Ann. I. H. Poincaré – AN 20, 5 (2003) 867–888 © 2003 Éditions scientifiques et médicales Elsevier SAS. All rights reserved 10.1016/S0294-1449(03)00015-5/FLA

NONHYPERBOLIC PERSISTENT ATTRACTORS NEAR THE MORSE–SMALE BOUNDARY

Eduardo M. MUÑOZ MORALES *,1, Bernardo SAN MARTÍN REBOLLEDO ², Jaime A. VERA VALENZUELA ²

Departamento de Matemáticas, Universidad Católica del Norte, Casilla 1280, Antofagasta, Chile Received 5 May 2002, accepted 30 September 2002

Dedicated to Professor Jacob Palis on his sixtieth birthday

ABSTRACT. – In this paper we present a submanifold \mathcal{N} of codimension two contained in the Morse–Smale regular boundary in $\mathcal{X}^r(M)$, the space of \mathcal{C}^r vector fields on a 3-dimensional manifold M, which exhibits singular cycles, such that for generic 2-parameter families going through, it has nonhyperbolic nontrivial attractors for a set of parameters with positive Lebesgue measure.

© 2003 Éditions scientifiques et médicales Elsevier SAS

RÉSUMÉ. – Dans cet article on construit une sous-variété \mathcal{N} de codimension deux contenue dans la frontière régulière de l'ensemble des champs de vecteurs Morse–Smale de $\mathcal{X}^r(M)$, l'espace de champs de vecteurs de classe C^r sur la variété de dimension trois M, laquelle présente des cycles singuliers tels que, toute famille générique à deux paramètres transversale à \mathcal{N} possède des ensembles attractifs non hyperboliques et non triviaux pour un ensemble de paramètres dont la mesure de Lebesgue est positive.

© 2003 Éditions scientifiques et médicales Elsevier SAS

1. Introduction

Morse–Smale systems constitute an open set of dynamically or structurally stable elements as classically established in Palis [12] and Palis–Smale [13].

On the other hand, singular cycles constitute a bifurcation mechanism for vector fields that, to the opposite to what happen with homoclinic tangencies in the diffeomorphisms

^{*} Corresponding author.

E-mail addresses: emunoz@ucn.cl (E.M. Muñoz Morales), sanmarti@ucn.cl (B. San Martín Rebolledo), jvera@ucn.cl (J.A. Vera Valenzuela).

¹ Partially supported by Fondecyt Grant 1981241, IMPA CNPq and UCN.

² Partially supported by Fondecyt Grant 1981241 and 1000047, IMPA CNPq, ICTP, Trieste, Italia and UCN.

case, the hyperbolicity from the parameter point of view, is a total Lebesgue measure phenomenom for generic unfolding of these type of vector fields, see [3,11]. These kind of cycles were studied for the first time in [7], as being part of the so called singular horseshoe, an example (on manifolds with boundary) of a structurally stable vector field with a nonhyperbolic nonwandering set.

An important ingredient to obtain these results about hyperbolicity, is the presence in the cycle of just one branch of the unstable manifold of the singularity. As we will observe later, the consideration of both branches on the cycle, plus a contractivity hyphotesis for the cycle itself it will imply that the generic unfolding exhibits nonhyperbolic nontrivial attractors which are persistent in the measure sense.

Before to give the precise statement of our result we establish the setting and concepts that we will need. Let M be a compact boundaryless 3-dimensional manifold and let $\mathcal{X}^r(M)$ be the Banach space of C^r vector fields on M. For $X \in \mathcal{X}^r(M)$, $\Gamma(X)$ denotes its chain recurrent set. We say that $X \in \mathcal{X}^r(M)$ is *simple* if $\Gamma(X)$ is a union of finitely many hyperbolic critical orbits. By a *critical orbit* we mean an orbit that is either periodic or singular. It is easy to see that the set S^r of simple C^r vector fields is an open subset of $\mathcal{X}^r(M)$.

An orbit γ of a vector field X is *transversal* if $\alpha(\gamma)$ and $\omega(\gamma)$ are critical hyperbolic elements and their invariant manifolds meet transversally along γ , here α and ω stand for the α and ω -limit sets, respectively. We say that X is a *Morse–Smale* system if X is simple and every orbit is transversal.

A cycle Λ of $X \in \mathcal{X}^r(M)$ is a compact invariant chain recurrent set consisting of finitely many orbits whose α and ω -limit sets are critical elements linked in a cyclical way. If the cycle contains at least one singularity, we will call it a *singular cycle*. If σ_i and σ_j are critical elements on the cycle, we write $\sigma_i \prec \sigma_j$ to mean that there is an orbit γ_{ij} on the cycle such that $\alpha(\gamma_{ij}) = \sigma_i$ and $\omega(\gamma_{ij}) = \sigma_j$.

In this paper we are concerned with those cycles Λ whose critical elements satisfy $\sigma_0 \prec \sigma_1 \prec \cdots \prec \sigma_k = \sigma_0$, where σ_0 is the unique singularity in the cycle, dim $W^u(\sigma_0) = 1$ and $W^u(\sigma_0) \subset \Lambda \cap W^s(\sigma_1)$. Furthermore, if we denote by γ_l and γ_r the branches of $W^u(\sigma_0)$, then $\Lambda_r = \Lambda \setminus \{\gamma_l\}$ and $\Lambda_l = \Lambda \setminus \{\gamma_r\}$ are contractive singular cycles as in [3] satisfying some extra condition. In fact, as in [3] Λ_r and Λ_l satisfy:

- (1) The eigenvalues $-\lambda_1, \lambda_2$ and $-\lambda_3$ of the singularity σ_0 are real and satisfy $-\lambda_3 < -\lambda_1 < 0 < \lambda_2$ and $\lambda_1 \lambda_2 > 0$, i.e., the cycle is contractive. Besides, we assume the additional condition $\beta > \alpha + 1$ where $\alpha = \lambda_1/\lambda_2$ and $\beta = \lambda_3/\lambda_2$.
- (2) For all $p \in \gamma_l \cup \gamma_r$ and any invariant manifold W^{cu} of X passing through σ_0 tangent to the eigenspace associated with the eigenvalues $-\lambda_1$ and λ_2 holds:

$$T_p(W^{cu}) + T_p(W^s(\omega(p))) = T_p M.$$

(3) There is a neighbourhood \mathcal{U} of X such that if $Y \in \mathcal{U}$ the continuations $\sigma_i(Y), 0 \leq i \leq k-1$, of critical orbits $\sigma_i(X)$ of the cycle are well defined, the vector field Y is C^2 -linearizing nearby $\sigma_0(Y)$ and the Poincaré maps of $\sigma_i(Y), 1 \leq i \leq k-1$, are C^2 -linearizing ones.

In linearizing coordinates for σ_0 the flow is given by $X^t(x_3, x_2, x_1) = (e^{-\lambda_3 t} x_3, e^{\lambda_2 t} x_2, e^{-\lambda_1 t} x_1)$ and then $W^s_{loc}(\sigma_0) = \{(x_3, 0, x_1)\}$ and $W^u_{loc}(\sigma_0) = \{(0, x_2, 0)\}$. We define $\Sigma_0 = \{(x_3, x_2, 1)\}$.

Let Σ_1 be a transversal section to the flow through a point q of σ_1 , then in linearizing coordinates, $\{(x, 0)\}$ and $\{(0, y)\}$ are the local stable and unstable manifolds of the fixed point q = (0, 0) of the Poincaré map induced by σ_1 , respectively. The flow generated by X also induces a Poincaré map $\mathcal{P}_{01}: \mathcal{D}om(\mathcal{P}_{01}) \subset \Sigma_0 \to \Sigma_1$. $\mathcal{P}_{01}(w)$ is the first point of intersection among the positive orbit of w and the section Σ_1 .

- (4) The orbit $\gamma_{k-1,0}$ in the cycle is not contained in the strong stable manifold of σ_0 , and we suppose that intersects Σ_0 .
- (5) The branch γ_l (resp. γ_r) of $W^u(\sigma_0(X))$ is the only nontransversal orbit in Λ_l (resp. Λ_r). Moreover, if $\sigma \subset \Lambda = \Lambda_l \cup \Lambda_r$ is a periodic orbit, then there is only one regular orbit γ in the cycle Λ such that $\sigma = \alpha(\gamma)$.
- (6) For *j* = *l*, *r*, Λ_j is *isolated*, i.e., there exists a neighbourhood U_j of Λ_j, which is called an *isolating block* such that ∩_t X^t(U_j) does not contains any orbit close to W^u(σ₀(X)). Here X^t : M ⇔ is the flow generated by X. U = U_l ∪ U_r is an isolating block for Λ.

In Fig. 1 a sketch of the geometrical shape of our cycle considering the coordinates system as mentioned in (3), is shown.

Let $C^k(I^2, \mathcal{X}^r(M))$, $k \ge 2$, be the space of families $\{X_{\mu,\nu}\}$ of C^r -vector fields, $r \ge 2$, depending on two parameters $(\mu, \nu) \in I^2 = [-1, 1] \times [-1, 1]$ such that the map $(\mu, \nu) \in I^2 \to X_{\mu,\nu} \in \mathcal{X}^r(M)$ is C^k . Let us endow $C^k(I^2, \mathcal{X}^r(M))$ with the C^k -topology. We are interested in families $\{X_{\mu,\nu}\} \in C^k(I^2, \mathcal{X}^r(M))$ such that (after a reparametrization, if necessary) $X_{0,0}$ has a singular cycle as described above.

We say that a compact invariant transitive set \underline{Q} of a vector field X is an *attractor* if there is a neighbourhood $U \supseteq Q$ such that $\overline{X^t(U)} \subseteq U$ for all t > 0 and $Q = \bigcap_{t \ge 0} X^t(U)$. If Q is a transitive set but not a regular orbit, we say that it is a *nontrivial attractor*. Now, we state our main result:



THEOREM 1. – For generic families $\{X_{\mu,\nu}\} \in C^k(I^2, \mathcal{X}^r(M))$ having a singular cycle for $X_{(\mu_0,\nu_0)} = X_{(0,0)}$ that satisfies the conditions (1)–(6) above, it holds that

 $m(\{(\mu, \nu): X_{\mu,\nu} \text{ has a nonhyperbolic nontrivial attractor}\}) > 0$

where m stands for the Lebesgue measure in I^2 . Moreover, this set accumulates (0,0).

Remark 1.-

- (1) The condition $\beta > \alpha + 1$ in (1) above had been established in order to obtain certain appropriate C^2 -invariant foliation (Lemma 1, Section 2), leading us to a one-dimensional dynamics reduction on our study. Besides, in Lemma 3 we need this class of differentiability in order to apply Mañé's Theorem [8].
- (2) Condition (2) above also means that $W^{cu}(\sigma_0)$ meets $W^s(\sigma_1)$ transversally.
- (3) The conditions (5) and (6) are required in order to guarantee that we can construct the vector field X in the boundary of the Morse–Smale systems. In fact, by condition (5) we can slightly perturb the system X "moving down" the first intersection point belonging to each of both branches of W^u(σ₀) with Σ₁ without create homoclinic points (this can not happen if we "move this point up"), and then the condition (6) (with some extra hypothesis about the existence of a filtration, see [18] for definitions and details) guarantee us that there are not changes (Ω-explosion) in the dynamics far away from the cycle. In this way we can attain the dynamics of X (the cycle) from the Morse–Smale systems and therefore this kind of cycles constitute a bifurcation mechanism for the Morse–Smale systems.
- (4) Condition (6) implies that the eigenvalues associated to the σ_1 are positives, otherwise a singular horseshoe appears [7].
- (5) Condition (5) contains the only two degeneracencies, so the cycle is a codimension two phenomena.
- (6) The attractors obtained are one side orientation preserving Lorenz-like atractors.
- (7) A similar result can be obtained when we consider cycles in which the branches of the singularity σ_0 are linked to distinct periodic orbits. This result is established in a forthcoming paper [17].
- (8) The case where the cycle reduces to a double homoclinic cycle has been studied in [9].

The paper is organized as follows. To understand the dynamics of $X_{\mu,\nu}$, $\mu, \nu > 0$, we restrict our analysis to the dynamics of $X_{\mu,\nu}$ on $\Gamma(X_{\mu,\nu}, U)$ where U is an isolating block. In Section 2 we establish the existence of a C^2 -invariant stable foliation (Lemma 1). This reduces our study to the dynamics of one-dimensional maps $f_{\mu,\nu}$, and some limit dynamic is obtained via renormalization (Lemma 2). In Section 3 we show that for a large enough positive integer *n*, there exists a positive Lebesgue measure set E_n in the (μ, ν) -parameter space such that the critical values of $g_{\mu,\nu} = (f_{\mu,\nu})^{n+1}$ have positive Lyapunov exponents and the clousure of the $g_{\mu,\nu}$ -iterates of these points is an interval (Theorem 2). From this result our Theorem 1 follows. Theorem 2 is proved by applying Benedicks–Carleson techniques, and in order to do that, a fundamental and crucial fact is to prove that the maximal orbits outside of a certain neighbourhood of

the critical points have exponential growth, provided that the number of elements in the orbit is big enough.

