $2^{e}$ ordre, les méthodes mixtes de Raviart-Thomas, Brezzi-Douglas-Marini et Marini-Pietra

## 1. INTRODUCTION

It has been observed that in many cases mixed finite element methods give better approximations for the flux variable associated with the solution of a second order elliptic problem than classical Galerkin methods [2], [3], [13].

[^0][^1]However, the mixed formulation is more difficult to handle and, in general, is more expensive from a computational point of view [2], [11]. For second order elliptic problems with piecewise constant coefficients, it has been observed [2], [11] that this drawback can be circumvented by observing a certain equivalence between mixed methods and some modified versions of standard nonconforming methods. Arnold and Brezzi [2] showed, for example, that the Raviart-Thomas mixed method of lowest order is equivalent to the usual $P_{1}$-nonconforming method modified by augmenting the classical $P_{1}$-nonconforming space with $P_{3}$-bubbles and then proved that the equivalence is useful not only for implementing the mixed method but also for deriving error estimates.

Variable coefficients may significantly complicate the equivalence above and thus the performance of the mixed methods. Indeed, in the method of Arnold and Brezzi [2], the weighted averages of the inverse of the coefficients enter the numerical schemes through a projection on the RaviartThomas space.

The main purpose of this paper is to develop, in a rather general setting with minimal hypotheses, error estimates and implementations of mixed finite element methods for second order elliptic problems with variable coefficients. We shall develop an abstract framework under which an equivalence between mixed methods and certain modified versions of conforming and nonconforming finite element methods can be established. It will be shown that our abstract theory includes not only the existing analysis for the Raviart-Thomas method but also provides an approach to the analysis of other mixed methods such as the Brezzi-Douglas-Marini and MariniPietra methods [3], [12]. More specifically, it is proven, by means of the present techniques, that the lowest-order Brezzi-Douglas-Marini and MariniPietra methods are equivalent to modified conventional $P_{3}$-conforming and $P_{2}$-nonconforming finite element methods, respectively. It should be emphasized that the field of application of our abstract results is quite large even through only three families of mixed finite elements are analyzed here. Recently, Arbogast [1] has independently considered many of the same questions with different emphases.

We shall also show that the difficulties with the variable coefficients noted above can be avoided by projecting these coefficients into the finite element space of the scalar variable and that the introduction of the projection of the coefficients in the mixed methods above does not result in a reduction of the order of convergence of the method and can lead to great savings in computational effort. Moreover, the desirable features for piecewise-constant coefficients are shared by the approximation procedure considered for variable coefficienis. In particular, based on the equivalence above, it is proven that the approximate solution for the flux variable produced by both methods can be computed from the solution of the usual conforming or
nonconforming methods in an inexpensive manner, that a superconvergent approximation of the scalar variable by means of post-processing can be obtained, and that new duality error estimates for the methods can be obtained under a certain assumption on the triangulation of the domain.

The rest of this paper is organized as follows. In the next section an abstract theory of an equivalence between mixed methods and modified conventional finite element methods is established. Then, in §3-6, an application of the results to the three families of mixed methods mentioned above is presented. Finally, in $\S 7$, the mixed methods for second order elliptic problems with variable coefficients are discussed.

Throughout this paper we shall use the notation $\|\cdot\|_{s, K}$ and $\|\cdot\|_{s, \infty, K}$ for the norms on the Sobolev spaces $H^{s}(K)$ and $W^{s, \infty}(K)$, respectively, for $s \geqslant 0$ and some set $K \subset \mathbb{R}^{2}$. We shall also denote by $\|\cdot\|_{K}$ and $(., .)_{K}$ the norm and the scalar product on $L^{2}(K)$. The subscript $K$ will be omitted when it is $\Omega$. Finally, the notation $\|\cdot\|_{-s}$ will indicate the norm on the dual space $H^{-s}(\Omega)=\left(H_{0}^{s}(\Omega)\right)^{\prime}, s \geqslant 0$.

## 2. AN ABSTRACT THEORY

In this section we shall first develop an abstract framework under which an equivalence between mixed finite element methods and certain modified versions of conventional finite element methods for (2.2) below can be established. Then, based on the equivalence, we shall obtain a duality error estimate for the methods introduced. In order to fix ideas, in the present and next four sections, the coefficient $a$ will be assumed piecewise-constant. The extension of the results to the case of variable coefficients will be discussed in the last section.

Let $\Omega$ be a domain in $\mathbb{R}^{2}$, let $f \in L^{2}(\Omega)$, and let $a$ be a smooth function on $\boldsymbol{\Omega}$ such that

$$
\begin{equation*}
0<a_{*} \leqslant a \leqslant a^{*}<\infty \quad \text { on } \Omega . \tag{2.1}
\end{equation*}
$$

Consider the Dirichlet problem

$$
\begin{align*}
-\operatorname{div}(a \nabla u) & =f  \tag{2.2a}\\
u & \text { in } \quad \Omega,  \tag{2.2b}\\
& \text { on } \quad \partial \Omega .
\end{align*}
$$

It is well known that problem (2.2) has a unique solution $u$.
Set

$$
V=H(\operatorname{div} ; \Omega)=\left\{\tau \in\left(L^{2}(\Omega)\right)^{2}: \operatorname{div} \tau \in L^{2}(\Omega)\right\},
$$

with the usual norm

$$
\|\tau\|_{V}=\left(\sum_{i=1}^{2}\left\|\tau_{t}\right\|^{2}+\|\operatorname{div} \tau\|^{2}\right)^{1 / 2}
$$

vol. $27, \mathrm{n}^{\circ} 1,1993$
where $\tau=\left(\tau_{1}, \tau_{2}\right)$, and let

$$
W=L^{2}(\Omega)
$$

Introducing $\sigma=-a \nabla u$, a formulation of (2.2) appropriate for the mixed method is then :

$$
\text { Find }(\sigma, u) \in V \times W \text { such that }
$$

$$
\begin{array}{rlrl}
(\alpha \sigma, \tau)-(u, \operatorname{div} \tau) & =0, & & \forall \tau \in V, \\
(\operatorname{div} \sigma, v) & =(f, v), & \forall v \in W, \tag{2.3b}
\end{array}
$$

where $\alpha=a^{-1}$.
For the discretization of (2.3), let $T_{h}=\{T\}$ be a regular partition of $\Omega$ into triangles or rectangles of diameter not greater than $h, 0<h<1$, such that if $T$ is a boundary element, the boundary edge can be curved, and let $E_{h}$ denote the set of edges of triangles or rectangles of $T_{h}$ with the decomposition

$$
E_{h}^{\partial}=\left\{e \in E_{h}: e \in \partial \Omega\right\}, \quad E_{h}^{0}=E_{h} \backslash E_{h}^{\partial}
$$

Associated with $T_{h}$, we introduce the finite element spaces

$$
\begin{array}{ll}
V_{h}=\left\{\tau \in(W)^{2}:\left.\tau\right|_{T} \in V(T),\right. & \left.\forall T \in T_{h}\right\}, \\
W_{h}=\left\{v \in W:\left.v\right|_{T} \in W(T),\right. & \left.\forall T \in T_{h}\right\}
\end{array}
$$

where $V(T)$ and $W(T)$ are finite dimensional, polynomial spaces on $T$ such that

$$
\begin{equation*}
\operatorname{div} V(T) \subseteq W(T) \tag{H1}
\end{equation*}
$$

Note that we do not require that $V_{h} \subset V$. Namely, the normal components of elements in $V_{h}$ are not assumed to be continuous across the interelement boundaries. For simplicity, we assume that there is an integer which bounds the degrees of the polynomials in the finite dimensional spaces introduced in this section.

It is well-known that, when dealing with discretizations of the mixed formulation (2.3), the linear algebraic systems produced by usual mixed finite element methods are generally indefinite. A way to overcome this difficulty is the introduction of Lagrange multipliers on the interelement boundaries in order to relax the continuity requirement on the normal components of the approximate solutions associated with the flux variables across these boundaries [10]. This leads to defining the multiplier space

$$
\Lambda_{h}=\left\{\mu \in L^{2}\left(E_{h}\right):\left.\mu\right|_{e} \in \Lambda(e), \quad \forall e \in E_{h}^{0} ;\left.\quad \mu\right|_{e}=0, \quad \forall e \in E_{h}^{\partial}\right\}
$$

where $\Lambda(e)$ is some polynomial space on the set $e$, and the norm of $\Lambda_{h}$ is given by

$$
\begin{equation*}
|\mu|_{h}^{2}=\sum_{e \in E_{h}^{0}}\|\mu\|_{e}^{2} \tag{2.4}
\end{equation*}
$$

We are now ready to state the mixed-hybrid formulation for approximating the solution of (2.2) [2] :

Find $\left(\sigma_{h}, u_{h}, \lambda_{h}\right) \in V_{h} \times W_{h} \times \Lambda_{h}$ such that

$$
\begin{align*}
\left(\alpha \sigma_{h}, \tau\right)-\sum_{T}\left\{\left(u_{h}, \operatorname{div} \tau\right)_{T}-\left(\lambda_{h}, \tau \cdot n_{T}\right)_{\mathrm{\partial}}\right\}=0, & \forall \tau \in V_{h}  \tag{2.5a}\\
\sum_{T}\left(v, \operatorname{div} \sigma_{h}\right)_{T}=(f, v), & \forall v \in W_{h}  \tag{2.5b}\\
\sum_{T}\left(\mu, \sigma_{h} \cdot n_{T}\right)_{\partial T}=0, & \forall \mu \in \Lambda_{h} \tag{2.5c}
\end{align*}
$$

where $n_{T}$ denotes the outward unit normal to $T$. We shall assume that the problem has a unique solution for each $f \in L^{2}(\Omega)$. This can easily be established [12] under the assumptions (H1) and that for each $v \in W(T)$ and $\mu \in \Lambda(\partial T)$ such that

$$
\|\nabla v\|_{T}+\|\mu\|_{\partial T} \neq 0
$$

there exists $\tau \in V(T)$ satisfying

$$
(\nabla v, \tau)_{T}+\left(\mu, \tau . n_{T}\right)_{\partial T} \neq 0
$$

and that

$$
\gamma_{0}(W(T)) \subseteq \Lambda(\partial T)
$$

where $\Lambda(\partial T)=\prod_{e \in \partial T} \Lambda(e)$ and $\gamma_{u}$ denotes the trace on $\partial T$.
We shall introduce another discrete formulation for approximating the solution of (2.2) which we shall prove to be equivalent to (2.5). To that end, we now define the «intermediate» multiplier space

$$
\tilde{\Lambda}_{h}=\left\{\mu \in L^{2}\left(E_{h}\right):\left.\mu\right|_{e} \in \tilde{\Lambda}(e), \forall e \in E_{h}^{0} ;\left.\mu\right|_{e}=0, \forall e \in E_{h}^{\mathfrak{a}}\right\}
$$

where again $\tilde{\Lambda}(e)$ is a polynomial space on the set $e$ such that

$$
\begin{equation*}
\Lambda(e) \subseteq \tilde{\Lambda}(e), \tau . n_{e} \in \tilde{\Lambda}(e), \tau \in V(T), e \in E_{h}^{0}, T \in T_{h} \tag{H2}
\end{equation*}
$$

with $n_{e}$ being a unit vector normal to $e$. Let now $P_{h}$ and $R_{h}\left(\tilde{R}_{h}\right)$ denote the orthogonal projections onto $W_{h}$ and $\Lambda_{h}\left(\tilde{\Lambda}_{h}\right)$ with respect to the norms vol. 27, n 1,1993
$\|\cdot\|$ and $|\cdot|_{h}$, respectively. Then, define $M_{h} \subset L^{2}(\Omega)$ to be a finite dimensional space such that the following two assumptions are satisfied :

$$
\begin{equation*}
\tilde{R}_{h} \psi \in \Lambda_{h}, \quad \forall \psi \in M_{h} \tag{H3}
\end{equation*}
$$

(H4) For each $v \in W_{h}$ and $\mu \in \Lambda_{h}$, there is a unique $\varphi$ in $M_{h}$ satisfying

$$
P_{h} \varphi=v, \quad \tilde{R}_{h} \varphi=\mu
$$

For a given $\tau \in(W)^{2}$, denote $P_{V} \tau$ the $L^{2}$-projection of $\tau$ in $V_{h}$. Another approximation procedure for (2.2) is then defined by seeking $\psi_{h} \in M_{h}$ such that [2]

$$
\begin{equation*}
\sum_{T}\left(a P_{V}\left(\nabla \psi_{h}\right), \nabla \varphi\right)_{T}=\left(P_{h} f, \varphi\right), \quad \forall \varphi \in M_{h} \tag{2.6}
\end{equation*}
$$

We are now in a position to prove the following equivalence theorem.
ThEOREM 2.1: Assume that assumptions (H1)-(H4) are satisfied. Let ( $\sigma_{h}, u_{h}, \lambda_{h}$ ) be the unique solution of system (2.5) and let $\psi_{h} \in M_{h}$ be determined by

$$
\begin{equation*}
P_{h} \psi_{h}=u_{h}, \quad \tilde{R}_{h} \psi_{h}=\lambda_{h} \tag{2.7}
\end{equation*}
$$

Then $\psi_{h}$ is the unique solution of (2.6). Moreover, $\sigma_{h}$ is related to $\dot{\psi}_{h}$ by

$$
\begin{equation*}
\sigma_{h}=-a P_{V}\left(\nabla \psi_{h}\right) . \tag{2.8}
\end{equation*}
$$

Proof: Note that $\psi_{h}$ is uniquely determined by (H4). From (H1), (H2), (2.7), and Green's formula, (2.5a) becomes

$$
\sum_{T}\left(\nabla \psi_{h}, \tau\right)_{T}+\left(\alpha \sigma_{h}, \tau\right)=0, \quad \forall \tau \in V_{h} .
$$

This shows that $\alpha \sigma_{h}$ is the $L^{2}$-projection of $-\nabla \psi_{h}$ in $V_{h}$ since $\alpha=a^{-1}$ is piecewise constant, and thus (2.8) holds.

Next, by (H1) and (2.5b) with $v=P_{h} \varphi, \varphi \in M_{h}$, we have

$$
\sum_{T}\left(\operatorname{div} \sigma_{h}, \varphi\right)_{T}=\left(P_{h} f, \varphi\right), \quad \forall \varphi \in M_{h}
$$

Hence, by Green's formula,

$$
\begin{aligned}
\sum_{T}\left\{\left(\varphi-\tilde{R}_{h} \varphi, \sigma_{h} \cdot n_{T}\right)_{\partial T}+\left(\tilde{R}_{h} \varphi, \sigma_{h} \cdot n_{T}\right)_{\partial T}-\right. & \left.\left(\sigma_{h}, \nabla \varphi\right)_{T}\right\}= \\
& =\left(P_{h} f, \varphi\right), \quad \forall \varphi \in M_{h}
\end{aligned}
$$

which together with (H2), (H3), and (2.5c) implies that

$$
-\sum_{T}\left(\sigma_{h}, \nabla \varphi\right)_{T}=\left(P_{h} f, \varphi\right), \quad \forall \varphi \in M_{h}
$$

This yields that $\psi_{h}$ is a solution of (2.6) by (2.8).
In order to prove the uniqueness, let $\tilde{\psi}_{h}$ be another solution of (2.6) and define $\left(\tilde{\sigma}_{h}, \tilde{u}_{h}, \tilde{\lambda}_{h}\right)$ by

$$
\begin{align*}
\tilde{\sigma}_{h} & =-a P_{V}\left(\nabla \tilde{\psi}_{h}\right),  \tag{2.9a}\\
\tilde{u}_{h} & =P_{h} \tilde{\psi}_{h},  \tag{2.9b}\\
\tilde{\lambda}_{h} & =\tilde{R}_{h} \tilde{\psi}_{h} . \tag{2.9c}
\end{align*}
$$

By (2.9) and Green's formula, we see that

$$
\begin{aligned}
\left(\alpha \tilde{\sigma}_{h}, \tau\right)- & \sum_{T}\left\{\left(\tilde{u}_{h}, \operatorname{div} \tau\right)_{T}-\left(\tilde{\lambda}_{h}, \tau \cdot n_{T}\right)_{\partial T}\right\} \\
& =-\sum_{T}\left\{\left(\nabla \tilde{\psi}_{h}, \tau\right)_{T}+\left(\tilde{\psi}_{h}, \operatorname{div} \tau\right)_{T}-\left(\tilde{\psi}_{h}, \tau \cdot n_{T}\right)_{\partial T}\right\} \\
& =0, \quad \forall \tau \in V_{h} .
\end{aligned}
$$

by (H1) and (H2). Next, for each $v \in W_{h}$, we define $\chi \in M_{h}$ such that

$$
\begin{equation*}
P_{h} \chi=v, \quad \tilde{R}_{h} \chi=0 \tag{2.10}
\end{equation*}
$$

Then, using (2.10), (H1), Green's formula, (H2), (2.9a), and (2.6),

$$
\begin{aligned}
\sum_{T}\left(\operatorname{div} \tilde{\sigma}_{h}, v\right)_{T} & =\sum_{T}\left(\operatorname{div} \tilde{\sigma}_{h}, \chi\right)_{T} \\
& =\sum_{T}\left\{\left(\tilde{\sigma}_{h}, n_{T}, \chi\right)_{\partial T}-\left(\tilde{\sigma}_{h}, \nabla \chi\right)_{T}\right\} \\
& =\sum_{T}\left(a P_{V}\left(\nabla \tilde{\psi}_{h}\right), \nabla \chi\right)_{T} \\
& =\left(P_{h} f, \chi\right) \\
& =(f, v), \quad \forall v \in W_{h}
\end{aligned}
$$

Finally, for any $\mu \in \Lambda_{h}$, choose $\chi \in M_{h}$ satisfying

$$
\begin{equation*}
P_{h} \chi=0, \quad \tilde{R}_{h} \chi=\mu . \tag{2.11}
\end{equation*}
$$

Hence,

$$
\begin{aligned}
\sum_{T}\left(\tilde{\sigma}_{h}, n_{T}, \mu\right)_{\partial T} & =\sum_{T}\left(\tilde{\sigma}_{h}, n_{T}, \chi\right)_{\partial T} \\
& =\sum_{T}\left\{\left(\tilde{\sigma}_{h}, \nabla \chi\right)_{T}+\left(\operatorname{div} \tilde{\sigma}_{h}, \chi\right)_{T}\right\} .
\end{aligned}
$$

Applying (2.6), (2.9a), (H1), and (2.11), we find that

$$
\sum_{r}\left(\tilde{\sigma}_{h} \cdot n_{T}, \mu\right)_{\partial T}=0, \quad \forall \mu \in \Lambda_{h} .
$$

Combine the results above; thus, $\left(\tilde{\sigma}_{h}, \tilde{u}_{h}, \tilde{\lambda}_{h}\right)$ is a solution of system (2.5). But, by the uniqueness of the solution of (2.5), we see that $\left(\tilde{\sigma}_{h}, \tilde{u}_{h}, \tilde{\lambda}_{h}\right)=\left(\sigma_{h}, u_{h}, \lambda_{h}\right)$, and so, by (H4), $\tilde{\psi}_{h}=\psi_{h}$. This completes the proof of the theorem.