2. Renormalization

Our first proposition, Proposition 1, points out that systems having our cycles constitute (locally) submanifolds of codimension two in $\mathcal{X}^r(M)$ contained in the boundary of the Morse–Smale systems and, consequently, the natural approach to study their bifurcations is by mean of two parameter families transversal to these codimension two submanifolds.

PROPOSITION 1. – Let Λ be a singular cycle for the vector field X as described in the introduction, and let U be an isolating block associated to it. Then, there exists a neighbourhood U of X in $\mathcal{X}^r(M)$ and a C^r -submanifold \mathcal{N} of codimension two in $\mathcal{X}^r(M)$ containing X such that if $Y \in \mathcal{N} \cap \mathcal{U}$, then $\Gamma(Y, U) = \bigcap_t Y^t(U)$ contains a cycle topologically equivalent to Λ .

Proof. – The proof is quite simple and it goes as follows: let γ_r and γ_l be the branches of $W^u(\sigma_0)$. Let Σ_1 be a transversal section to X at $q \in \sigma_1$. For each Y close to X we consider linearizing coordinates (x, y) on Σ_1 for the Poincaré map associated to σ_1 (which vary smoothly with Y) such that $W^s_{loc}(\sigma_1) \cap \Sigma_1 = \{(x, 0)\}$. If \mathcal{U} is a small enough neighbourhood of X, then we can define the C^r -map $G_i : \mathcal{U} \to \mathbb{R}$ by $G_i(Y) = \pi_2 \circ p_i(Y)$, where $p_i(Y)$ is the first intersection of γ_i with Σ_1 , i = l, r, and π_2 denotes the projection over the second term on the coordinates of Σ_1 . It is clear that $B_r = G_r^{-1}(0)$ is a submanifold of codimension one of \mathcal{U} that contains X. Analogously we obtain $B_l = G_l^{-1}(0)$ which is transversal to B_r , and thus we get that both, $\mathcal{N} = B_r \cap B_l$ and \mathcal{U} , satisfy the conditions of the proposition. In fact, if $G : \mathcal{U} \to \mathbb{R}^2$ is defined by $G(Y) = (G_r(Y), G_l(Y))$ then G is a submersion and $\mathcal{N} = G^{-1}(0, 0)$. \Box

Let X be a vector field having a cycle Λ as above, and let S be a two-dimensional C^k -submanifold of $\mathcal{X}^r(M)$ transversal to \mathcal{N} at X. The map $G: S \to \mathbb{R}^2$ defined by $G(Y) = (G_r(Y), G_l(Y))$ is a C^k -diffeomorphism in a neighbourhood of X in S and then, it determines a parametrization of S in a neighbourhood of X, i.e., $Y = X_{\mu,\nu}$ if and only if $G(Y) = (\mu, \nu)$. Moreover, $X = X_{0,0}$ and $X_{\mu,\nu}$ is a Morse–Smale system if and only if μ and ν are both negatives. Furthermore, from [3,6,11,16] we know that in the Lebesgue sense, for almost every (μ, ν) close enough to (0, 0) and such that $\mu\nu < 0$ we have that $X_{\mu,\nu}$ is an Axiom A system. Therefore, our work will be focus on the region corresponding to positive values for the parameters μ and ν .

Let $\{X_{\mu,\nu}\} \in C^k(I^2, \mathcal{X}^r(M))$, and let Λ be a singular cycle associated to the vector field $X = X_{0,0}$. Using transversality arguments we can change the cycle for other one, choosing a new cycle contained in $\Gamma(X, U)$, still named Λ , having only one periodic orbit, i.e., $\Lambda = \{\sigma_0, \gamma_r, \gamma_l, \sigma_1\}$ linked by $\sigma_0 \prec \sigma_1 \prec \sigma_2 = \sigma_0$.

Let Σ_1 be a cross section to the flow of X at $q \in \sigma_1(X)$ and choose C^2 -linearizing coordinates (x, y) for the Poincaré map induced by $\sigma_1(X_{\mu,\nu})$, depending smoothly on the parameters. In this coordinates we have:

- (1) $Q = \{(x, y)/|x|, |y| \leq 2\} \subset \Sigma_1.$
- (2) $\{(x, 0)/|x| \leq 1\} \subset W^s(\sigma_1(X_{\mu,\nu})).$
- (3) $\{(0, y)/|y| \leq 1\} \subset W^u(\sigma_1(X_{\mu, \nu})).$
- (4) $\Gamma(X_{\mu,\nu}, U) \cap \Sigma_1 \subset \{(x, y)/y \ge 0\}.$
- (5) The Poincaré map $P_{\mu,\nu}$, defined by $\sigma_1(X_{\mu,\nu})$ is linear and $P_{\mu,\nu}(x, y) = (\tau_{\mu,\nu}x, \rho_{\mu,\nu}y)$ where $0 < \tau_{\mu,\nu} < 1$ and $\rho_{\mu,\nu} > 1$ (see (5) and (6) in Section 1).

A closed subset $R \subset \Sigma_1$ is called a *horizontal strip* if it is bounded (in Q) by two disjoint continuous curves that connect the vertical sides of Q.

There are two horizontal strips R_0 and R_1 in Q such that for (μ, ν) close to (0, 0), the $X_{\mu,\nu}$ positive orbit of any point in $R_0 \cup R_1$ meets Σ_1 and then, there is a (μ, ν) -dependent first return map $F_{\mu,\nu}: R_0 \cup R_1 \to \Sigma_1$. In fact, R_0 is a horizontal strip containing the $P_{\mu,\nu}$ -local stable manifold of $q_{\mu,\nu}$, and $F_{\mu,\nu}$ restricted to this strip coincides with the Poincaré map $P_{\mu,\nu}$. R_1 is a horizontal strip containing in its interior the connected component $W_{0,0}^s$ of $W^s(\sigma_0(X_{0,0})) \cap Q$ which contains the last intersection point among $\gamma_{1,0}$ (time oriented) and Q. Here we note that, as in [3] the strip R_1 , there denoted by R_Y^0 , depends on each vector field $Y = X_{\mu,\nu}$ close to X. In [3] one of the horizontal sides of R_Y^0 is contained in the continuous variation $W_{\mu,\nu}^s \cap W_{0,0}^s$ with respect to the vector field $X_{\mu,\nu}$. In that situation, the return map is defined only at one side of $W_{\mu,\nu}^s$, whereas in our case, the return map is defined at both sides of $W_{\mu,\nu}^s \subset R_1$ (this fact introduces a discontinuity in the Poincaré return map). In fact, the positive orbit of a point in R_1 has a segment close to one of the branches γ_l or γ_r of $W^u(\sigma_0(X_{\mu,\nu}))$ and then intersects Σ_1 , as it is shown in Fig. 2.

For simplicity, we observe that after multiplication by (μ, ν) -dependent constants in the coordinates we can suppose $W_{\mu,\nu}^s \cap \{(0, y)\} = \{(0, 1)\}.$

In order to find some analytical expression for $F_{\mu,\nu}$ we choose $(x_3, x_2, x_1) C^2$ linearizing coordinates for the singularity σ_0 depending smoothly on the vector field. In these coordinates the $X_{\mu,\nu}$ flow is given by $X_{\mu,\nu}^t(x_3, x_2, x_1) = (e^{-\lambda_3 t} x_3, e^{\lambda_2 t} x_2, e^{-\lambda_1 t} x_1)$ where $-\lambda_3 < -\lambda_1 < 0 < \lambda_2$ are the (μ, ν) -dependent eingenvalues of $DX_{\mu,\nu}(\sigma_0(X_{\mu,\nu}))$.



Fig. 2.

Next, we take the sections $\Sigma_0 = \{(x_3, x_2, 1)\}$, $\Sigma_l = \{(x_3, -1, x_1)\}$ and $\Sigma_r = \{(x_3, 1, x_1)\}$ (see Fig. 2), and consider the (μ, ν) -dependent Poincaré map $P_1 : \Sigma_0 \setminus \{(x_3, 0, 1)\} \rightarrow \Sigma_l \cup \Sigma_r$ induced by the $X_{\mu,\nu}$ flow given by $P_1(x_3, x_2, x_1) = (x_3|x_2|^{\beta_{\mu,\nu}}, |x_2|/x_2, |x_2|^{\alpha_{\mu,\nu}})$, where $\alpha_{\mu,\nu} = \lambda_1/\lambda_2$ and $\beta_{\mu,\nu} = \lambda_3/\lambda_2$.

Finally, we consider the (μ, ν) -dependent diffeomorphisms $P_0: \mathcal{D}om(P_0) \subset \Sigma_1 \to \Sigma_0$ and $P_2: \mathcal{D}om(P_2) \subset \Sigma_l \cup \Sigma_r \to \Sigma_1$ induced by the $X_{\mu,\nu}$ flow. Then, for $F_{\mu,\nu}$ restricted to R_1 we have

$$F_{\mu,\nu}(x, y) = (P_2 \circ P_1 \circ P_0)(x, y) = (u_{\mu,\nu}(x, y), v_{\mu,\nu}(x, y))$$

By using the chain rule it is inmediate that there is a positive constant C independent from the parameter (μ, ν) such that

$$C^{-1}|x_{2}|^{\alpha_{\mu,\nu}-1} \leq \left|\partial_{y}v_{\mu,\nu}(x,y)\right| \leq C|x_{2}|^{\alpha_{\mu,\nu}-1},$$

$$C^{-1}|x_{2}|^{\alpha_{\mu,\nu}-2} \leq \left|\partial_{y}^{2}v_{\mu,\nu}(x,y)\right| \leq C|x_{2}|^{\alpha_{\mu,\nu}-2}.$$
(1)

The following lemma, whose proof is obtained in the same way than in Lemma 1 of [3] (see also p. 695 in [7]), implies that after a C^2 -change of coordinates the second component of $F_{\mu,\nu}$ does not depend on x, i.e., $v_{\mu,\nu}(x, y) = f_{\mu,\nu}(y)$ and therefore, our study is reduced to a setting in one-dimensional dynamics.

LEMMA 1. – If $\beta > \alpha + 1$ then for each pair (μ, ν) small enough there exists a C^2 invariant stable foliation $\mathcal{F}^s_{\mu,\nu}$ defined on Σ_1 with a C^k -dependence on the parameter (μ, ν) .

Using this lemma, the formula for P_1 and Taylor expansion for P_0 and P_2 , we have that after a C^2 -change of coordinates given by proyection along the leaves of $\mathcal{F}^s_{\mu,\nu}$, and depending C^k on the parameter (μ, ν) , the return map $F_{\mu,\nu}$ can be expressed by

$$F_{\mu,\nu}(x,y) = \begin{cases} (\widetilde{\mu}_{\mu,\nu}(x,y), f_{\mu,\nu}(y)), & (x,y) \in R_1, \\ (\tau_{\mu,\nu}x, f_{\mu,\nu}(y)), & (x,y) \in R_0, \end{cases}$$

where $\{f_{\mu,\nu}\}$ is the C^k -two parameter family of one-dimensional C^2 -maps induced by the stable foliation $\mathcal{F}^s_{\mu,\nu}$ defined by

$$f_{\mu,\nu}(y) = \begin{cases} \rho_{\mu,\nu}y, & 0 \leq y \leq \rho_{\mu,\nu}^{-1}b(\mu,\nu), \\ \mu - k_{\mu,\nu}(y)(1-y)^{\alpha_{\mu,\nu}}, & a(\mu,\nu) \leq y < 1, \\ \nu - \tilde{k}_{\mu,\nu}(y)(y-1)^{\alpha_{\mu,\nu}}, & 1 < y \leq b(\mu,\nu). \end{cases}$$

The functions $(\mu, \nu, y) \to k_{\mu,\nu}(y)$ and $(\mu, \nu, y) \to \tilde{k}_{\mu,\nu}(y)$ are positives, bounded and bounded away from zero (see (1)) and besides, C^k on the parameter (μ, ν) and C^2 on $y \in [a(\mu, \nu), 1) \cup (1, b(\mu, \nu)]$. The extreme points $a = a(\mu, \nu)$ and $b = b(\mu, \nu)$ are defined by the relations

$$\mu - k_{\mu,\nu}(a)(1-a)^{\alpha_{\mu,\nu}} = 0$$
 and $\nu - \tilde{k}_{\mu,\nu}(b)(b-1)^{\alpha_{\mu,\nu}} = 0.$

Note that this relation is given to leave out those values of y such that $f_{\mu,\nu}(y) < 0$, because they are not related to $\Gamma(X_{\mu,\nu}, U)$, since their orbits come back to Σ_1 at most a finite number of times, see Fig. 3.