We now turn to the derivation of an error estimate for the problem (2.6). To that end, we shall state further assumptions which will be required in the proof of our next theorem.
(H5) For any smooth function $\phi$ there exists a function $\phi_{h}$ in $M_{h}$ such that

$$
P_{h}\left(\phi-\phi_{h}\right)=0, \quad \tilde{R}_{h}\left(\phi-\phi_{h}\right)=0 .
$$

(H6) For $\varphi \in M_{h}$,

$$
\sum_{e \in E_{h}^{0}}(\mu,[\varphi])_{e}=0, \quad \forall \mu \in \Lambda_{h},
$$

where [ $\varphi$ ] stands for the value of the jump discontinuity of $\varphi$ on the interelement boundaries.

$$
\begin{equation*}
\bar{P}_{0}(e) \subseteq \Lambda(e), \quad e \in E_{h}^{0} \tag{H7}
\end{equation*}
$$

where $P_{0}(e)$ denotes the set of constants on $e$, so that

$$
\begin{equation*}
\left|v-R_{h} v\right|_{h} \leqslant C h^{1 / 2}\left(\sum_{T}\|\nabla v\|_{T}^{2}\right)^{1 / 2}, \quad \forall v \in \prod_{T} H^{1}(T), T \in T_{h} \tag{2.12}
\end{equation*}
$$

The solution $\psi_{h}$ of (2.6) is assumed to satisfy the relation

$$
\begin{equation*}
P_{V}\left(\nabla \psi_{h}\right)=\nabla \psi_{h} \tag{H8}
\end{equation*}
$$

and the approximation property
(H9) $\left\|\nabla\left(u-\psi_{h}\right)\right\|_{h}=\left(\sum_{T}\left\|\nabla\left(u-\psi_{h}\right)\right\|_{T}^{2}\right)^{1 / 2} \leqslant C h^{r-1}\|u\|_{r}$,
for some $r \geqslant 1$.
The domain $\Omega$ will be said to be 2-regular if the Dirichlet problem

$$
\begin{aligned}
-\operatorname{div}(a \nabla p) & =q \quad \text { in } \quad \Omega \\
p=0 & \text { on } \quad \partial \Omega
\end{aligned}
$$

is uniquely solvable for $q \in L^{2}(\Omega)$ and if

$$
\|p\|_{2} \leqslant C\|q\| .
$$

Theorem 2.2: Under the hypotheses of (H1)-(H9), if $u$ and $\psi_{h}$ are the solutions of (2.2) and (2.6), respectively, then,

$$
\begin{equation*}
\left\|u-\psi_{h}\right\| \leqslant C\left(h^{r}\|u\|_{r}+\left\|f-P_{h} f\right\|_{-2}\right), \tag{2.13}
\end{equation*}
$$

provided that $\Omega$ is 2-regular, where $C$ is a generic constant independent of $h$.
Proof: Let $w=u-\psi_{h}$. Consider the auxiliary Dirichlet problem :

$$
\begin{aligned}
& \text { Find } \phi \in H_{0}^{1}(\Omega) \text { such that } \\
& -\operatorname{div}(a \nabla \phi)=w \text { in } \Omega \text {, } \\
& \phi=0 \quad \text { on } \quad \partial \Omega .
\end{aligned}
$$

By the assumed elliptic regularity, we have

$$
\begin{equation*}
\|\phi\|_{2} \leqslant C\|w\| . \tag{2.14}
\end{equation*}
$$

Now,

$$
\begin{align*}
\|w\|^{2} & =\sum_{T}(a \nabla \phi, \nabla w)_{T}-\sum_{T}\left(a \nabla \phi \cdot n_{T}, w\right)_{\partial T}  \tag{2.15}\\
& \equiv R_{1}-R_{2}
\end{align*}
$$

Let $\phi_{h}$ be the function of $\phi$ in $M_{h}$ according to (H5). Then, by (2.2), (2.6), and (H8),

$$
\begin{align*}
R_{1} & =\sum_{T}(a \nabla \phi, \dot{\nabla} w)_{T}  \tag{2.16}\\
& =\left(f-P_{h} f, \phi\right)+\left(P_{h} f, \phi-\phi_{h}\right)-\sum_{T}\left(a \nabla\left(\phi-\phi_{h}\right), \nu \psi_{h}\right)_{T} .
\end{align*}
$$

Since, by (H1), (H2), (H5), and (2.8),

$$
\begin{aligned}
& \sum_{T}\left(a \nabla\left(\phi-\phi_{h}\right), \nabla \psi_{h}\right)_{T}= \\
&=\sum_{T}\left\{\left(\phi-\phi_{h}, \operatorname{div} \sigma_{h}\right)_{T}-\left(\phi-\phi_{h}, \sigma_{h}, n_{T}\right)_{\mathrm{d} T}\right\}=0,
\end{aligned}
$$

and

$$
\left(P_{h} f, \phi-\phi_{h}\right)=0,
$$

we find that

$$
\begin{equation*}
\left|R_{1}\right|=\left|\left(f-P_{h} f, \phi\right)\right| \leqslant\left\|f-P_{h} f\right\|_{-2}\|\phi\|_{2} . \tag{2.17}
\end{equation*}
$$

vol. $27, n^{\circ} 1,1993$

Next, using (H6),

$$
\begin{aligned}
R_{2} & =\sum_{T}\left(a \nabla \phi \cdot n_{T}, w\right)_{\partial T} \\
& =\sum_{T}\left(a\left(\nabla \phi \cdot n_{T}-R_{h}\left\{\nabla \phi \cdot n_{T}\right\}\right), w-R_{h} w\right)_{\partial T}
\end{aligned}
$$

so that, applying (2.12),

$$
\begin{equation*}
\left|R_{2}\right| \leqslant C h\|\phi\|_{2}\|\nabla w\|_{h} . \tag{2.18}
\end{equation*}
$$

Now, combine (2.14)-(2.15), (2.17)-(2.18), and (H9) to obtain the desired result (2.13), and the proof of the theorem has been completed.

Remark : If $\tilde{\Lambda}_{h}=\Lambda_{h}$, the assumption (H5) follows immediately follows from (H4). As seen in the next section, this is the case in most applications. The hypothesis (H6) requires that the elements of $M_{h}$ have a certain continuity across the interelement boundaries. The assumption (H7) is trivially satisfied for all the existing mixed spaces. The relation (H8) may be shown for some mixed spaces under a certain assumption on the triangulation of the domain. The estimate ( H 9 ) can be easily verified by the equivalence above between (2.5) and (2.6) and a known error estimate for the mixed method (2.5). Finally, the duality estimate (2.13) has been proved using the discretization formulation (2.6), which cannot be naturally derived from the original mixed formulation (2.5).

## 3. THE LOWEST-ORDER RAVIART-THOMAS METHOD I

In this section and the next three sections we shall apply the results of the previous section to several examples. We shall consider the lowest-order Raviart-Thomas and Brezzi-Douglas-Marini methods [13], [3] and the mixed method recently introduced by Marini and Pietra [12] since these methods are the most useful in practice. But, as mentioned in the introduction, the results in the previous section can be applied to other more general mixed spaces with higher indexes. For simplicity, we shall assume that $\Omega$ is a convex, polygonal domain. However, it will become clear that the same analysis can also be done for more general domains where $u \in H^{2}(\Omega)$.

Let $T_{h}=\{T\}$ be a triangular decomposition of $\Omega$. The spaces $V(T)$, $W(T), \Lambda(\partial T)$, and $\tilde{\Lambda}(\partial T)$ are defined by

$$
\begin{aligned}
V(T) & =\left(P_{0}(T)\right)^{2}+(x, y) P_{0}(T) \\
W(T) & =P_{0}(T) \\
\Lambda(\partial T) & =\tilde{\Lambda}(\partial T)=P_{0}(\partial T)=\prod_{e \in \partial T} P_{0}(e)
\end{aligned}
$$

where $P_{k}(T)$ denotes the set of polynomials of degree not greater than $k, k \geqslant 0$, on $T$. The assumption (H1) follows from the relation

$$
\operatorname{div} V(T)=W(T)
$$

and it is easy to see that (H2) holds.
We now turn to define the space $M_{h}$. For each $T$ in $T_{h}$, let ( $\lambda_{1}, \lambda_{2}, \lambda_{3}$ ) represent the barycentric coordinates of a point of $T$, and let

$$
N_{3 h}=\left\{v:\left.v\right|_{T}=\gamma_{T} \lambda_{1} \lambda_{2} \lambda_{3}, \gamma_{T} \in \mathbb{R}, \quad \forall T \in T_{h}\right\}
$$

Then, define $M_{h}$ by [2]

$$
M_{h}=M_{N C} \oplus N_{3 h},
$$

where $M_{N C}$ is the usual nonconforming space; i.e.,

$$
M_{N C}=\left\{v \in L^{2}(\Omega):\left.v\right|_{T} \in P_{1}(T), \quad \forall T \in T_{h},\right.
$$

$v$ is continuous at the midpoints of sides in $E_{h}^{0}$ and vanishes at the midpoints of sides in $\left.E_{h}^{\mathrm{a}}\right\}$.
Note that the space $M_{h}$ is the classical $P_{1}$-nonconforming space augmented with $P_{3}$-bubbles. The hypothesis (H3) is trivial, since $\Lambda_{h}=\tilde{\Lambda}_{h}$ and thus $R_{h}=\tilde{R}_{h}$. The assumption (H4) was shown in [2].

Consequently, Theorem 2.1 is applicable and shows that the lowest-order Raviart-Thomas method is equivalent to a modified $P_{1}$-nonconforming method. It follows from (2.8) that

$$
\sigma_{h}=-a P_{V}\left(\nabla \psi_{h}\right) .
$$

Furthermore, it can be shown [11] using the equivalence between systems (2.5) and (2.6) that the approximate solution $\sigma_{h}$ can be computed by the simple formula

$$
\begin{equation*}
\sigma_{h}(x)=-a \nabla z_{h}+\left(P_{h} f\right)_{T}\left(x-x_{T}\right) / 2, \quad x \in T \tag{3.1}
\end{equation*}
$$

where $x_{T}$ is the barycenter of the triangle $T$, and $z_{h} \in M_{N C}$ is the solution of

$$
\begin{equation*}
\sum_{T}\left(a \nabla z_{h}, \nabla v\right)_{T}=\left(P_{h} f, v\right), \quad \forall v \in M_{N C} \tag{3.2}
\end{equation*}
$$

Since it is easily seen from the definition of $N_{3 h}$ that the assumption (H8) is not valid, the duality estimate (2.13) cannot be derived naturally in the present case.

## 4. THE LOWEST-ORDER RAVIART-THOMAS METHOD II

In this section we shall reanalyze the lowest-order Raviart-Thomas method by means of modifying the space $M_{h}$. We shall show that, while the features of the previous section are preserved here, the new approach allows for Theorem 2.2 to be used to derive a new duality estimate for the scalar variable.

The spaces $V_{h}, W_{h}, \Lambda_{h}$, and $\tilde{\Lambda}_{h}$ are defined as in the previous section, but $M_{h}$ is modified as follows. On the triangle $T$ there exists a quadratic function (unique up to a multiplicative constant) $\phi_{0, T}(x)$ which vanishes at the two Gaussian quadrature points of each side of $T$. It can be written explicitly as [9]

$$
\phi_{0, T}(x)=2-3\left(\lambda_{1}^{2}+\lambda_{2}^{2}+\lambda_{3}^{2}\right)
$$

which has been scaled so that its value is unity at the barycenter of $T$. Then, we introduce

$$
N_{2 h}=\left\{v:\left.v\right|_{T}=\gamma_{T} \phi_{0, T}(x), \quad \gamma_{T} \in \mathbb{R}, \quad \forall T \in T_{h}\right\}
$$

and

$$
M_{h}=M_{N C} \oplus N_{2 h}
$$

where $M_{N C}$ is defined as in the previous section. $M_{h}$ is now the usual nonconforming space augmented with the $P_{2}$-bubble functions.

Again, (H3) is trivially valid and (H4) can be seen from the next iemma.
LEMMA 4.1: Let $T \in T_{h}$ be a triangle with edges $e_{l}(i=1,2,3)$. Then for all $p_{i} \in L^{2}\left(e_{l}\right) \quad(i=1,2,3)$ and $q \in L^{2}(T)$, there exists a unique $\chi \in M(T)=\left\{\left.v\right|_{T}: v \in M_{h}\right\}$ satisfying

$$
\begin{align*}
& \left(\chi-p_{i}, 1\right)_{e_{i}}=0, \quad i=1,2,3  \tag{4.1}\\
& (\chi-q, 1)_{T}=0 \tag{4.2}
\end{align*}
$$

Proof: Clearly, the system given by (4.1) and (4.2) is a square linear system with four equations and unknowns. Hence, to prove existence and uniqueness of $\chi$, it suffices to show that $\chi=0$ if $q=0$ and $p_{i}=0$ ( $i=1,2,3$ ).
Let $\chi=\chi_{1}+\gamma_{T} \phi_{0, T}$ such that $\chi_{1} \in M_{N C}$ and $\gamma_{T} \in \mathbb{R}$. Then, conditions (4.1) with $p_{i}=0$ and the vanishing of the average value of $\phi_{0, T}$ on each edge imply that

$$
\left(\chi_{1}, 1\right)_{e_{t}}=0, \quad i=1,2,3
$$

Consequently, it follows that $\chi_{1}=0$. As a result of this, $\chi=\gamma_{T} \phi_{0, T}$. Hence, by (4.2) with $q=0, \gamma_{T}=0$ and $\chi=0$, and the proof has been completed.

Consequently, Theorem 2.1 shows again that the Raviart-Thomas method of lowest-order is equivalent to a modified $P_{1}$-nonconforming method amplified by $P_{2}$-bubbles this time. Moreover, based on the present equivalence, it can also be shown that the simple implementation (3.1) for $\sigma_{h}$ is preserved here [4].

To apply Theorem 2.2, we must check that hypotheses (H5)-(H9) are valid. First, notice that (H5) is valid by Lemma 4.1. Next, it is immediate from the definition of $M_{h}$ that the jumps of elements in $M_{h}$ have zero mean values on interelement boundaries, so that (H6) is valid. Also, (H7) is obviously satisfied. It thus remains to check (H8) and (H9). For (H8), we need the next result.

Lemma 4.2 : If all triangles in the triangulation $T_{h}$ of $\Omega$ are equilateral, then

$$
\begin{equation*}
P_{V}(\nabla v)=\nabla v, \quad \forall v \in M_{h} . \tag{4.3}
\end{equation*}
$$

Proof: It suffices to prove (4.3) for the $P_{2}$-bubbles by the definition of $M_{h}$. Let $T \in T_{h}$ be an equilateral triangle with vertices $\left(x_{1}, y_{1}\right),\left(x_{2}, y_{2}\right)$, and $\left(x_{3}, y_{3}\right)$. Since $\lambda_{3}=1-\lambda_{1}-\lambda_{2}$, it follows from the definition of $\phi_{0, T}(x)$ that

$$
\phi_{0, T}(x)=-1+6\left(\lambda_{1}+\lambda_{2}\right)-6\left(\lambda_{1}^{2}+\lambda_{1} \lambda_{2}+\lambda_{2}^{2}\right) .
$$

Hence, it suffices to consider the function

$$
\phi(x)=\lambda_{1}^{2}+\lambda_{1} \lambda_{2}+\lambda_{2}^{2} .
$$

Let

$$
\begin{gathered}
a_{1}=y_{2}-y_{3}, \quad a_{2}=y_{3}-y_{1}, \\
b_{1}=x_{3}-x_{2}, \quad b_{2}=x_{1}-x_{3}, \\
D=\left|a_{1} b_{2}-a_{2} b_{1}\right| .
\end{gathered}
$$

Then, a calculation shows that

$$
\nabla \phi=\left(\phi_{11} x+\phi_{12} y+\phi_{1}, \phi_{12} x+\phi_{22} y+\phi_{2}\right),
$$

where

$$
\begin{aligned}
& \phi_{11}=2\left(a_{1}^{2}+a_{1} a_{2}+a_{2}^{2}\right) / D^{2}, \\
& \phi_{12}=\left(2 a_{1} b_{1}+a_{1} b_{2}+a_{2} b_{1}+2 a_{2} b_{2}\right) / D^{2}, \\
& \phi_{22}=2\left(b_{1}^{2}+b_{1} b_{2}+b_{2}^{2}\right) / D^{2},
\end{aligned}
$$

and $\phi_{1}$ and $\phi_{2}$ are some constants. Since $T$ is equilateral, it can be easily calculated that $\phi_{12}=0$ and $\phi_{11}=\phi_{22}=3 h_{T}^{2} / 2 D^{2}$. That is, $\nabla \phi \in V(T)$, and the proof is completed.

As a result of (4.3), (H8) is valid and the system (2.6) may be rewritten as :
Find $\psi_{h} \in M_{h}$ such that

$$
\begin{equation*}
\sum_{T}\left(a \nabla \psi_{h}, \nabla v\right)_{T}=\left(P_{h} f, v\right), \quad \forall v \in M_{h} \tag{4.4}
\end{equation*}
$$

Hence, we see that the method is just a slightly modified version of the usual nonconforming method.

THEOREM 4.3: If $u$ and $\psi_{h}$ are the solutions of (2.2) and (4.4), respectively, and all triangles in the triangulation $T_{h}$ of $\Omega$ are equilateral, then,

$$
\begin{gather*}
\left\|\nabla u-\nabla \psi_{h}\right\|_{h}=\left(\sum_{T}\left\|\nabla u-\nabla \psi_{h}\right\|_{T}^{2}\right)^{1 / 2} \leqslant C h\|u\|_{2}  \tag{4.5}\\
\left\|u-\psi_{h}\right\| \leqslant C h^{2}\|f\|_{1} \tag{4.6}
\end{gather*}
$$

with $C$ independent of $u$ and $h$.
Proof: First, (4.5) follows directly from the relation

$$
\sigma_{h}=-a \nabla \psi_{h},
$$

and the known error estimate [2], [7]

$$
\left\|\sigma-\sigma_{h}\right\| \leqslant C h\|u\|_{2}
$$

consequently, (H9) is satisfied with $r=2$. Hence, apply Theorem 2.2 and an approximation property of $P_{h}$ to obtain

$$
\begin{aligned}
\left\|u-\psi_{h}\right\| & \leqslant C\left(h^{2}\|u\|_{2}+\left\|f-P_{h} f\right\|_{-2}\right) \\
& \leqslant C\left(h^{2}\|u\|_{2}+\left\|f-P_{h} f\right\|_{-1}\right) \\
& \leqslant C h^{2}\|f\|_{1} .
\end{aligned}
$$

and the proof is complete.