Fig. 3.

Furthermore, from (1) we conclude that there are positive constants K_1 and K_2 , independents from the parameter (μ, ν) such that

$$K_{1}|1-y|^{\alpha_{\mu,\nu}-1} < \left|\partial_{y}f_{\mu,\nu}(y)\right| < K_{2}|1-y|^{\alpha_{\mu,\nu}-1},$$

$$K_{1}|1-y|^{\alpha_{\mu,\nu}-2} < \left|\partial_{y}^{2}f_{\mu,\nu}(y)\right| < K_{2}|1-y|^{\alpha_{\mu,\nu}-2}.$$
(2)

For a positive integer *n*, consider $g_{n,\mu,\nu} = (f_{\mu,\nu}^{n+1})|_{[a,b]\setminus\{1\}}$ where $a = a(\mu, \nu)$ and $b = b(\mu, \nu)$. Note that $g_{n,\mu,\nu}$ may be not defined for every $y \in [a,b]\setminus\{1\}$. However, if we suppose $\mu \leq \rho_{\mu,\nu}^{-n}a$ and $\nu \leq \rho_{\mu,\nu}^{-n}b$ then a straightforward evaluation gives

$$g_{n,\mu,\nu}(y) = \begin{cases} \rho_{\mu,\nu}^{n} [\mu - k_{\mu,\nu}(y)(1-y)^{\alpha_{\mu,\nu}}], & a \leq y < 1, \\ \rho_{\mu,\nu}^{n} [\nu - \tilde{k}_{\mu,\nu}(y)(y-1)^{\alpha_{\mu,\nu}}], & 1 < y \leq b. \end{cases}$$

In Fig. 3(a), (b) and (c), we have been drawn the graphs of $f_{\mu,\nu}$; $f_{\mu,\nu}^{n+1}$ for $\rho_{\mu,\nu}^n \mu < 1$ and $\rho_{\mu,\nu}^n \nu < 1$; and the graph of $f_{\mu,\nu}^{n+1}$ for $1 < \rho_{\mu,\nu}^n \mu < b$ and $1 < \rho_{\mu,\nu}^n \nu < b$. In the next section our attention will be focus on the case sketched in Fig. 3(c).

The next lemma is similar to Theorem 1, p. 47 in [14], and it means that after a suitable *n*-dependent reparametrization there is a well defined limit dynamic for g_n when *n* converges to infinite. These reparametrizations will be done in neighbourhoods of some special parameters values μ'_n and ν'_n . In order to do that, we have the following:

CLAIM 1. – Given n big enough, there are unique parameters μ'_n and ν'_n such that $\mu'_n = \nu'_n$ and $\rho^n_{\mu'_n,\nu'_n}\mu'_n = 1$ hold.

Proof. – For μ such that $\rho_{\mu,\mu}^n \mu < 2$, let $w_n(\mu) = \rho_{\mu,\mu}^n \mu$. Taking derivative we obtain, for some constant,

$$\partial_{\mu}w_{n}(\mu) = \rho_{\mu,\mu}^{n} \left[n \rho_{\mu,\mu}^{-1} \partial_{\mu}(\rho_{\mu,\mu}) \mu + 1 \right] \ge C \rho_{\mu,\mu}^{n}$$

and then $w_n(\mu)$ is an increasing function of μ . Consequently, there is a unique $\mu = \mu'_n$ such that $\rho^n_{\mu'_n,\mu'_n}\mu'_n = 1$. \Box

LEMMA 2 (Renormalization). – For $\alpha = \alpha_{0,0} = \lambda_1/\lambda_2$ and $K = \tilde{k}_{0,0}(1)/k_{0,0}(1)$ consider the two-parameter family

$$h_{\bar{\mu},\bar{\nu}}(\bar{y}) = \begin{cases} -(-\bar{y})^{\alpha} + \bar{\mu}, & \bar{y} < 0, \\ -K(\bar{y})^{\alpha} + \bar{\nu}, & 0 < \bar{y}. \end{cases}$$

Define $F(y, \mu, \nu) = (f_{\mu,\nu}(y), \mu, \nu)$, and $\Phi(\bar{y}, \bar{\mu}, \bar{\nu}) = (h_{\bar{\mu},\bar{\nu}}(\bar{y}), \bar{\mu}, \bar{\nu})$. Then, for each positive integer *n*, there are *n*-dependent reparametrizations $\mu = \mu_n(\bar{\mu}, \bar{\nu})$, $\nu = \nu_n(\bar{\mu}, \bar{\nu})$ of the (μ, ν) -parameter, and $(\bar{\mu}, \bar{\nu})$ -dependent changing coordinate transformations $\psi_{n,\bar{\mu},\bar{\nu}}$ such that if we consider the functions

$$\phi_n(\bar{y},\bar{\mu},\bar{\nu}) = (\psi_{n,\bar{\mu},\bar{\nu}}(\bar{y}),\mu_n(\bar{\mu},\bar{\nu}),\nu_n(\bar{\mu},\bar{\nu})),$$

$$\Phi_n = (\phi_n)^{-1} \circ F^{n+1} \circ \phi_n,$$

then $\Phi_n(\bar{y}, \bar{\mu}, \bar{\nu})$ converges C^2 to $\Phi(\bar{y}, \bar{\mu}, \bar{\nu})$ on any compact set of $\mathbb{R}\setminus\{0\} \times \mathbb{R}^2$. Moreover, the ϕ_n -image of a compact set $K \subset \mathbb{R}^3$, converges, for $n \to \infty$, in the (y, μ, ν) space to (1, 0, 0). Furthermore, if $h_{n,\bar{\mu},\bar{\nu}} = (\psi_{n,\bar{\mu},\bar{\nu}})^{-1} \circ g_{n,\mu,\nu} \circ \psi_{n,\bar{\mu},\bar{\nu}}$ then there are positive constants K_1 and K_2 independent of either $n \in \mathbb{N}$ and the $\bar{\mu}, \bar{\nu}$ parameters, such that

$$K_{1}|\bar{y}|^{\alpha_{\mu,\nu-1}} < \left|\partial_{\bar{y}}h_{n,\bar{\mu},\bar{\nu}}(\bar{y})\right| < K_{2}|\bar{y}|^{\alpha_{\mu,\nu-1}},$$

$$K_{1}|\bar{y}|^{\alpha_{\mu,\nu-2}} < \left|\partial_{\bar{y}}^{2}h_{n,\bar{\mu},\bar{\nu}}(\bar{y})\right| < K_{2}|\bar{y}|^{\alpha_{\mu,\nu-2}}.$$
(3)

Proof. – Take $n \in \mathbb{N}$. Given $\mu, \nu > 0$ define

$$c_n(\mu, \nu) = (k_{0,0}(1)\rho_{\mu,\nu}^n)^{\frac{1}{1-\alpha_{\mu,\nu}}}, \quad c_n = c_n(\mu, \nu) > 0.$$

Since $\alpha_{\mu,\nu} > 1$ we have

$$\rho_{\mu,\nu}^{n} c_{n}^{\alpha_{\mu,\nu-1}} = \frac{1}{k_{0,0}(1)} \quad \text{and} \quad c_{n} \to 0 \quad \text{as } n \to \infty.$$
(4)

Now, define

$$\bar{\mu} = c_n^{-1} (\rho_{\mu,\nu}^n \mu - 1) \quad \text{and} \quad \bar{\nu} = c_n^{-1} (\rho_{\mu,\nu}^n \nu - 1).$$
(5)

A straightforward computation give us

$$\begin{aligned} \partial_{\mu}\bar{\mu} &= c_{n}^{-2} \left[\left(n\rho_{\mu,\nu}^{n-1} \partial_{\mu}(\rho_{\mu,\nu}) \mu + \rho_{\mu,\nu}^{n} \right) c_{n} - \left(\rho_{\mu,\nu}^{n} \mu - 1 \right) \partial_{\mu} c_{n} \right], \\ \partial_{\nu}\bar{\mu} &= c_{n}^{-2} \left[n\rho_{\mu,\nu}^{n-1} \partial_{\nu}(\rho_{\mu,\nu}) \nu c_{n} - \left(\rho_{\mu,\nu}^{n} \nu - 1 \right) \partial_{\nu} c_{n} \right], \\ \partial_{\mu}\bar{\nu} &= c_{n}^{-2} \left[n\rho_{\mu,\nu}^{n-1} \partial_{\mu}(\rho_{\mu,\nu}) \nu c_{n} - \left(\rho_{\mu,\nu}^{n} \nu - 1 \right) \partial_{\mu} c_{n} \right], \\ \partial_{\nu}\bar{\nu} &= c_{n}^{-2} \left[\left(n\rho_{\mu,\nu}^{n-1} \partial_{\nu}(\rho_{\mu,\nu}) \nu + \rho_{\mu,\nu}^{n} \right) c_{n} - \left(\rho_{\mu,\nu}^{n} \nu - 1 \right) \partial_{\nu} c_{n} \right], \end{aligned}$$
(6)

and taking μ'_n and ν'_n as in the claim 1, then for $\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)} = \frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}(\mu'_n,\nu'_n)$ we obtain

876 E.M. MUÑOZ MORALES ET AL. / Ann. I. H. Poincaré – AN 20 (2003) 867–888

$$\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)} = \frac{1}{c_n(\mu'_n,\nu'_n)} \begin{bmatrix} n\rho_{\mu'_n,\nu'_n}^{-1}\partial_{\mu}(\rho_{\mu,\nu}) + \rho_{\mu'_n,\nu'_n}^n & n\rho_{\mu'_n,\nu'_n}^{-1}\partial_{\nu}(\rho_{\mu,\nu}) \\ n\rho_{\mu'_n,\nu'_n}^{-1}\partial_{\mu}(\rho_{\mu,\nu}) & n\rho_{\mu'_n,\nu'_n}^{-1}\partial_{\nu}(\rho_{\mu,\nu}) + \rho_{\mu'_n,\nu'_n}^n \end{bmatrix}.$$

From here

$$\det\left[\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}\right] = \left(c_n(\mu'_n,\nu'_n)\right)^{-2}\rho_{\mu'_n,\nu'_n}^{2n}\left[n\rho_{\mu'_n,\nu'_n}^{-n-1}\left(\partial_{\mu}(\rho_{\mu,\nu}) + \partial_{\nu}(\rho_{\mu,\nu})\right) + 1\right]$$

and then

$$(c_n(\mu'_n,\nu'_n))^2 \rho_{\mu'_n,\nu'_n}^{-2n} \det\left[\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}\right] \approx 1.$$

This implies $(\mu, \nu) \rightarrow (\bar{\mu}, \bar{\nu})$ is invertible in a neighbourhood of (μ'_n, ν'_n) .

Now, note that for $\|(\bar{\mu}, \bar{\nu})\| < r$, $\max\{\mu, \nu\} < 2\rho_{\mu,\nu}^{-n}$ for *n* large enough depending on *r*. Hence the domain of the map $(\bar{\mu}, \bar{\nu}) \rightarrow (\mu, \nu)$ converges to \mathbb{R}^2 , and then $\mu = \mu_n(\bar{\mu}, \bar{\nu})$ and $\nu = \nu_n(\bar{\mu}, \bar{\nu})$ is a well defined reparametrization in any large compact set.