## 5. THE LOWEST-ORDER BREZZI-DOUGLAS-MARINI METHOD

Let $T_{h}=\{T\}$ be again a decomposition of $\Omega$ into triangles, and set

$$
\begin{aligned}
& V(T)=\left(P_{1}(T)\right)^{2} \\
& W(T)=P_{0}(T) \\
& \Lambda(\partial T)=\tilde{\Lambda}(\partial T)=\prod_{e \in \partial T} P_{1}(e) .
\end{aligned}
$$

The assumptions (H1) and (H2) hold trivially.
In order to introduce $M_{h}$, let $\Gamma_{h}$ denote the collection of the vertices of triangles of $T_{h}$, and let $\tilde{\Delta}$ be a function from $\Gamma_{h}$ into $\mathbb{R}$. Then, define

$$
M_{h}^{\tilde{\Delta}}=\left\{v \in C^{0}(\bar{\Omega}):\left.v\right|_{T} \in P_{3}(T), \forall T \in T_{h} ; v(i)=\tilde{\Delta}(i), \forall i \in \Gamma_{h}\right\}
$$

and

$$
M_{h}=H_{0}^{1}(\Omega) \cap M_{h}^{\tilde{\Delta}}
$$

Observe that $M_{h}$ depends on the function $\tilde{\Delta}$. For example, let $\tilde{\Delta}$ be the zero function on $\Gamma_{h}$; then,

$$
\begin{aligned}
M_{h}^{0}= & \left\{v \in C^{0}(\bar{\Omega}):\left.v\right|_{T} \in P_{3}(T), \forall T \in T_{h} ; v(i)=0, \forall i \in \Gamma_{h}\right\}, \\
= & \left\{v \in C^{0}(\bar{\Omega}):\left.v\right|_{T} \in \operatorname{span}\left\{9 \lambda_{\imath} \lambda_{j}\left(3 \lambda_{\imath}-1\right) / 2,27 \lambda_{1} \lambda_{2} \lambda_{3},\right.\right. \\
& \left.i, j=1,2,3, \quad i \neq j\}, \forall T \in T_{h}\right\} .
\end{aligned}
$$

The next lemma is just one form of the standard uniqueness theorem for determining a cubic polynomial on a triangle.

Lemma 5.1. Let $\tilde{\Delta}$ be any given function on $\Gamma_{h}$. Then for $v \in W_{h}$ and $\mu \in \Lambda_{h}$, there is a unique $\chi \in M_{h}$ such that

$$
P_{h} \chi=v, \quad R_{h} \chi=\mu
$$

As a consequence of the lemma, the assumption (H4) is valid. Therefore, as $\Lambda_{h}=\tilde{\Lambda}_{h}$ and thus (H3) is true, we apply Theorem 2.1 to conclude that the lowest-order Brezzi-Douglas-Marini method is equivalent to a modified $P_{3}$-conforming finite element method.

Note that there is no simple formula for the computation of the approximate solution $\sigma_{h}$ produced by the Brezzi-Douglas-Marini method like (3.1) in the case of the Kaviart-Thomas method. This may account for a difference in the computation between these two mixed methods.

## 6. THE MARINI-PIETRA METHOD

With $T_{h}$ defined as before, for each triangle $T \in T_{h}$ we let

$$
\begin{aligned}
& V(T)=\operatorname{span}\left\{\tau^{1}, \tau^{2}, \tau^{3}\right\}, \\
& W(T)=P_{0}(T), \\
& \Lambda(\partial T)=P_{0}(\partial T), \\
& \tilde{\Lambda}(\partial T)=P_{1}(\partial T),
\end{aligned}
$$

where

$$
\tau^{1}=(1,0), \quad \tau^{2}=(0,1), \quad \tau^{3}=\left(\tau_{1}^{3}, \tau_{2}^{3}\right)
$$

vol. 27, n ${ }^{\circ} 1,1993$
such that $\tau^{3} \in\left(P_{1}(T)\right)^{2}$ and, for a chosen edge $e$ of $T$,

$$
\begin{gather*}
\left(\tau^{3} \cdot n_{e}, 1\right)_{e}=1  \tag{6.1a}\\
\left(\tau^{3} \cdot n_{e}, 1\right)_{\tilde{e}}=0, \quad \forall \tilde{e} \neq e  \tag{6.1b}\\
\left(\tau_{1}^{3}, 1\right)_{T}=\left(\tau_{2}^{3}, 1\right)_{T}=0 \tag{6.1c}
\end{gather*}
$$

The condition (6.1) determines a one-dimensional manifold; $\tau^{3}$ can be chosen arbitrarily as an element of this manifold. In particular, $\tau^{3}$ can be chosen as the element of minimum norm [12], for example.

From the choice above, it is obvious that (H1) and (H2) are valid.
To construct $M_{h}$, let $R_{h}=R_{h}^{0}$ and $\tilde{R}_{h}=R_{h}^{1}$ indicate the usual orthogonal projections onto $\Lambda_{h}$ and $\tilde{\Lambda}_{h}$, respectively, with respect to the norm |. $\left.\right|_{h}$. For each $T \in T_{h}$, let $a_{i 1}^{T}$ and $a_{i 2}^{T}$ be the two Gaussian quadrature points of each side $e_{i}$ of $T, i=1,2,3$, and let

$$
M(T)=\left\{v: v \in P_{2}(T), v\left(a_{11}^{T}\right)=v\left(a_{t 2}^{T}\right), \quad i=1,2,3\right\}
$$

Note that, since the six nodal values satisfy [9]

$$
\sum_{t=1}^{3}\left\{v\left(a_{t 2}^{T}\right)-v\left(a_{t 1}^{T}\right)\right\}=0, \quad \forall v \in P_{2}(T)
$$

$\dot{M}(T)$ is four-dimensionai. Then, we introduce

$$
M_{h}=\left\{v \in L^{2}(\Omega):\left.v\right|_{T} \in M(T), \quad \forall T \in T_{h},\right.
$$

$v$ is continuous at the two Gaussian quadrature points of sides in $E_{h}^{0}$ and vanishes at the two Gaussian quadrature points of sides in $\left.E_{h}^{\boldsymbol{p}}\right\}$.
Hence, $M_{h}$ is a modified $P_{2}$-nonconforming space.
Lemma 6.1: Let $T \in T_{h}$ be a triangle with sides $e_{i}(i=1,2,3)$. Then for any $p_{t} \in P_{0}\left(e_{t}\right)(i=1,2,3)$ and $q \in L^{2}(T)$, there exists a unique $\chi \in M(T)$ such that

$$
\begin{align*}
& \left(\chi-p_{\imath}, z\right)_{e_{i}}=0, \quad \forall z \in P_{1}\left(e_{\imath}\right), \quad i=1,2,3  \tag{6.2a}\\
& (\chi-q, 1)_{T}=0 \tag{6.2b}
\end{align*}
$$

Remark : By the definition of $M(T)$, equation (6.2a) in fact has only three linearly independent equations. Also, as the $P_{2}$-bubble $\phi_{0, T}(x)$ vanishes at the six Gaussian quadrature points, $(6.2 b)$ is needed to uniquely determine $\chi \in M(T)$.

Proof of Lemma 6.1 : Clearly, by the definition of $M(T)$, the system given by (6.2) is a square linear system with four equations and unknowns. Hence, to prove existence and uniqueness of $\chi$, it suffices to show that $\chi=0$ if $q=0$ and $p_{i}=0(i=1,2,3)$.

First, condition ( $6.2 a$ ) with $p_{l}=0$ implies that

$$
(\chi, z)_{e_{i}}=0, \quad \forall z \in P_{1}\left(e_{\imath}\right), \quad i=1,2,3 ;
$$

consequently, there is $\gamma_{T} \in \mathbb{R}$ such that $\chi=\gamma_{T} \phi_{0, T}(x)$. Hence, by (6.2b) with $q=0, \gamma_{T}=0$. Namely, $\chi=0$, and the proof has been completed.

It now becomes apparent that the hypothesis (H4) is just a simple consequence of the above lemma.

Lemma 6.2: For all $v \in M_{h}$,

$$
R_{h}^{1} v=R_{h}^{0} v
$$

This result is immediate from the definition of $M_{h}$, and so (H3) is satisfied.
Using Theorem 2.1, we see that the Marini-Pietra method is equivalent to a modified $P_{2}$-nonconforming method.

Set

$$
X_{h}=\left\{v \in C^{0}(\bar{\Omega}):\left.v\right|_{T} \in M(T), \forall T \in T_{h},\left.v\right|_{\partial \Omega}=0\right\}
$$

Then, it is interesting to note that [9]

$$
M_{h}=X_{h} \oplus N_{2 h},
$$

where $N_{2 h}$ is defined as in the second example.

## 7. VARIABLE COEFFICIENTS

In this section we shall briefly extend the results of the previous sections un piecewise constants to the case of variabie coefficients. Às an example, we shall consider the lowest-order Raviart-Thomas method in detail. Other methods can be analyzed analogously. For more information on variable coefficient mixed finite element methods, refer to [4].

### 7.1. Basic error estimates

Let $T_{h}=\{T\}$ be a quasiregular partition of $\Omega$ into triangles. Set

$$
\bar{V}_{h}=\left\{\tau \in V:\left.\tau\right|_{T} \in V(T), \quad \forall T \in T_{h}\right\}
$$

with $V(T)$ defined as in § 3, and set $\alpha_{h}=P_{h} \alpha$. Note that the normal components of each element in $\bar{V}_{h}$ are now required to be continuous across the interelement boundaries.

We now introduce a modified mixed formulation for approximating the solution of (2.2) :

Find $\left(\bar{\sigma}_{h}, \bar{u}_{h}\right) \in \bar{V}_{h} \times W_{h}$ such that

$$
\begin{array}{lll}
\left(\alpha_{h} \bar{\sigma}_{h}, \tau\right)-\left(\operatorname{div} \tau, \bar{u}_{h}\right) & =0, & \\
\left(\operatorname{div} \bar{\sigma}_{h}, v\right) & & \forall \tau \in \bar{V}_{h},  \tag{7.1b}\\
(f, v), & \forall v \in W_{h},
\end{array}
$$

where we have projected the coefficient into the space $W_{h}$. Observe that, when $a$ is piecewise constant, (7.1) is the standard mixed finite element method [2], [7].