Next, we define a map $\psi_{n,\bar{\mu},\bar{\nu}}$ from \mathbb{R} to \mathbb{R} by $y = \psi_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = c_n \bar{y} + 1$, where $c_n = c_n (\mu, \nu)$ (hence c_n depends on $\bar{\mu}$ and $\bar{\nu}$), and we also define $h_{n,\bar{\mu},\bar{\nu}} = (\psi_{n,\bar{\mu},\bar{\nu}})^{-1} \circ g_{n,\mu,\nu} \circ \psi_{n,\bar{\mu},\bar{\nu}}$. From (4) and (5) and a direct evaluation we obtain

$$h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = \begin{cases} -\frac{k_{\mu,\nu}(y)}{k_{0,0}(1)}(-\bar{y})^{\alpha_{\mu,\nu}} + \bar{\mu}, & \frac{a-1}{c_n} \leqslant \bar{y} < 0, \\ -\frac{\bar{k}_{\mu,\nu}(y)}{k_{0,0}(1)}(\bar{y})^{\alpha_{\mu,\nu}} + \bar{\nu}, & 0 < \bar{y} \leqslant \frac{b-1}{c_n} \end{cases}$$

and

$$\partial_{\bar{y}}h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = \begin{cases} \left(\alpha_{\mu,\nu}\frac{k_{\mu,\nu}(y)}{k_{0,0}(1)} + \frac{\partial_{y}k_{\mu,\nu}(y)}{k_{0,0}(1)}c_{n}\bar{y}\right)(-\bar{y})^{\alpha_{\mu,\nu}-1}, & \frac{a-1}{c_{n}} \leqslant \bar{y} < 0, \\ \left(-\alpha_{\mu,\nu}\frac{\bar{k}_{\mu,\nu}(y)}{k_{0,0}(1)} - \frac{\partial_{y}\bar{k}_{\mu,\nu}(y)}{k_{0,0}(1)}c_{n}\bar{y}\right)\bar{y}^{\alpha_{\mu,\nu}-1}, & 0 < \bar{y} \leqslant \frac{b-1}{c_{n}}. \end{cases}$$

When *n* tends to infinite, the parameters $\mu = \rho_{\mu,\nu}^{-n}(c_n\bar{\mu}+1)$ and $\nu = \rho_{\mu,\nu}^{-n}(c_n\bar{\nu}+1)$, and c_n go to 0, and thus, from the expression above we obtain that $h_{n,\bar{\mu},\bar{\nu}}$ converges C^1 uniformly on any compact set in the \bar{y} -space to $h_{\bar{\mu},\bar{\nu}}$. In a similar way C^2 convergence on \bar{y} is obtained.

The $(\bar{\mu}, \bar{\nu})$ -convergence in the C^0 sense is clear from the formulas. For the C^1 and C^2 $(\bar{\mu}, \bar{\nu})$ -convergence we need to compute $\partial_r (\partial_{\bar{y}}^i h_{n,\bar{\mu},\bar{\nu}}(\bar{y}))$ for i = 0, 1, 2 and $r = \bar{\mu}, \bar{\nu}$. We will compute only $\partial_{\bar{\mu}} h_{n,\bar{\mu},\bar{\nu}}(\bar{y})$. The other ones go on in a similar way and we left them to the reader.

For $\bar{y} < 0$, a straightforward computation give us

$$\partial_{\bar{\mu}}h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = \frac{\partial_{y}k_{\mu,\nu}(y)}{k_{0,0}(1)} [\partial_{\mu}c_{n}\partial_{\bar{\mu}}\mu + \partial_{\nu}c_{n}\partial_{\bar{\mu}}\nu](-\bar{y})^{\alpha_{\mu,\nu}+1} - \left[\partial_{\mu}k_{\mu,\nu}(c_{n}\bar{y}+1)\partial_{\bar{\mu}}\mu + \partial_{\nu}k_{\mu,\nu}(y)\partial_{\bar{\mu}}\nu\right](-\bar{y})^{\alpha_{\mu,\nu}} - \frac{k_{\mu,\nu}(y)}{k_{0,0}(1)}(-\bar{y})^{\alpha_{\mu,\nu}}[\partial_{\mu}\alpha_{\mu,\nu}\partial_{\bar{\mu}}\mu + \partial_{\nu}\alpha_{\mu,\nu}\partial_{\bar{\mu}}\nu]\log(-\bar{y}) + 1,$$

and for $\bar{y} > 0$

$$\begin{aligned} \partial_{\bar{\mu}}h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) &= -\frac{\partial_{y}\tilde{k}_{\mu,\nu}(y)}{k_{0,0}(1)} [\partial_{\mu}c_{n}\partial_{\bar{\mu}}\mu + \partial_{\nu}c_{n}\partial_{\bar{\mu}}\nu]\bar{y}^{\alpha_{\mu,\nu}+1} \\ &- \left[\partial_{\mu}\tilde{k}_{\mu,\nu}(y)\partial_{\bar{\mu}}\mu + \partial_{\nu}\tilde{k}_{\mu,\nu}(y)\partial_{\bar{\mu}}\nu\right]\bar{y}^{\alpha_{\mu,\nu}} \\ &- \frac{\tilde{k}_{\mu,\nu}(y)}{k_{0,0}(1)}\bar{y}^{\alpha_{\mu,\nu}}[\partial_{\mu}\alpha_{\mu,\nu}\partial_{\bar{\mu}}\mu + \partial_{\nu}\alpha_{\mu,\nu}\partial_{\bar{\mu}}\nu]\log(\bar{y}) + 1. \end{aligned}$$

We need to estimate $\partial_{\mu}c_n$, $\partial_{\nu}c_n$, $\partial_{\bar{\mu}}\mu$ and $\partial_{\bar{\mu}}\nu$. A direct computation gives:

$$\partial_{\mu}c_{n} = c_{n} \left[\frac{1}{1 - \alpha_{\mu,\nu}} n \rho_{\mu,\nu}^{-1} \partial_{\mu} \rho_{\mu,\nu} + \frac{\partial_{\mu} \alpha_{\mu,\nu}}{(1 - \alpha_{\mu,\nu})^{2}} \log(k_{0,0} \rho_{\mu,\nu}^{n}) \right]$$

and

$$\partial_{\nu}c_{n} = c_{n} \left[\frac{1}{1 - \alpha_{\mu,\nu}} n \rho_{\mu,\nu}^{-1} \partial_{\nu} \rho_{\mu,\nu} + \frac{\partial_{\nu} \alpha_{\mu,\nu}}{(1 - \alpha_{\mu,\nu})^{2}} \log(k_{0,0} \rho_{\mu,\nu}^{n}) \right].$$

From here we see that $\partial_{\mu}c_n = c_n U(n)$ and $\partial_{\nu}c_n = c_n V(n)$, where $U(n) = np(\mu, \nu) + q(\mu, \nu)$ and $V(n) = n\tilde{p}(\mu, \nu) + \tilde{q}(\mu, \nu)$ with $p(\mu, \nu)$, $q(\mu, \nu)$, $\tilde{p}(\mu, \nu)$, $\tilde{q}(\mu, \nu)$ bounded functions. From this and (6) we obtain

$$\det\left[\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}\right] = \left(\frac{\rho_{\mu,\nu}^n}{c_n}\right)^2 \left[1 + \theta(n)\right]$$

where $\theta(n)$ depends on μ and ν , and goes to zero when *n* goes to infinite. From this and the formula

$$\frac{\partial(\mu,\nu)}{\partial(\bar{\mu},\bar{\nu})} = \left[\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}\right]^{-1} = \left(\det\left[\frac{\partial(\bar{\mu},\bar{\nu})}{\partial(\mu,\nu)}\right]\right)^{-1} \left[\begin{array}{cc}\partial_{\nu}\bar{\nu} & -\partial_{\nu}\bar{\mu}\\ -\partial_{\mu}\bar{\nu} & \partial_{\mu}\bar{\mu}\end{array}\right]$$

we obtain

$$\partial_{\bar{\mu}}\mu = \left(\frac{c_n}{\rho_{\mu,\nu}^n}\right)^2 \frac{1}{1+\theta(n)} \partial_{\nu}\bar{\nu} \quad \text{and} \quad \partial_{\bar{\mu}}\nu = \left(\frac{c_n}{\rho_{\mu,\nu}^n}\right)^2 \frac{-1}{1+\theta(n)} \partial_{\mu}\bar{\nu},$$

from this expression and (6) we conclude that

$$\partial_{\bar{\mu}}\mu \approx \frac{c_n}{\rho_{\mu,\nu}^n}$$
 and $\partial_{\bar{\mu}}\nu \approx \frac{c_n}{\rho_{\mu,\nu}^n}$

and consequently $\partial_{\bar{\mu}}\mu$ and $\partial_{\bar{\mu}}\nu$ go to zero when *n* goes to infinite and the $C^1 - (\mu, \nu)$ convergence follows.

To finish the proof of lemma, for $c_n^{-1}(a-1) \leq \bar{y} < 0$, recalling that $y = c_n \bar{y} + 1$ and using definition for $h_{n,\bar{u},\bar{v}}(\bar{y})$, we have

$$\partial_{\bar{y}}h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = \partial_{\bar{y}} ((\psi_{n,\bar{\mu},\bar{\nu}})^{-1} \circ g_{n,\mu,\nu} \circ \psi_{n,\bar{\mu},\bar{\nu}})(\bar{y})$$

$$= \rho_{\mu,\nu}^{n} \partial_{y} f_{\mu,\nu}(y)$$

$$= \rho_{\mu,\nu}^{n} \frac{\partial_{y} f_{\mu,\nu}(y)}{(1-y)^{\alpha_{\mu,\nu}-1}} (1-y)^{\alpha_{\mu,\nu}-1},$$

but then $\partial_{\bar{y}} h_{n,\bar{\mu},\bar{\nu}}(\bar{y}) = \frac{1}{k_{0,0}(1)} \frac{\partial_{y} f_{\mu,\nu}(y)}{|1-y|^{\alpha_{\mu,\nu}-1}} |\bar{y}|^{\alpha_{\mu,\nu}-1}.$

The same conclusion is obtained for $0 < \overline{y} < c_n^{-1}(b-1)$. Now, from (2) the first inequality in the lemma follows. The same argument works for the second inequality. \Box

3. Proof of the main theorem

In this section we will prove a one dimensional version of Theorem 1 from which Theorem 1 follows.

THEOREM 2. – For each $n \in \mathbb{N}$ large enough there exists a Lebesgue positive measure set E_n in the (μ, ν) -parameter space such that:

- (1) For every $(\mu, \nu) \in E_n$, the critical values $\rho_{\mu,\nu}^n \mu$ and $\rho_{\mu,\nu}^n \nu$ of the function $g_{n,\mu,\nu} = (f_{\mu,\nu})^{n+1}$ have positive Lyapunov exponents, in other words, there exists $\lambda > 1$ such that $|((g_{n,\mu,\nu})^k)'(\rho_{\mu,\nu}^n \mu)| > \lambda^k$, and $|((g_{n,\mu,\nu})^k)'(\rho_{\mu,\nu}^n \nu)| > \lambda^k$ for every k > 0.
- (2) For every $(\mu, \nu) \in E_n$, both critical orbits, i.e. $\{((g_{n,\mu,\nu})^k)'(\rho_{\mu,\nu}^n\mu)\}_{k\in\mathbb{N}}$ and $\{((g_{n,\mu,\nu})^k)'(\rho_{\mu,\nu}^n\nu)\}_{k\in\mathbb{N}}$ are dense subsets of the interval $\Theta_{\mu,\nu} = [g_{n,\mu,\nu}(\eta), \eta]$ where $\eta = \max\{\rho_{\mu,\nu}^n\mu, \rho_{\mu,\nu}^n\nu\}$.

Remark 2. – Note that our Theorem 1 follows from Theorem 2. In fact, for each $(\mu, \nu) \in E_n$, the set $Q_{\mu,\nu} = \bigcap_{T>0} \bigcup_{t>T} X^t_{\mu,\nu}(\pi^{-1}(\Theta_{\mu,\nu}))$ is the announced attractor since $\Theta_{\mu,\nu}$ is transitive and the foliation defining $f_{\mu,\nu}$ is contracting.

The proof is quite long and we need to introduce some terminology. For a positive large integer *n*, define the μ -dependent points $p(\mu)$ and $q(\mu)$ by the equations $g_{n,\mu,\mu}(p(\mu)) = p(\mu) = g_{n,\mu,\mu}(q(\mu))$.

Note that $a(\mu, \mu) < p(\mu) < 1 < q(\mu) < b(\mu, \mu)$. Let μ_1, μ_2 be defined by

 $\mu_1 = \sup\{\mu \mid \rho_{\mu,\mu}^n \mu \leqslant 1\} \quad \text{and} \quad \mu_2 = \inf\{\mu \mid b(\mu,\mu) \leqslant \rho_{\mu,\mu}^n \mu\}.$

The map $w: [\mu_1, \mu_2] \to \mathbb{R}$ defined by $w(\mu) = \rho_{\mu,\mu}^n \mu - q(\mu)$ is continuous, $w(\mu_1) < 0$ and $w(\mu_2) > 0$, hence there is a parameter value $\mu_n \in [\mu_1, \mu_2]$ such that $\rho_{\mu_n,\mu_n}^n \mu_n = q(\mu_n)$, see Fig. 4.



From the definitions of p_n, q_n, μ_n and g_{n,μ_n,μ_n} , we have that for *n* large enough we get

$$1 - p_n = \left(\frac{k_{\mu_n,\mu_n}(p_n)}{\tilde{k}_{\mu_n,\mu_n}(q_n)}\right)^{1/\alpha} (q_n - 1)$$

and

$$1 - p_n = \left(\rho_{\mu_n,\mu_n}^{-n} \left[\left(\frac{k_{\mu_n,\mu_n}(p_n)}{\tilde{k}_{\mu_n,\mu_n}(q_n)}\right)^{1/\alpha} + 1 \right] \frac{1}{k_{\mu_n,\mu_n}(p_n)} \right)^{1/(\alpha-1)}$$

hence, using (4) $\frac{1-p_n}{c_n}$ and $\frac{q_n-1}{c_n}$ are uniformly bounded. This implies that there are parameter values $\bar{\mu}_n = \bar{\nu}_n$ and points $\bar{q}_n = \bar{\mu}_n$, \bar{p}_n in the $(\bar{y}, \bar{\mu}, \bar{\nu})$ -space, such that $h_{n,\bar{\mu}_n,\bar{\nu}_n}(\bar{p}_n) = \bar{p}_n = h_{n,\bar{\mu}_n,\bar{\nu}_n}(\bar{q}_n)$.

Furthermore, from the definition of $h_{\bar{\mu},\bar{\nu}}$ we can see that there are a unique parameter value $\bar{\mu}_{\infty} = \bar{\nu}_{\infty}$ and points $\bar{q} = \bar{\nu}_{\infty}$ and \bar{p} such that $h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}}(\bar{p}) = \bar{p} = h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}}(\bar{q})$ and since $h_{n,\bar{\mu},\bar{\nu}}$ converges uniformly to $h_{\bar{\mu},\bar{\nu}}$ on any compact set in \mathbb{R} we conclude that $\bar{p}_n \to \bar{p}$, $\bar{q}_n \to \bar{q}$ and $\bar{\mu}_n = \bar{\nu}_n \to \bar{\mu}_{\infty} = \bar{\nu}_{\infty}$ when $n \to \infty$.

For each $\theta > 0$ consider the straight line $L_{n,\theta}$ in the $(\bar{\mu}, \bar{\nu})$ -plane given by

$$L_{n,\theta} = \left\{ (\bar{\mu}, \bar{\nu}) \mid \bar{\nu} = \theta (\bar{\mu} - \bar{\mu}_n) + \bar{\nu}_n \right\}.$$

Let *m* and m_{θ} be the Lebesgue measure in the $(\bar{\mu}, \bar{\nu})$ -plane and in the straight line $L_{n,\theta}$ respectively. Using the Benedick–Carleson techniques, for *n* large (see [1], [2, Section 2], [10, Section 3], [15]) we will be able to show the existence of a m_{θ} -positive measure subset $\tilde{E}_{n,\theta}$ of $L_{n,\theta}$, having $(\bar{\mu}_n, \bar{\nu}_n)$ as a density point and satisfying the requisites for Theorem 2 in the $(\bar{\mu}, \bar{\nu})$ -parameter setting. Once we have done this, we consider $\tilde{E}_n = \bigcup_{\theta} \tilde{E}_{n,\theta}$. Fubini's theorem implies $m(\tilde{E}_n) > 0$. Now, since the reparametrization $(\mu, \nu) \rightarrow (\bar{\mu}, \bar{\nu})$ is a diffeomorphism, we get a set E_n in the (μ, ν) plane satisfying the requirements of Theorem 2 because $\psi_{n,\bar{\mu},\bar{\nu}}$ is an affine map.

Before begining the set up of the $E_{n,\theta}$, following the steps on [15] we will prove that the maximal orbits outside of a neighbourhood of the critical point, having length bigger than K_0 (some K_0) have exponential growth. This fact is established in Proposition 2 and it is a key step in order to apply Benedick–Carleson techniques. In Lemma 5 we will show that under the basic assumption (BA) the increasing in the derivative after a binding period (see definitions before Lemma 5) fully compensates the small factor introduced in the derivative by the orbit point of the critical value passing close to the critical point. As a consequence, we can establish via an inductive argument conditions (condition (FA), given before Proposition 3) to guarantee the exponential growing of the derivatives at the critical values. This is stated in Proposition 3.

To prove that the maximal orbits outside of a neighbourhood of the critical point have exponential growth we follow [15]. The nonsymmetry of the invariant interval is the main difference here.

LEMMA 3. – There exists a constant C > 0 such that, for every small δ_0 there are constants $\lambda_1 = \lambda_1(\delta_0) > 1$, $N_1 = N_1(\delta_0)$, $K_0 = K_0(\delta_0)$ and $a_1 = a_1(\delta_0)$ satisfying the following property: For every $n > N_1$, for every $(\bar{\mu}, \bar{\nu})$ such that $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_1$, and \bar{y} such that $|h_{n,\bar{\mu},\bar{\nu}}^k(\bar{y})| \ge \delta_0$, $i = 0, \ldots, k - 1$, but $|h_{n,\bar{\mu},\bar{\nu}}^k(\bar{y})| < \delta_0$, then 880 E.M. MUÑOZ MORALES ET AL. / Ann. I. H. Poincaré – AN 20 (2003) 867–888

$$\left| \left(h_{n,\bar{\mu},\bar{\nu}}^k \right)'(\bar{y}) \right| > \begin{cases} C & \text{if } k < K_0, \\ \lambda_1^k & \text{if } k \geqslant K_0. \end{cases}$$

Proof. – Given $\varepsilon > 0$ small, $0 < \delta_0 < \frac{1}{2} \min\{|p|, q\}$ and $k \in \mathbb{N}$, let $W = [p - \varepsilon, q + \varepsilon] \setminus (-\delta_0, \delta_0)$, $U_k(\bar{\mu}, \bar{\nu}) = \{\bar{y} \in W \mid h^i_{\bar{\mu}, \bar{\nu}}(\bar{y}) \in W, i = 0, \dots, k\}$ and $U_k(n, \bar{\mu}, \bar{\nu}) = \{\bar{y} \in W \mid h^i_{n, \bar{\mu}, \bar{\nu}}(\bar{y}) \in W, i = 0, \dots, k\}$.

The set $\Delta = \bigcap_{k \ge 0} U_k(\bar{\mu}_{\infty}, \bar{\nu}_{\infty})$ is compact, $h_{\bar{\mu}_{\infty}, \bar{\nu}_{\infty}}$ -invariant, without sinks (Singer's Theorem), its periodic points are hyperbolic (Minimum Principle) and away from the critical point. From Mañé's Theorem [8] (see also Misiurewicz's Theorem in the same reference), we have that Δ is hyperbolic. This means that there are constants $\tilde{C} > 0$ and $\lambda > 1$ such that for every $\bar{y} \in \Delta$ and $k \in \mathbb{N}$ we have $|(h_{\bar{\mu}_{\infty}, \bar{\nu}_{\infty}}^k)'(\bar{y})| \ge \tilde{C}\lambda^k$. From this we can conclude that there are constants $\lambda_1 > 1$, k_0 , N_1 and a_1 depending on δ_0 , and an open neighbourhood V of Δ such that for every $k \ge k_0$, $n \ge N_1$ and $(\bar{\mu}, \bar{\nu})$ satisfying $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_1$ then

$$\forall \bar{y}, \ h_{n,\bar{\mu},\bar{\nu}}^{i}(\bar{y}) \in V, \ 0 \leq i < k-1 \quad \Rightarrow \quad \left| \left(h_{n,\bar{\mu},\bar{\nu}}^{k} \right)^{\prime}(\bar{y}) \right| \geq \lambda_{1}^{k}.$$

$$\tag{7}$$

Note that we can take $k_0 = \min\{k \mid \tilde{C}\lambda^k > 1\}$ and $1 < \lambda_1 < \lambda(\tilde{C})^{1/k_0}$.

Furthermore, since $U_{k+1}(\bar{\mu}_{\infty}, \bar{\nu}_{\infty}) \subset U_k(\bar{\mu}_{\infty}, \bar{\nu}_{\infty})$ for a_1 small enough and N_1 large, there is a k_1 (depending only on δ_0) such that $U_{k_1}(n, \bar{\mu}, \bar{\nu}) \subset V \cap W$. Let $\bar{y} \in [p - \varepsilon, q + \varepsilon] \setminus V$ (and hence $\bar{y} \notin U_{k_1}(n, \bar{\mu}, \bar{\nu})$), then there exists a time $j < k_1$ such that $|h_{n,\bar{\mu},\bar{\nu}}^j(\bar{y})| < \delta_0$. Let consider $K_0 = k_0 + k_1$ and $C = \frac{1}{2} \min\{\frac{|p|}{q}, \frac{q}{|p|}\}$. We have the following claim: If \bar{y} is such that for some $i \leq K_0$ and $|h_{n,\bar{\mu},\bar{\nu}}^i(\bar{y})| < \delta_0$ then

$$\left| \left(h_{n,\bar{\mu},\bar{\nu}}^{i} \right)'(\bar{y}) \right| \ge C.$$
(8)

In fact, the Minimum Principle applied to $h = h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}}$ gives that $|(h^i)'(\bar{y})| > \frac{3}{2}C$, indeed, suppose that for some $\bar{y}, |h^i(\bar{y})| < \delta_0$ and $|(h^i)'(\bar{y})| < \frac{3}{2}C$. Let $[\bar{y}_1, \bar{y}_2]$ be the small interval containing \bar{y} such that \bar{y}_1 and \bar{y}_2 are local extremes of h^i . From the choice of C and δ_0 we have that the map $\bar{y} \to |(h^i)'(\bar{y})|$ has a minimum in (\bar{y}_1, \bar{y}_2) but, by the Minimun Principle this is not possible. Taking a_1 small enough and N_1 large the claim follows.

Now, let \bar{y} be as in the statement of the lemma, then for a suitable choice of m, indeed $m = \min\{\{k\} \cup \{i/h_{n,\bar{\mu},\bar{\nu}}^i(\bar{y}) \notin V\}\}$ we have that $h_{n,\bar{\mu},\bar{\nu}}^m(\bar{y}) \notin U_{k_1}(n,\bar{\mu},\bar{\nu})$ and then $j = k - m < k_1$. Applying the chain rule, (7) and (8) we obtain

$$\begin{split} \left| \left(h_{n,\bar{\mu},\bar{\nu}}^k \right)'(\bar{y}) \right| &\geqslant \begin{cases} C & \text{if } m < k_0, \\ |(h_{n,\bar{\mu},\bar{\nu}}^j)'_{(h_{n,\bar{\mu},\bar{\nu}}^m(\bar{y}))}||(h_{n,\bar{\mu},\bar{\nu}}^m)'(\bar{y})| & \text{if } m \geqslant k_0, \end{cases} \\ &\geqslant \begin{cases} C & \text{if } m < k_0, \\ C\lambda_1^m & \text{if } m \geqslant k_0. \end{cases} \end{split}$$

Thus, for k_0 large enough we obtain (by decreasing λ_1), a new constant $\lambda_1 > 1$ such that

$$\left| \left(h_{n,\bar{\mu},\bar{\nu}}^k \right)'(\bar{y}) \right| \ge \begin{cases} C & \text{if } k < K_0, \\ \lambda_1^k & \text{if } k \ge K_0. \end{cases} \quad \Box$$

LEMMA 4. – Given δ_0 small enough and C > 0 there exists $\lambda_2 > 1$, not depending on δ_0 , with the following property: for $\delta < \delta_0$ there are $N_2 = N_2(\delta) \in \mathbb{N}$ and $a_2 = a_2(\delta) > 0$ such that for every $n > N_2$, for any $(\bar{\mu}, \bar{\nu})$ with $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_2$, and $\bar{\nu}$, $\delta \leq |\bar{\nu}| < \delta_0$, there is $l = l(a_2, \bar{\nu}, \delta_0)$ such that $|h_{n,\bar{\mu},\bar{\nu}}^j(\bar{\nu})| > \delta_0$ for j = 1, ..., l and $|(h_{n,\bar{\mu},\bar{\nu}}^l)'(\bar{\nu})| > \lambda_2^l$. Furthermore, there are constants $L = L(\delta_0)$ and $M = M(\delta, \delta_0)$ such that L < l < M and besides $C\lambda_2^L > 1$.