We now state some error estimates.
THEOREM 7.1 : Problem (7.1) has a unique solution ( $\bar{\sigma}_{h}, \bar{u}_{h}$ ). Moreover, there exists $C>0$, independent of $h$, such that

$$
\begin{align*}
\left\|\sigma-\bar{\sigma}_{h}\right\| & \leqslant C\left(\left\|\alpha-\alpha_{h}\right\|+\left\|\sigma-\Pi_{h} \sigma\right\|\right)  \tag{7.2}\\
\left\|\operatorname{div}\left(\sigma-\bar{\sigma}_{h}\right)\right\| & \leqslant C\left\|\left(I-P_{h}\right) \operatorname{div} \sigma\right\| \leqslant C h\|f\|_{1},  \tag{7.3}\\
\left\|u-\bar{u}_{h}\right\| & \leqslant C h\left(\|u\|_{2}+|\alpha|_{1}\right)  \tag{7.4}\\
\left\|P_{h} u-\bar{u}_{h}\right\| & \leqslant C_{1} h^{2}\left(\|u\|_{2}+\|f\|_{1}+|\alpha|_{1}\right) \tag{7.5}
\end{align*}
$$

where $(\sigma, u)$ is the solution of (2.3), I is the identity operator, $C_{1}=C_{1}\left(\|a\|_{1, \infty}\right)$, and $\Pi_{h}$ will be defined below.

Note that it follows from (2.1) that

$$
\begin{equation*}
\alpha_{h} \geqslant\left(a^{*}\right)^{-1}>0 \tag{7.6}
\end{equation*}
$$

so that the existence and uniqueness of a solution to (7.1) can be demonstrated in a standard way (see, e.g., [7]). Estimates (7.2)-(7.5) can be obtained by making use of the duality ideas of Douglas and Roberts [7]. We shall here use a more direct approach to obtain these estimates. This approach is easy to understand and is simpler than that given in [7], [8], [13].

Let $\Pi_{h}: H^{1}(\Omega) \rightarrow \bar{V}_{h}$ be the Raviart-Thomas projection, [7], [13], which satisfies

$$
\begin{align*}
\left\|\tau-\Pi_{h} \tau\right\| & \leqslant C\|\tau\|_{1} h, \quad \tau \in\left(H^{1}(\Omega)\right)^{2}  \tag{7.7}\\
\left\|\operatorname{div}\left(\tau-\Pi_{h} \tau\right)\right\| & \leqslant C\|\operatorname{div} \tau\|_{1} h, \quad \tau \in\left(H^{1}(\Omega)\right)^{2}, \operatorname{div} \tau \in H^{1}(\Omega)  \tag{7.8}\\
\operatorname{div} \Pi_{h} & =P_{h} \operatorname{div},\left(H^{1}(\Omega)\right)^{2} \rightarrow W_{h} \tag{7.9}
\end{align*}
$$

We shall also require the approximation property

$$
\begin{equation*}
\left\|v-P_{h} v\right\|_{-s} \leqslant C\|v\|_{1} h^{s+1}, \quad s=0, \mathbf{1} \tag{7.10}
\end{equation*}
$$

if $v \in H^{1}(\Omega)$.

Proof of Theorem 7.1: Let $\quad x=\sigma-\bar{\sigma}_{h}=\left(\sigma-\Pi_{h} \sigma\right)+$ $\left(\Pi_{h} \sigma-\bar{\sigma}_{h}\right)=y+z$ and $\xi=u-\bar{u}_{h}=\left(u-P_{h} u\right)+\left(P_{h} u-\bar{u}_{h}\right)=\eta+\zeta$. These errors satisfy the equations given by subtracting (7.1) from (2.3) and using (7.9) :

$$
\begin{align*}
\left(\alpha_{h} x, \tau\right)-(\zeta, \operatorname{div} \tau) & =\left(\left\{\alpha_{h}-\alpha\right\} \sigma, \tau\right), & & \forall \tau \in \bar{V}_{h}  \tag{7.11a}\\
(v, \operatorname{div} z) & =0, & & \forall v \in W_{h} \tag{7.11b}
\end{align*}
$$

Take the test functions $\tau=z$ in (7.11a) and $v=\zeta$ in (7.11b) and add to have

$$
\begin{aligned}
\left(\alpha_{h} x, z\right) & =\left(\left\{\alpha_{h}-\alpha\right\} \sigma, z\right) \\
& \leqslant C\left\|\alpha_{h}-\alpha\right\|^{2}+\varepsilon\|z\|^{2}
\end{aligned}
$$

where $\varepsilon$ is a positive constant which may be taken as small as we please. Consequently, since $\left(\alpha_{h} z, z\right)=\left(\alpha_{h} x, z\right)-\left(\alpha_{h} y, z\right)$, it follows that

$$
\|z\|^{2} \leqslant C\left\{\left\|\alpha_{h}-\alpha\right\|^{2}+\|y\|^{2}\right\}+\varepsilon\|z\|^{2}
$$

and that

$$
\|x\| \leqslant\|y\|+\|z\| \leqslant C\left(\left\|\alpha_{h}-\alpha\right\|+\|y\|\right) ;
$$

i.e., (7.2) holds.

Next, (7.11b) shows that $\operatorname{div} \bar{\sigma}_{h}=P_{h} \operatorname{div} \sigma$; consequently,

$$
\|\operatorname{div} x\|=\left\|\operatorname{div} \sigma-P_{h} \operatorname{div} \sigma\right\|,
$$

which gives (7.3) by (7.10) with $s=0$.
Now, let $\tilde{\tau} \in \bar{V}_{h}$ in (7.11a) be a function associated with $\zeta$ such that [13]

$$
\operatorname{div} \tilde{\tau}=\zeta \quad \text { and } \quad\|\tilde{\tau}\|_{H(\operatorname{div} ; \Omega)} \leqslant C\|\zeta\| ;
$$

then,

$$
\begin{aligned}
\|\zeta\|^{2} & =(\zeta, \operatorname{div} \tilde{\tau})=\left(\alpha_{h} x, \tilde{\tau}\right)-\left(\left\{\alpha_{h}-\alpha\right\} \sigma, \tilde{\tau}\right) \\
& \leqslant C\left\{\|x\|^{2}+\left\|\alpha_{h}-\alpha\right\|^{2}\right\}+\varepsilon E\|\tilde{\tau}\|^{2} \\
& \leqslant C\left\{\|x\|^{2}+\left\|\alpha_{h}-\alpha\right\|^{2}\right\}+\varepsilon E\|\zeta\|^{2}
\end{aligned}
$$

i.e., by (7.2), (7.7), and (7.10), (7.4) holds.

In order to show (7.5), we rewrite (7.11) as

$$
\begin{align*}
(\alpha x, \tau)-(\zeta, \operatorname{div} \tau) & =\left(\left\{\alpha_{h}-\alpha\right\} \bar{\sigma}_{h}, \tau\right), & & \forall \tau \in \bar{V}_{h}  \tag{7.12a}\\
(v, \operatorname{div} x) & =0, & & \forall v \in W_{h} \tag{7.12b}
\end{align*}
$$

vol. $27, n^{\circ} 1,1993$

Note that, by (7.2), (7.7), (710), quasiregulanty of $T_{h}$, and the boundedness of $\Pi_{h}$,

$$
\begin{align*}
\left\|\bar{\sigma}_{h}\right\|_{\infty} & \leqslant\|z\|_{\infty}+\left\|\Pi_{h} \sigma\right\|_{\infty} \\
& \leqslant C h^{-1}(\|x\|+\|y\|)+\left\|\Pi_{h} \sigma\right\|_{\infty}  \tag{7.13}\\
& \leqslant C
\end{align*}
$$

We are now in a position to prove the estimate (7.5) by means of a duality argument different from that given in [7].

Let $\psi \in L^{2}(\Omega)$, and let $\varphi \in H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$ be such that $-\operatorname{div}(a \nabla \varphi)=\psi$. then, it follows from (79) and (712) that with $\beta=\alpha_{h}-\alpha:$

$$
\begin{aligned}
(\zeta, \psi)= & (\zeta,-\operatorname{div}(a \nabla \varphi)) \\
= & \left(\zeta,-\operatorname{div}\left(\Pi_{h}\{a \nabla \varphi\}\right)\right) \\
= & \left(\beta \bar{\sigma}_{h}, \Pi_{h}\{a \nabla \varphi\}\right)-\left(\alpha x, \Pi_{h}\{a \nabla \varphi\}\right) \\
= & \left(\beta \bar{\sigma}_{h}, \Pi_{h}\{a \nabla \varphi\}-a \nabla \varphi\right)-(\beta x, a \nabla \varphi)+(\beta \sigma, a \nabla \varphi) \\
& +\left(\alpha x, a \nabla \varphi-\Pi_{h}\{a \nabla \varphi\}\right)+\left(\operatorname{div} x, \varphi-P_{h} \varphi\right),
\end{aligned}
$$

so that, by $(77),(710),(713)$, and the approximation property

$$
\|\beta\|_{\infty} \leqslant C h\|\alpha\|_{1 \infty}
$$

we have

$$
\begin{aligned}
|(\zeta, \psi)| \leqslant & \|\beta\|\left\|\bar{\sigma}_{h}\right\|_{\infty}\left\|\Pi_{h}\{a \nabla \varphi\}-a \nabla \varphi\right\| \\
& +\|\beta\|_{\infty}\|x\|\|a \nabla \varphi\|+\|\beta\|-1\|a \sigma \nabla \varphi\|_{1} \\
& +\|\alpha\|_{\infty}\|x\|\left\|a \nabla \varphi-\Pi_{h}\{a \nabla \varphi\}\right\|+\|\operatorname{div} x\|\left\|\varphi-P_{h} \varphi\right\| \\
\leqslant & C_{1}\left(h\|\beta\|+\|\beta\|_{-1}+h\|x\|+h\|\operatorname{div} x\|\right)\|\varphi\|_{2},
\end{aligned}
$$

with $C_{1}=C_{1}\left(\|a\|_{1 \infty}\right)$. Finally, combine (72), (7.3), (77), (710), and the assumed ellıptic regularity to obtain (7.5), and the proof has been completed.