Proof. – Let define $M_{\varepsilon} = M_{\varepsilon}(n, \bar{\mu}, \bar{\nu}) = \max\{|(h_{n,\bar{\mu},\bar{\nu}})'(x)|: p_n - \varepsilon \leq x \leq p_n + \varepsilon\}$ and $m_{\varepsilon} = m_{\varepsilon}(n, \bar{\mu}, \bar{\nu}) = \min\{|(h_{n,\bar{\mu},\bar{\nu}})'(x)|: p_n - \varepsilon \leq x \leq p_n + \varepsilon\}$. When *n* tends to infinite, $h_{n,\bar{\mu},\bar{\nu}}$ converges uniformly on compact sets to $h_{\bar{\mu},\bar{\nu}}$ and thus, given $\zeta > 0$, there are constants $\tau > 0, \bar{\varepsilon} > 0$ and $N_2 \in \mathbb{N}$ such that $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_{\infty}, \bar{\nu}_{\infty})| < \tau, \varepsilon < \bar{\varepsilon}, n \geq N_2 \Rightarrow 1 - \zeta < m_{\varepsilon}/M_{\varepsilon} < 1$.

From this, and since $\alpha = \alpha_{\mu,\nu} \approx \alpha_{0,0}$, we can choose constants $a_2 > 0$, $\varepsilon > 0$, $N_2 \in \mathbb{N}$ and $\lambda > 1$, such that $n \ge N_2$, $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_2 \Rightarrow m_{\varepsilon} M_{\varepsilon}^{(1-\alpha)/\alpha} > \lambda > 1$. Note that λ is independent of δ_0 . Next, let δ_0 be a small positive constant and fix δ , $0 < \delta < \delta_0$. Write $h = h_{n,\bar{\mu},\bar{\nu}}$. Given $\bar{y}, \delta \le \bar{y} \le \delta_0$, let l be the first positive integer such that $h^l(h^2(\bar{y})) > p_n + \varepsilon$. Notice that from the uniform convergence of $h_{n,\bar{\mu}_n,\bar{\nu}_n}$ to $h_{\bar{\mu}_\infty,\bar{\nu}_\infty}$, there are constants $N_2 \in \mathbb{N}$, $L = L(\delta_0)$ and $M = M(\delta, \delta_0)$ such that for $n \ge N_2$ we have L < l < M. We also observe that for N_2 large and a_2 small enough we can get $|h(0^{\pm}) - q_n| + |h^l(h(q_n)) - p_n|$ as small as we want because $h_{n,\bar{\mu}_n,\bar{\nu}_n}(q_n) = p_n$, where $h(0^{\pm}) = \lim_{\bar{\nu} \to 0^{\pm}} h(\bar{y})$.

Now, using the above comments, the chain rule, the mean value theorem and constants K_1 and K_2 given by inequality (3) in the renormalization lemma, we obtain the following estimates:

$$\begin{split} \varepsilon &< h^l \big(h^2(\bar{y}) \big) - p_n \\ &\leq \big| h^l \big(h^2(\bar{y}) \big) - h^l \big(h^2(0^{\pm}) \big) \big| + \big| h^l \big(h^2(0^{\pm}) \big) - h^l \big(h(q_n) \big) \big| + \big| h^l \big(h(q_n) \big) - p_n \big| \\ &\leq M_{\varepsilon}^l C_2 K_2 |\bar{y}|^{\alpha} + M_{\varepsilon}^l C_2 \big| h(0^{\pm}) - q_n \big| + \big| h^l \big(h(q_n) \big) - p_n \big|, \end{split}$$

where C_2 is a positive constant close to $|(h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}})'(q)|$. Solving this inequality for $|\bar{y}|$ we obtain

$$|\bar{y}| > \left\{ \frac{1}{M_{\varepsilon}^{l}C_{2}K_{2}} \left[\varepsilon - M_{\varepsilon}^{l}C_{2} |h(0^{\pm}) - q_{n}| - |h^{l}(h(q_{n})) - p_{n}| \right] \right\}^{1/\alpha}$$

Furthermore, for a positive constant C_1 close to $|(h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}})'(q)|$ we also have $|(h^{l+2})'(\bar{y})| \ge m_{\varepsilon}^{l} C_1 K_1 |\bar{y}|^{\alpha-1}$ and then,

$$\left| \left(h^{l+2} \right)'(\bar{y}) \right| \ge m_{\varepsilon}^{l} C_{1} K_{1} K_{2}^{\frac{1-\alpha}{\alpha}} \left\{ \frac{\varepsilon}{M_{\varepsilon}^{l} C_{2}} - \left| h(0^{\pm}) - q_{n} \right| - \frac{\left| h^{l}(h(q_{n})) - p_{n} \right|}{M_{\varepsilon}^{l} C_{2} K_{2}} \right\}^{\frac{\alpha-1}{\alpha}}.$$

Now, if we consider $\bar{\mu} = \bar{\mu}_n$, $\bar{\nu} = \bar{\nu}_n$ and *L* large (this last condition occurs with δ_0 small), then for some $\lambda_1 > 1$ the above inequality and the choice of ε leads us to

$$\begin{split} \left| \left(h^{l+2} \right)'(\bar{y}) \right| &\ge \left(m_{\varepsilon} M_{\varepsilon}^{\frac{1-\alpha}{\alpha}} \right)^{l} \varepsilon^{\frac{\alpha-1}{\alpha}} K_{2}^{\frac{1-\alpha}{\alpha}} C_{1} C_{2}^{\frac{1-\alpha}{\alpha}} K_{1} \\ &\geqslant \bar{C} \lambda^{l}, \quad \text{with } \bar{C} = \varepsilon^{\frac{\alpha-1}{\alpha}} K_{2}^{\frac{1-\alpha}{\alpha}} C_{1} C_{2}^{\frac{1-\alpha}{\alpha}} K_{1} \\ &\geqslant \lambda_{1}^{l+2}. \end{split}$$

Since L < l < M, taking a_2 small enough, we can find λ_2 , $1 < \lambda_2 < \lambda_1$, independent of δ_0 , such that for $(\bar{\mu}, \bar{\nu})$ with $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_2$ then $|(h^{l+2})'(\bar{\nu})| \ge \lambda_2^{l+2}$. Finally, given any positive constant *C* we can choose δ_0 small enough to obtain $C\lambda_2^L > 1$. This ends the proof of the lemma. \Box

PROPOSITION 2. – There exist $K_0 \in \mathbb{N}$, C > 0 and $\lambda_0 > 1$ satisfying the following property: For every $\delta > 0$ small enough, there exist $a_0 = a_0(\delta)$ and $N_0 = N_0(\delta)$ such that for every $n > N_0$, $(\bar{\mu}, \bar{\nu})$ with $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_0$, and \bar{y} such that $|h_{n,\bar{\mu},\bar{\nu}}^i(\bar{y})| \ge \delta$ for $i = 0, \ldots, k - 1$ but $|h_{n,\bar{\mu},\bar{\nu}}^k(\bar{y})| < \delta$, we have

$$\left| \left(h_{n,\bar{\mu},\bar{\nu}}^{k} \right)'(\bar{y}) \right| \ge \begin{cases} C & \text{if } k < K_{0}, \\ \lambda_{0}^{k} & \text{if } k \ge K_{0}. \end{cases}$$

Proof. – Take C > 0 as in Lemma 3 and take $\delta_0 > 0$ small enough so that Lemma 4 is true for the pair (C, δ_0) . Let $\lambda_1 = \lambda_1(\delta_0), a_1(\delta_0), N_1(\delta_0), K_0 = K_0(\delta_0)$ and $L = L(\delta_0)$ as in Lemma 3 applied to the pair (C, δ_0) . Given $\delta < \delta_0$, let λ_2 , $a_2(\delta)$ and $N_2(\delta)$ be determined by Lemma 4 applied to the pair (C, δ_0) . Define $a_0 = \min\{a_1(\delta_0), a_2(\delta)\}$ and $N_0 = \max\{N_1(\delta_0), N_2(\delta)\}$. Let $n > N_0$ and $(\bar{\mu}, \bar{\nu})$ with $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a_0$. Write $h = h_{n,\bar{\mu},\bar{\nu}}$, and let $\bar{\nu}$ be as in the statement of the proposition. We consider two possibilities: (i) $\delta \leq |\bar{y}| < \delta_0$ and (ii) $\delta_0 \leq |\bar{y}|$. In the first one we decompose the orbit block $\{h^i(\bar{y})\}_{i=0}^{i=k}$ in blocks $\{\bar{y} = y_0, h(y_0), \dots, h^{l_0}(y_0) =$ $\begin{array}{l} x_0\}, \ \{h(x_0), \dots, h^{k_0}(x_0) = y_1\}, \ \{h(y_1), \dots, h^{l_1}(y_1) = x_1\}, \ \{h(x_1), \dots, h^{k_1}(x_1) = y_2\}, \\ \{h(y_2), \dots, h^{l_2}(y_2) = x_2\}, \ \{h(x_2), \dots, h^{k_2}(x_2) = y_3\}, \dots, \{h(y_w), \dots, h^{l_w}(y_w) = x_w\}, \end{array}$ $\{h(x_w), \ldots, h^{k_w}(x_w) = h^k(\bar{y})\}$, where $l_i \ge L$ and k_i are inductively defined by applying Lemmas 4 and 3 to y_i and x_i respectively. Let k_{i_1}, \ldots, k_{i_r} be the integers satisfying $k_{i} \ge K_0$, and rewrite those integers k_i such that $k_i < K_0$ as $s_0, s_1, \ldots, s_{u-1}$ preserving the appearance order $l_0 \prec k_0 \prec l_1 \prec k_1 \prec \cdots \prec l_w \prec k_w$, that is, s_0 is the first integer k_i such that $k_i < K_0$, s_1 is the second one and so on. It is clear that $u \leq w, s_i < K_0$ and $k = (l_0 + s_0) + \dots + (l_{u-1} + s_{u-1}) + l_u + \dots + l_w + k_{i_1} + \dots + k_{i_r}$. Applying the chain rule, Lemmas 3 and 4 and after some rearrangement, we obtain

$$\left| \left(h^k \right)'(\bar{y}) \right| \ge \left(\lambda_2^{l_0} C \right) \cdots \left(\lambda_2^{l_{u-1}} C \right) \lambda_2^{l_u + \dots + l_w} \lambda_1^{k_{i_1} + \dots + k_{i_r}}$$

From Lemma 4 we have that if l > L then $C\lambda_2^L > 1$, and we obtain $\lim_{\lambda \to 1^+} l \log(\lambda_2/\lambda) = l \log \lambda_2 > -\log C = \lim_{\lambda \to 1^+} \log(\lambda^{K_0}/C)$, and consequently we can take $\lambda_0, 1 < \lambda_0 < \min\{\lambda_1, \lambda_2\}$ such that if l > L then $C\lambda_2^l > \lambda_0^{l+K_0}$. In particular, for $l_i > L$ and since $s_i < K_0, i = 0, \ldots, u - 1$, we have $C\lambda_2^{l_i} > \lambda_0^{l_i+s_i}$, and then $|(h^k)'(\bar{y})| \ge \lambda_0^k$.

Note that if $k < K_0$ then $k_1 = k$ and in this case we have $|(h^k)'(\bar{y})| \ge C$.

Finally, if $\delta_0 \leq |\bar{y}|$, the same arguments as in the $\delta \leq |\bar{y}| \leq \delta_0$ case work, the only difference here is that we do not have the first block $\{\bar{y} = y_0, h(y_0), \dots, h^{l_0}(y_0) = x_0\}$, and the second one must be changed by $\{\bar{y} = x_0, h(x_0), \dots, h^{k_0}(x_0) = y_1\}$. This ends the proof of proposition. \Box

Given $(\bar{\mu}, \bar{\nu}) \in L_{n,\theta} = \{(\bar{\mu}, \bar{\nu})/\bar{\nu} = \theta(\bar{\mu} - \bar{\mu}_n) + \bar{\mu}\}$ set $h_{\bar{\mu}} = h_{n,\bar{\mu},\bar{\nu}}$. For a small $\gamma > 0$, let define $\tilde{E}_{n,\theta}(\gamma) \subset L_{n,\theta}$ as the set of parameters $(\bar{\mu}, \bar{\nu})$ satisfying the following basic assumption:

$$\left|h_{\bar{\mu}}^{j}(\bar{\mu})\right| \ge e^{-\gamma j} \text{ and } \left|h_{\bar{\mu}}^{j}(\bar{\nu})\right| \ge e^{-\gamma j} \quad \forall j \ge 1.$$
 (BA)

Let $\beta > 0$, $(\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta}(\gamma)$, and $\delta > 0$ small. For a positive integer k such that $|h_{\bar{\mu}}^k(\bar{\mu})| \leq \delta$, define the *binding period* associated to the parameter $(\bar{\mu}, \bar{\nu})$ for the return $h_{\bar{\mu}}^k(\bar{\mu})$ of the critical value $\bar{\mu}$ as the maximal interval [k+1, k+s] such that for $1 \leq j \leq s$

$$|h_{\bar{\mu}}^{k+j}(\bar{\mu}) - h_{\bar{\mu}}^{j-1}(\bar{\mu})| < e^{-\beta j}, \text{ if } h_{\bar{\mu}}^{k}(\bar{\mu}) < 0,$$

or

$$\left|h_{\bar{\mu}}^{k+j}(\bar{\mu}) - h_{\bar{\mu}}^{j-1}(\bar{\nu})\right| < e^{-\beta j}, \text{ if } h_{\bar{\mu}}^{k}(\bar{\mu}) > 0$$

holds.

Thus, during the binding period, the orbit of $h_{\bar{\mu}}^{k+1}(\bar{\mu})$ is close to that of $\bar{\mu}$ or $\bar{\nu}$ depending on above conditions. In the same way, we define the *binding period* associated to the parameter $(\bar{\mu}, \bar{\nu})$ for the return $h_{\bar{\mu}}^k(\bar{\nu})$ of the critical value $\bar{\nu}$.

LEMMA 5. – For a convenient choosing of γ , β and δ , there are positive constants \tilde{K} , D, A, k_1 and τ depending only on γ and β , such that for both critical values $\bar{\eta} = \bar{\mu}$ or $\bar{\eta} = \bar{\nu}$ of $h_{\bar{\mu}} = h_{n,\bar{\mu},\bar{\nu}}$, with $(\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta}(\gamma)$, we have: If $h_{\bar{\mu}}^k(\bar{\eta}) \in (e^{-\gamma k}, \delta)$ for some $k > k_1$ and $|(h_{\bar{\mu}}^j)'(\bar{\eta})| > \lambda_1^j$, $j = 1, \ldots, k - 1$ for some λ_1 , $1 < \lambda_1 < \lambda_0$ (λ_0 as in Proposition 2), then

(1)

$$\frac{1}{A} < \frac{\left|(h_{\bar{\mu}}^{j})'(x)\right|}{\left|(h_{\bar{\mu}}^{j})'(y)\right|} < A \quad for \ all \ x, \ y \in \left[h_{\bar{\mu}}^{k+1}(\bar{\eta}), \bar{\nu}\right] \ and \ 1 \leqslant j \leqslant s,$$

where *s* is the binding period associated to the parameter $(\bar{\mu}, \bar{\nu})$ for the return $h^k_{\bar{\mu}}(\bar{\eta})$ of the critical value $\bar{\eta}$ of $h_{\bar{\mu}}$.

(2)

$$\frac{r}{\beta + \log D} - 1 \leqslant s \leqslant \frac{r\alpha_{\mu,\nu} - \log(\bar{K}A^{-1})}{\beta + \log\lambda_1}, \quad where \ e^{-r} = \left|h_{\bar{\mu}}^k(\bar{\eta})\right|.$$

(3)

$$\left|\left(h_{\bar{\mu}}^{s+1}\right)'\left(h_{\bar{\mu}}^{k}(\bar{\eta})\right)\right| \ge \tau \exp\left[\left(\frac{\log\lambda_{1}}{\alpha_{\mu,\nu}} - \frac{\alpha_{\mu,\nu} - 1}{\alpha_{\mu,\nu}}\beta\right)(s+1)\right] > 1.$$

Similar results can be obtained if $h^k_{\bar{\mu}}(\bar{\eta}) \in (-\delta, -e^{-k\gamma})$, but in this case we have to change $[h^{k+1}_{\bar{\mu}}(\bar{\eta}), \bar{\nu}]$ by $[h^{k+1}_{\bar{\mu}}(\bar{\eta}), \bar{\mu}]$ in conclusion (1) above.

Proof. – First, we observe that taking $\gamma < \beta$ then $e^{-(j+1)\beta} < e^{-j\gamma}$. Since $|h_{\bar{\mu}}^{j}(\bar{\mu})| > e^{-j\gamma}$ and *s* is the binding period associated to $(\bar{\mu}, \bar{\nu})$ for the return $h_{\bar{\mu}}^{k}(\bar{\eta})$ of the critical value $\bar{\eta}$ of $h_{\bar{\mu}}$, we conclude that $0 \notin (h_{\bar{\mu}}^{k+1+j}(\bar{\eta}), h_{\bar{\mu}}^{j}(\bar{\nu}))$ for j < s. Consequently, if $x \in (h_{\bar{\mu}}^{k+1}(\bar{\eta}), \bar{\nu})$ then $0 \notin (h_{\bar{\mu}}^{j}(x), h_{\bar{\mu}}^{j}(\bar{\nu})) \subset (h_{\bar{\mu}}^{k+1+j}(\bar{\eta}), h_{\bar{\mu}}^{j}(\bar{\nu}))$. From the second inequality in (3), for all $j \leq s - 1$ we can find $\xi_{j} \in (h_{\bar{\mu}}^{j}(x), h_{\bar{\mu}}^{j}(\bar{\nu}))$ such that

$$\begin{aligned} \left| h_{\bar{\mu}}^{'} \big(h_{\bar{\mu}}^{j}(x) \big) - h_{\bar{\mu}}^{'} \big(h_{\bar{\mu}}^{j}(\bar{\nu}) \big) \right| &\leq K_{2} |\xi|^{\alpha_{\mu,\nu}-2} \left| h_{\bar{\mu}}^{j}(x) - h_{\bar{\mu}}^{j}(\bar{\nu}) \right| \\ &\leq K_{2} |\xi|^{\alpha_{\mu,\nu}-2} \left| h^{k+1+j}(\bar{\eta}) - h^{j}(\bar{\nu}) \right|, \end{aligned}$$

and from here

$$\begin{aligned} \left| h_{\bar{\mu}}^{'} \left(h_{\bar{\mu}}^{j}(x) \right) - h_{\bar{\mu}}^{'} \left(h_{\bar{\mu}}^{j}(\bar{\nu}) \right) \right| &\leq K_{2} \mathrm{e}^{-\gamma(\alpha_{\mu,\nu}-2)j} \left| h^{k+1+j}(\bar{\eta}) - h^{j}(\bar{\nu}) \right| \\ &\leq K_{2} \mathrm{e}^{-\gamma(\alpha_{\mu,\nu}-2)j - (j+1)\beta}. \end{aligned}$$

Furthermore, $|h'_{\bar{\mu}}(h^{j}_{\bar{\mu}}(\bar{\eta}))| \ge K_{1}|h^{j}_{\bar{\mu}}(\bar{\eta})|^{\alpha_{\mu,\nu}-1} \ge K_{1}e^{-j\gamma(^{\alpha_{\mu,\nu}-1})}$. Now,

$$\begin{split} \left| \frac{(h_{\bar{\mu}}^{j})'(x)}{(h_{\bar{\mu}}^{j})'(\bar{\eta})} \right| &= \prod_{i=0}^{j-1} \left[1 + \frac{|h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(x))| - |h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(\bar{\eta}))|}{|h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(\bar{\eta}))|} \right] \\ &\leqslant \exp \left[\sum_{i=0}^{j-1} \frac{|h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(x)) - h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(\bar{\eta}))|}{|h_{\bar{\mu}}^{'}(h_{\bar{\mu}}^{i}(\bar{\eta}))|} \right] \\ &\leqslant \exp \left[\sum_{i=0}^{j-1} \frac{K_{2} e^{-\gamma(\alpha_{\mu,\nu}-2)j - (j+1)\beta}}{K_{1} e^{-i\gamma(\alpha_{\mu,\nu}-1)}} \right] \\ &< \exp \left[\frac{K_{2}}{K_{1}} e^{-\beta} \sum_{\infty}^{j-1} e^{(\gamma-\beta)i} \right] \\ &< \sqrt{A}. \end{split}$$

In a similar way we obtain

$$\frac{|(h_{\bar{\mu}}^{j})'(x)|}{|(h_{\bar{\mu}}^{j})'(\bar{\eta})|} > \frac{1}{\sqrt{A}},$$

and conclusion (1) of the lemma follows.

Integrating the first inequality in (3), we see that there is a constant \widetilde{K} , $0 < \widetilde{K} < 1$, such that $|h^{k+1}(\overline{\eta}) - \overline{\nu}| > \widetilde{K} |h^k(\overline{\eta})|^{\alpha_{\mu,\nu}}$. Now, applying the mean value theorem to h^j , conclusion (1) and the hyphotesis of lemma, we obtain

$$\left| h^{k+1+j}(\bar{\eta}) - h^j(\bar{\nu}) \right| > A^{-1} \lambda_1^j \widetilde{K} \mathrm{e}^{-r \alpha_{\mu,\nu}}, \quad j < k.$$

From the binding period definition, for $j \leq s$ and j < k we have $e^{-\beta\gamma} \geq A^{-1}\lambda_1^j \widetilde{K} \times e^{-r\alpha_{\mu,\nu}}$, which implies that

$$j \leqslant \frac{r\alpha_{\mu,\nu} - \log(KA^{-1})}{\beta + \log(\lambda_1)}$$

Next, choosing γ and β such that $2\gamma \alpha_{\mu,\nu} < \beta$ and fixing

$$k_1 > \frac{-2\log(\tilde{K}A^{-1})}{\beta + \log(\lambda_1) - 2\gamma\alpha_{\mu,\nu}}$$

then for j and k as in the lemma, we obtain

$$j \leqslant \frac{r\alpha_{\mu,\nu} - \log(KA^{-1})}{\beta + \log(\lambda_1)} < \frac{\gamma k\alpha_{\mu,\nu} - \log(KA^{-1})}{\beta + \log(\lambda_1)} < \frac{k}{2}.$$

If $s \ge k$ then the inequality above holds for j = k - 1 but this implies k - 1 < k/2, which is a contradiction. Hence, s < k and we can replace j = s in the inequality above, obtaining the right inequality

$$s \leqslant \frac{r\alpha_{\mu,\nu} - \log(\widetilde{K}A^{-1})}{\beta + \log(\lambda_1)}.$$

To establish the left inequality of conclusion (2) of the lemma, let

$$D = \max\{|h'_{\bar{\mu}}(x)|: p_n - \varepsilon \leq x \leq q_n + \varepsilon, \bar{\mu}\} > 1.$$

Applying the mean value theorem and the binding period definition once again, we obtain a point $\xi \in (0, h_{\bar{u}}^k(\bar{\eta}))$ such that:

$$D^{s+1}|h_{\bar{\mu}}^{k}(\bar{\eta})| \ge |(h_{\bar{\mu}}^{s+1})'(\xi)||h_{\bar{\mu}}^{k}(\bar{\eta}) - 0| = |h_{\bar{\mu}}^{s+1}(h_{\bar{\mu}}^{k}(\bar{\eta})) - h_{\bar{\mu}}^{s}(\bar{\nu})| > e^{-(s+1)\beta}$$

wich implies $-(s+1)\beta < (s+1)\log(D) - r$, and the left inequality follows.