### 7.2. Post-Processing

Associated with (2.5), we have the modified mixed-hybrid method.
Find $\left(\bar{\sigma}_{h}, \bar{u}_{h}, \bar{\lambda}_{h}\right) \in V_{h} \times W_{h} \times \Lambda_{h}$ such that

$$
\begin{array}{rlrl}
\left(\alpha_{h} \bar{\sigma}_{h}, \tau\right)-\sum_{T}\left\{\left(\bar{u}_{h}, \operatorname{div} \tau\right)_{T}-\left(\bar{\lambda}_{h}, \tau \cdot n_{T}\right)_{\partial T}\right\} & =0, & \forall \tau \in V_{h}, \\
\sum_{T}\left(v, \operatorname{div} \bar{\sigma}_{h}\right)_{T}=(f, v), & \forall v \in W_{h}, \\
\sum_{T}\left(\mu, \bar{\sigma}_{h} \cdot n_{T}\right)_{\partial T} & =0, & \forall \mu \in \Lambda_{h} . \tag{714c}
\end{array}
$$

Again, by (7.6), the existence and uniqueness of solution to (7.14) can be easily shown in a standard way. Furthermore, since equation ( $7.14 c$ ) imposes the required continuity on $\bar{\sigma}_{h}$, the pair $\left(\bar{\sigma}_{h}, \bar{u}_{h}\right)$ of (7.14) is the same as that of (7.1).

Set

$$
|\mu|_{-1 / 2, h}^{2}=\sum_{e \in E_{h}^{0}}\|\mu\|_{e}^{2} h_{e}
$$

The following lemma can be proved by the argument given in [2].
LEMMA 7.2 : There is a constant $C$, independent of $u$ and $h$, such that, for every $T \in T_{h}$ and every edge $e$ of $T$,

$$
\begin{gather*}
\left\|\bar{\lambda}_{h}-R_{h} u\right\|_{e} \leqslant C\left(h_{T}^{1 / 2}\|x\|_{T}+h_{T}^{-1 / 2}\|\zeta\|_{T}+h_{T}^{1 / 2}\left\|\alpha-\alpha_{h}\right\|_{T}\right)  \tag{7.15}\\
\left|\bar{\lambda}_{h}-R_{h} u\right|_{-1 / 2, h} \leqslant C\left(h\|x\|+\|\zeta\|+h\left\|\alpha-\alpha_{h}\right\|\right) \tag{7.16}
\end{gather*}
$$

As mentioned before, the advantage of the system associated with (7.14) is that the stiffness matrix is positive definite. Moreover, the multiplier above can be used to obtain by post-processing an approximate solution to $u$ which is asymptotically more accurate than the approximation $\bar{u}_{h}$.

THEOREM 7.3: Let

$$
W_{h}^{1}=\left\{v \in L^{2}(\Omega):\left.v\right|_{T} \in P_{1}(T), \quad \forall T \in T_{h}\right\}
$$

and let $u_{h}^{*} \in W_{h}^{1}$ be defined by

$$
\begin{equation*}
R_{h} u_{h}^{*}=\bar{\lambda}_{h} \tag{7.17}
\end{equation*}
$$

Then,

$$
\begin{equation*}
\left\|u-u_{h}^{*}\right\| \leqslant C_{1} h^{2}\left(\|u\|_{2}+\|f\|_{1}+|\alpha|_{1}\right) \tag{7.18}
\end{equation*}
$$

if $u$ is the solution of (2.2), where $C_{1}=C_{1}\left(\|a\|_{1, \infty}\right)$.
Proof: The existence and uniqueness of $u_{h}^{*}$ are obvious. We also define $\tilde{u}_{h} \in W_{h}^{1}$ by

$$
\begin{equation*}
R_{h} \tilde{u}_{h}=R_{h} u \tag{7.19}
\end{equation*}
$$

by standard arguments (see, e.g., [6]),

$$
\begin{equation*}
\left\|u-\tilde{u}_{h}\right\| \leqslant C h^{2}\|u\|_{2} \tag{7.20}
\end{equation*}
$$

vol. 27, $\mathrm{n}^{\circ} 1,1993$

Observe that, by (7.17) and (7.19),

$$
R_{h}\left(u_{h}^{*}-\tilde{u}_{h}\right)=\bar{\lambda}_{h}-R_{h} u
$$

Then, by a simple scaling argument (see, e.g., [5]),

$$
\begin{equation*}
\left\|u_{h}^{*}-\tilde{u}_{h}\right\|_{0, T} \leqslant C h_{T}^{1 / 2} \sum_{i=1}^{3}\left\|\bar{\lambda}_{h}-R_{h} u\right\|_{e_{i}} \tag{7.21}
\end{equation*}
$$

for all $T \in T_{h}$ with edges $e_{i}, i=1,2,3$. Combine (7.15), (7.20) and (7.21) with (7.2), (7.5), (7.7), and (7.10) to obtain the desired result (7.18), and the proof is complete.

Note that $u_{h}^{*}$ approximates $u$ with a higher order of accuracy than $\bar{u}_{h}$, as required, and is continuous at the midpoints of sides in $E_{h}^{0}$ and zero at the midpoints of sides in $E_{h}^{\partial}$.

### 7.3. Implementation

Let $M_{h}$ be defined as in $\S 3$ or in $\S 4$. Corresponding to (2.6), we define the analogue :

Find $\bar{\psi}_{h} \in M_{h}$ such that

$$
\begin{equation*}
\sum_{T}\left(\alpha_{h}^{-1} P_{V}\left(\nabla \bar{\psi}_{h}\right), \nabla v\right)_{T}=\left(P_{h} f, v\right), \quad \forall v \in M_{h} \tag{7.22}
\end{equation*}
$$

Then, in the same argument as in § 2 , we have
THEOREM 7.4: Let $\left(\bar{\sigma}_{h}, \bar{u}_{h}, \bar{\lambda}_{h}\right)$ be the solution of (7.14) and let $\bar{\psi}_{h} \in M_{h}$ be given by

$$
\begin{equation*}
P_{h} \bar{\psi}_{h}=\bar{u}_{h}, \quad \tilde{R}_{h} \bar{\psi}_{h}=\bar{\lambda}_{h} . \tag{7.23}
\end{equation*}
$$

Then $\bar{\psi}_{h}$ is the unique solution of (7.22) and

$$
\begin{equation*}
\bar{\sigma}_{h}=-\alpha_{h}^{-1} P_{V}\left(\nabla \bar{\psi}_{h}\right) . \tag{7.24}
\end{equation*}
$$

Let us now discuss the structure of (7.22). Let $N_{h}=N_{2 h}$ or $N_{3 h}$. For $v \in M_{N C}$, we have $P_{V}(\nabla v)=\nabla v$, a piecewise constant. Moreover, the gradient of a bubble function in $N_{h}$ has zero mean value on each $T$. Indeed, for $v \in N_{h}$ and $q=(1,0)$ or ( 0,1 ), we have

$$
(\nabla v, q)_{T}=\left(v, q \cdot n_{T}\right)_{\partial T}-(v, \operatorname{div} q)_{T}=0
$$

since $v$ vanishes at the two Gaussian quadrature points of or on each side of $T$. Therefore, the solution of (7.22) may be determined as $z_{h}+\xi_{h}$ where

[^2]$\left(z_{h}, \xi_{h}\right) \in M_{N C} \times N_{h}$ is the unique solution of
\[

$$
\begin{gather*}
\sum_{T}\left(\alpha_{h}^{-1} \nabla z_{h}, \nabla v\right)_{T}=\left(P_{h} f, v\right), \quad \forall v \in M_{N C},  \tag{7.25a}\\
\sum_{T}\left(\alpha_{h}^{-1} P_{V}\left(\nabla \xi_{h}\right), \nabla \varphi\right)_{T}=\left(P_{h} f, \varphi\right), \quad \forall \varphi \in N_{h} . \tag{7.25b}
\end{gather*}
$$
\]

In summary, we have

$$
\begin{align*}
& \bar{\sigma}_{h}=-\alpha_{h}^{-1}\left(\nabla z_{h}+P_{V}\left(\nabla \xi_{h}\right)\right),  \tag{7.26}\\
& \bar{u}_{h}=P_{h}\left(z_{h}+\xi_{h}\right),  \tag{7.27}\\
& \bar{\lambda}_{h}=R_{h} z_{h}, \tag{7.28}
\end{align*}
$$

where $\left(z_{h}, \xi_{h}\right)$ satisfies (7.25). From the solution of (7.25) one can deduce the solution of (7.14). Moreover, based on (7.25)-(7.28), we have the simple solution given in the following proposition by the argument in [11].

PROPOSITION 7.5 : In each $T, \sigma_{h}$ at a point $x$ is evaluated by the simple formula

$$
\bar{\sigma}_{h}=-\alpha_{h}^{-1} \nabla z_{h}+\left(P_{h} f\right)_{T}\left(x-x_{T}\right) / 2, \quad x \in T,
$$

where $z_{h}$ satisfies ( $7.25 a$ ).
We shall now derive error estimates for (7.22).
THEOREM 7.6: If $u$ and $\bar{\psi}_{h}$ are the solutions of (2.2) and (7.22), respectively, and if

$$
\begin{equation*}
P_{V}\left(\nabla \bar{\psi}_{h}\right)=\nabla \bar{\psi}_{h}, \tag{7.29}
\end{equation*}
$$

then,

$$
\begin{align*}
& \left\|\nabla\left(u-\bar{\psi}_{h}\right)\right\|_{h}=\left(\sum_{T}\left\|\nabla u-\nabla \bar{\psi}_{h}\right\|_{T}^{2}\right)^{1 / 2} \leqslant C_{2} h\left(\|u\|_{2}+\|\alpha\|_{1}\right)  \tag{7.30}\\
& \left\|v-\bar{\psi}_{\eta}\right\| \leqslant C_{3} h^{2}\left(\|a\|_{1}+\|f\|_{1}\right) \tag{7.31}
\end{align*}
$$

where $C_{2}=C_{2}\left(\|u\|_{1, \infty}\right)$ and $C_{3}=C_{3}\left(\|u\|_{2, \infty},\|a\|_{1, \infty}\right)$.
Proof : First, by (2.1) and (7.10), note that

$$
\begin{align*}
& \left\|a-\alpha_{h}^{-1}\right\|_{-s} \leqslant C\|\alpha\|_{1} h^{s+1}, \quad s=0,1,  \tag{7.32}\\
& \left\|a-\alpha_{h}^{-1}\right\|_{\infty} \leqslant C\|\alpha\|_{1, \infty} h . \tag{7.33}
\end{align*}
$$

vol. $27, \mathrm{n}^{\circ} 1,1993$

Then, using (2.1), (7.32) with $s=0$, (7.2), (7.7), and the triangle inequality,

$$
\begin{aligned}
\left\|\nabla\left(u-\bar{\psi}_{h}\right)\right\|_{h} & \leqslant C\left\|\alpha_{h}^{-1} \nabla\left(u-\bar{\psi}_{h}\right)\right\|_{h} \\
& \leqslant C\left(\left\|\left(a-\alpha_{h}^{-1}\right) \nabla u\right\|+\left\|\sigma-\bar{\sigma}_{h}\right\|\right) \\
& \leqslant C 2 h\left(\|u\|_{2}+\|\alpha\|_{1}\right)
\end{aligned}
$$

i.e., (7.30) holds with $C_{2}=C_{2}\left(\|u\|_{1, \infty}\right)$.

In order to prove (7.31), we shall adapt the duality argument given in § 2. Let $w=u-\bar{\psi}_{h}$, and let $\phi \in H_{0}^{1}(\Omega)$ be such that

$$
-\operatorname{div}(a \nabla \phi)=w \quad \text { in } \Omega,
$$

and

$$
\begin{equation*}
\|\phi\|_{2} \leqslant C\|w\| . \tag{7.34}
\end{equation*}
$$

As in (2.15), we write

$$
\begin{align*}
\|w\|^{2}= & \sum_{T}(a \nabla \phi, \nabla w)_{T}-\sum_{T}\left(a \nabla \phi \cdot n_{T}, w\right)_{\partial T} \\
= & \sum_{T}\left(\left\{a-\alpha_{h}^{-1}\right\} \nabla \phi, \nabla w\right)_{T}+\sum_{T}\left(\alpha_{h}^{-1} \nabla \phi, \nabla w\right)_{T}  \tag{7.35}\\
& -\sum_{T}\left(a \nabla \phi \cdot n_{T} w\right)_{\partial T} \\
\equiv & R_{1}+R_{2}+R_{3}
\end{align*}
$$

Using (7.33), we see that

$$
\begin{align*}
\left|R_{1}\right| & \leqslant\left\|a-\alpha_{h}^{-1}\right\|_{\infty}\|\nabla \phi\|\|\nabla w\|_{h}  \tag{7.36}\\
& \leqslant C_{1} h^{2}\|\nabla w\|_{h}^{2}+\varepsilon\|\phi\|_{2}^{2}
\end{align*}
$$

where $C_{1}=C_{1}\left(\|a\|_{1, \infty}\right)$. The term $R_{3}$ can be treated in the same manner as in the second section to obtain

$$
\begin{equation*}
\left|R_{3}\right| \leqslant C h\|\phi\|_{2}\|\nabla w\|_{h} . \tag{7.37}
\end{equation*}
$$

For $R_{2}$, observe that

$$
\begin{align*}
R_{2} & =\sum_{T}\left(\alpha_{h}^{-1} \nabla \phi, \nabla\left(u-\bar{\psi}_{h}\right)\right)_{T} \\
& =\sum_{T}\left(\left\{\alpha_{h}^{-1}-a\right\} \nabla \phi, \nabla u\right)_{T}+\sum_{T}\left\{(a \nabla \phi, \nabla u)_{T}-\left(\alpha_{h}^{-1} \nabla \phi, \nabla \bar{\psi}_{h}\right)_{T}\right\} \\
& \equiv R_{2}^{1}+R_{2}^{2} . \tag{7.38}
\end{align*}
$$

Applying the same ideas as in (2.16)-(2.17), we get

$$
\begin{equation*}
\left|R_{2}^{2}\right| \leqslant C h^{2}\|f\|_{1} \tag{7.39}
\end{equation*}
$$

Finally, $R_{2}^{1}$ can be bounded as follows :

$$
\begin{align*}
\left|R_{2}^{1}\right| & =\left|\sum_{T}\left(\left\{\alpha_{h}^{-1}-a\right\} \nabla \phi, \nabla u\right)_{T}\right| \\
& \leqslant\left\|\alpha_{h}^{-1}-a\right\|_{-1}\|\nabla u \cdot \nabla \phi\|_{1}  \tag{7.40}\\
& \leqslant C_{3}\left\|\alpha_{h}^{-1}-a\right\|_{-1}^{2}+\varepsilon\|\phi\|_{2}^{2}
\end{align*}
$$

where $C_{3}=C_{3}\left(\|u\|_{2, \infty}\right)$. Now, combine (7.30), (7.32), and (7.34)-(7.40) to yield the desired result (7.31) if $h$ is sufficiently small, and the proof has been finished.

## ACKNOWLEDGMENTS

This paper contains most of the results of the author's doctoral dissertation under the supervision of Professor Jim Douglas, Jr. at Purdue University [4]. The author wish to thank him for many helpful suggestions in this paper. Also, the author would like to thank Professors Franco Brezzi and Paola Pietra for several useful conversations.

## REFERENCES

[1] T. Arbogast, A new formulation of mixed finite element methods for second order elliptic problems (to appear).
[2] D. N. Arnold and F. Brezzi, Mixed and nonconforming finite element methods : implementation postprocessing and error estimates, RAIRO Modél. Math. Anal. Numér., 19 (1985), pp. 7-32.
[3] F. Brezzi, J. Douglas Jr and L. Donatella Marini, Two families of mixed finite elements for second order elliptic problems, Numer. Math., 47 (1985), pp. 217-235.
[4] Z. Chen, On the relationship between mixed and Galerkin finite element methods, Ph. D. thesis, Purdue University, West Lafayette, Indiana, August (1991).
[5] F. Brezzi and M. Fortin, Hybrid and Mixed Finite Element Methods, to appear.
[6] P. Ciarlet, The Finite Element Method for Elliptic Problems, North-Holland, Amsterdam, 1978.
[7] J. Douglas Jr and J. E. Roberts, Global estimates for mixed methods for second order elliptic problems, Math. Comp., 45 (1985), pp. 39-52.
[8] R. Falk and J. Osborn, Error estimates for mixed methods, RAIRO, Modél. Math. Anal. Numér., 14 (1980), pp. 249-277.
[9] M. Fortin and M. Soulie, A non-conforming piecewise quadratic finite element on triangles, Internat. J. Numer. Methods Engrg., 19 (1983), pp. 505520.
[10] B. X. Fraeijs de Veubeke, Displacement and equilibrium models in the finite element method, in Stress Analysis, O. C. Zienkiewicz and G. Holister (eds.), John Wiley, New York, 1965.
[11] L. Donatella Marini, An inexpensive method for the evaluation of the solution of the lowest order Raviart-Thomas mixed method, SIAM J. Numer. Anal., 22 (1985), pp. 493-496.
[12] L. Donatella Marini and P. Pietra, An abstract theory for mixed approximations of second order elliptic problems, Mat. Apl. Comput., 8 (1989), pp. 219-239.
[13] P. A. Raviart and J. M. Thomas, A mixed finite element method for second order elliptic problems, in Mathematical Aspects of the Finite Element Method, Lecture Notes in Math. 606, Springer-Verlag, Berlin and New York (1977), pp. 292-315.


[^0]:    (*) Manuscrit received September 1991.
    AMS(MOS) subject classifications (1985 revision). Primary 65N15, 65N30; Secondary $41 \mathrm{~A} 10,41 \mathrm{~A} 25$.
    (1) AHPCRC, University of Minnesota, 1100 Washington Ave. South, Minneapolis, MN 55415, U.S.A.

[^1]:    $\mathrm{M}^{2}$ AN Modélisation mathématique et Analyse numérique 0764-583X/93/01/9/26/\$4.60 Mathematical Modelling and Numerical Analysis (C) AFCET Gauthier-Villars

[^2]:    $\mathbf{M}^{2}$ AN Modélisation mathématique et Analyse numérique
    Mathematical Modelling and Numerical Analysis