In order to obtain conclusion (3), let put

$$t = \frac{\alpha_{\mu,\nu}}{\alpha_{\mu,\nu} - 1} = 1 + \frac{1}{\alpha_{\mu,\nu} - 1}$$

and take $\xi \in (\bar{\nu}, h^{k+1}_{\bar{\mu}}(\bar{\eta}))$ such that

$$(h_{\bar{\mu}}^{s})'(\xi) = \frac{|h_{\bar{\mu}}^{s+k+1}(\bar{\eta}) - h_{\bar{\mu}}^{s}(\bar{\nu})|}{|h_{\bar{\mu}}^{k+1}(\bar{\eta}) - \bar{\nu}|}.$$

Note that $|(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))| > A^{-1}|(h_{\bar{\mu}}^{s})'(\xi)|$ implies

$$|(h^{s}_{\bar{\mu}})'(h^{k+1}_{\bar{\mu}}(\bar{\eta}))| > A^{-1}|(h^{s}_{\bar{\mu}})'(\bar{\nu})| > A^{-1}\lambda_{1}^{s}$$

(since s < k). Then

$$\begin{split} |(h_{\bar{\mu}}^{s+1})'(h_{\bar{\mu}}^{k}(\bar{\eta}))|^{t} &= |(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))|^{t}|h_{\bar{\mu}}'(h_{\bar{\mu}}^{k}(\bar{\eta}))|^{t} \\ &\geqslant |(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))||(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))|^{1/(\alpha_{\mu,\nu}-1)}K_{1}^{t}|h_{\bar{\mu}}^{k}(\bar{\eta})|^{\alpha_{\mu,\nu}} \\ &\geqslant A^{-1}|(h_{\bar{\mu}}^{s})'(\xi)||(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))|^{1/(\alpha_{\mu,\nu}-1)}K_{1}^{t}\widetilde{K}|h_{\bar{\mu}}^{k+1}(\bar{\eta}) - \bar{\nu}| \\ &> A^{-1}K_{1}^{t}\widetilde{K}|h_{\bar{\mu}}^{s+k+1}(\bar{\eta}) - \bar{\nu}||(h_{\bar{\mu}}^{s})'(h_{\bar{\mu}}^{k+1}(\bar{\eta}))|^{1/(\alpha_{\mu,\nu}-1)} \\ &> A^{-1}K_{1}^{t}\widetilde{K}e^{-\beta(s+1)}A^{-1/(\alpha_{\mu,\nu}-1)}\lambda_{1}^{s/(\alpha_{\mu,\nu}-1)} \\ &= A^{-t}K_{1}^{t}\widetilde{K}e^{-\log\lambda_{1}/(\alpha_{\mu,\nu}-1)}\exp\left[\left(\frac{\log\lambda_{1}}{\alpha_{\mu,\nu}-1} - \beta\right)(s+1)\right] \\ &= \tau^{t}\exp\left[\left(\frac{\log\lambda_{1}}{\alpha_{\mu,\nu}-1} - \beta\right)(s+1)\right]. \end{split}$$

Finally, if β is small the coefficient of (s + 1) in the exponential above is positive, and the inequality in conclusion (3) of lemma follows for *s* large, which is obtained making δ small (compatible with the previous conditions imposed on α and β). \Box

A free return index is a return index that does not belong to a binding period associated to a previous return. Take γ , β and δ as in Lemma 5, and for any positive integer k let consider $i_1 < i_2 < \cdots < i_l \leq k$ all free return indices, i.e., $e^{-\gamma i_l} < |h_{\bar{\mu}}^{i_l}(\bar{\nu})| \leq \delta$, and let $s_1, s_2, \ldots, s_{l-1}$, (s_l) be the binding periods associated to returns $h_{\bar{\mu}}^{i_l}(\bar{\nu}), \ldots, h_{\bar{\mu}}^{i_{l-1}}(\bar{\nu}), h_{\bar{\mu}}^{i_l}(\bar{\nu})$. Note that s_l will be or will be not considered according to $i_l < i_l + s_l \leq k$ or $i_l \leq k < i_l + s_l$. We say that $j, j \leq k$ is a free time for the critical value $\bar{\nu}$ asociated to the parameter $(\bar{\mu}, \bar{\nu})$ if $|h_{\bar{\mu}}^j(\bar{\nu})| > \delta$ and $j \notin [i_t, i_t + s_t], t = 1, \ldots, l-1, (l)$. In the same way we define free time for the critical value $\bar{\mu}$ associated to the parameter $(\bar{\mu}, \bar{\nu})$.

Now, for $\bar{\eta} = \bar{\mu}$ or $\bar{\eta} = \bar{\nu}$ denote by $F(\bar{\eta}, k)$ the number given by:

 $F(\bar{\eta}, k) = \#\{j \leq k: j \text{ is a free time for the critical value}\}$

 $\bar{\eta}$ associated to the parameter $(\bar{\mu}, \bar{\nu})$.

Finally, fix $\varepsilon > 0$ and define $\tilde{E}_{n,\theta} = \tilde{E}_{n,\theta}(\gamma, \beta, \delta, \varepsilon)$ by

$$\tilde{E}_{n,\theta} = \left\{ (\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta} (\gamma) \colon \min \left\{ F(\bar{\mu}, k), F(\bar{\nu}, k) \right\} \ge (1 - \varepsilon)k, k \in \mathbb{N} \right\}$$

Indeed, $\tilde{E}_{n,\theta}$ is the set of parameter values $(\bar{\mu}, \bar{\nu}) \in L_{n,\theta}$ satisfying the basic assumption (BA) and, in addition, the number of free times for the critical values $\bar{\nu}, \bar{\mu}$ associated to the parameter $(\bar{\mu}, \bar{\nu})$ is "almost" k, in fact bigger than $(1 - \varepsilon)k$, for any $k \in \mathbb{N}$. Note that this last condition, called *free time assumption* (FA) implies $\sum_{t=1}^{l} s_t < \varepsilon k$.

As we have seen in the proof of Lemma 5, by taking δ small we can increase the binding periods and we can assume $s_t > K_0$, K_0 as in Proposition 2.

The set of free times is the union of the free intervals $V_1 = [1, ..., i_1 - 1], V_2 = [i_1 + s_1 + 1, ..., i_2 - 1], ..., V_l = [i_{l-1} + s_{l-1} + 1, i_l - 1], (V_{l+1} = [i_l + s_l + 1, k])$. Denote *B* as the set of indices *t* such that $\#V_t < K_0$ and *G* as the set of indices *t* such that $\#V_t \ge K_0$. We have

$$#B < \sum_{t \in B} #V_t < \sum_{t \in B} s_t < \sum_t s_t < \varepsilon k, \text{ and}$$

$$\sum_{t \in G} #V_t \ge (1 - 2\varepsilon)k.$$
(9)

PROPOSITION 3. – For γ , β , δ and ε small enough there are λ , $1 < \lambda < \lambda_0$ (λ_0 as in Proposition 2), $N \in \mathbb{N}$ and a > 0 such that: for $n \ge N$ and $|(\bar{\mu}, \bar{\nu}) - (\bar{\mu}_n, \bar{\nu}_n)| < a$, if $(\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta}$ then $|(h_{\bar{\mu}}^k)'(\bar{\mu})| > \lambda^k$ and $|(h_{\bar{\mu}}^k)'(\bar{\nu})| > \lambda^k$, k = 1, 2, ...

Proof. – To start with, we proceed by induction on k. Take $1 < \lambda < \lambda_0$ such that $|h'_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}}(p)| > \lambda$ (here p is the orientation preserving fixed point of $h_{\bar{\mu}_{\infty},\bar{\nu}_{\infty}}$). Because $h_{\bar{\mu}} = h_{n,\bar{\mu},\bar{\nu}}$ converges uniformly to $h_{\bar{\mu},\bar{\nu}}$ we can find $N \in \mathbb{N}$ and a > 0 such that for $n \ge N$ and $(\bar{\mu},\bar{\nu})$ with $|(\bar{\mu},\bar{\nu}) - (\bar{\mu}_n,\bar{\nu}_n)| < a$ then $|(h^k_{\bar{\mu}})'(\bar{\eta})| > \lambda^k$, $\bar{\eta} = \bar{\mu}, \bar{\nu}, k \le k_1$, where k_1 is taken as in Lemma 5. And then, the first step in the induction works for

 $k \le k_1$. Next, take $k > k_1$ and assume that the conclusion of the proposition is true for any iterate of $\bar{\eta} = \bar{\mu}$, $\bar{\nu}$ less than k. Because $k > k_1$ and $(\bar{\mu}, \bar{\nu})$ satisfies the basic assumption (BA), Lemma 5 works and it implies that during the binding periods there is not lost of derivative, in fact, the total derivative during those periods is bigger than 1. From Proposition 2, it follows that during the free time intervals of length less than K_0 $(t \in B)$ we have a decreasing in the derivative of at most C (C as in Proposition 2), whereas during the free time intervals of length bigger or equal than K_0 ($t \in G$) we have an increasing in the derivative of at least $\lambda_0^{\#V_t}$. Denote $C = e^{-u}$. We can write $k = \sum_{t \in B} \#V_t + \sum_{t \in G} \#V_t + \sum_t \sum_t s_t$ and from the discusion above and (9) we obtain

$$|(h_{\bar{\mu}}^{k})'(\bar{\eta})| \ge \prod_{t \in B} C \prod_{t \in G} \lambda_{0}^{\#V_{t}} \prod_{t=1}^{l} 1 > e^{-u \cdot \#B} \lambda_{0}^{\Sigma_{t \in G}} \#V_{t}$$
$$> e^{-u\varepsilon k} \lambda_{0}^{(1-2\varepsilon)k} = \exp([\log \lambda_{0}(1-2\varepsilon) - u\varepsilon]k).$$

Taking ε small we obtain $\log \lambda < \log \lambda_0 (1 - 2\varepsilon) - u\varepsilon$ (this choice works for any *k*) and then $|(h_{\bar{u}}^k)'(\bar{\eta})| \ge \lambda^k$. This ends the proof of the proposition. \Box

To finish up the proof of Theorem 2 we have to prove that it is possible to define a set $\tilde{E}_{n,\theta}$ as was mentioned before that satisfies:

- (1) $m_{\theta}(\tilde{E}_{n,\theta}) > 0$, where m_{θ} stands for the Lebesgue measure in $E_{n,\theta}$, and
- (2) For m_{θ} -almost every $(\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta}$, $\overline{\{h_{\bar{\mu}}^{k}(\bar{\eta})\}}_{k \in \mathbb{N}} = [h_{\bar{\mu}}(\bar{\xi}), \bar{\xi}]$, for $\bar{\eta} \in \{\bar{\mu}, \bar{\nu}\}$ where $\bar{\xi} = \max\{\bar{\mu}, \bar{\nu}\}$.

The set $\tilde{E}_{n,\theta}$ is constructed as in [2, Section 2]. This construction is nicely outlined in [10, Section 3], and here we refer to it. $\tilde{E}_{n,\theta}$ is obtained as the intersection of some sets E_k in a small parameter interval having $\bar{\mu}_n$ in the right extreme, $E_1 \supset E_2 \supset \cdots \supset$ $E_k \supset \cdots$, satisfying the basic assumption (BA) and the free time assumption (FA). In order to define these sets E_k , some partition \mathcal{P}_{k-1} of E_{k-1} is required. The crucial point here is that the derivative of the maps $\bar{\mu} \rightarrow h_{\bar{\mu}}^k(\bar{\mu})$ and $\bar{\mu} \rightarrow h_{\bar{\mu}}^k(\bar{\nu})$ (remember that $\bar{\nu} = \theta(\bar{\mu} - \bar{\mu}_n) + \bar{\nu}_n$), require to have bounded distorsion on each interval ω in \mathcal{P}_{k-1} ([10, Lemma 3.3] and [1, Lemma 5(7.2)]). This requirement and Lemma 5 enable us to estimate the Lebesgue measure of the set of points that will be excluded from E_{k-1} to define E_k , all which is gotten as outlined in [10] (also see [15]). Thus, just following the same steps as in [10, Section 3], we obtain a set $\tilde{E}_{n,\theta}$ satisfying the basic assumption (BA), the free time assumption (FA) and $m_{\theta}(\tilde{E}_{n,\theta}) > 0$.

Finally, in order to prove that the orbits of $\bar{\mu}$ and $\bar{\nu}$ are dense in $[h_{\bar{\mu}}(\bar{\xi}), \bar{\xi}]$ where $\bar{\xi} = \max\{\bar{\mu}, \bar{\nu}\}$ for almost every $(\bar{\mu}, \bar{\nu}) \in \tilde{E}_{n,\theta}$, we observe that the set of parameters such that the $h_{\bar{\mu}}$ -orbit of $\bar{\mu}$ do not visit some fixed open interval is a zero Lebesgue measure set because $\bar{\mu} \rightarrow h_{\bar{\mu}}^k(\bar{\mu})$ is a distortion bounded map for all k in the respective domain. The same conclusion holds for the critical value $\bar{\nu}$. Finally, using the fact that the topology of $L_{n,\theta}$ has a countable basis we can conclude that the set of parameters such that the critical values are not dense is a zero Lebesgue measure set. The proof of Theorem 2 is now complete.

REFERENCES

- [1] M. Benedicks, L. Carleson, On iterations of $1 ax^2$ on (-1, 1), Ann. Math. 122 (1985) 1–24.
- [2] M. Benedicks, L. Carleson, The dynamics of the Hénon map, Ann. Math. 133 (1991) 73– 169.
- [3] R. Bamón, R. Labarca, R. Mañé, M.J. Pacífico, The explosion of singular cycles, Publ. Math. IHES 78 (1993) 207–232.
- [4] P. Collet, J.P. Eckmann, On the abundance of aperiodic behaviour for maps on the interval, Comm. Phys. 73 (1980) 115–160.
- [5] J. Guckenheimer, Sensitive dependence to initial conditions for one dynamical maps, Comm. Math. Phys. 70 (1979) 133–160.
- [6] R. Labarca, Bifurcation of contracting singular cycles, Ann. Scient. Ec. Norm. Sup. 4e Série 28 (1995) 705–745.
- [7] R. Labarca, M.J. Pacífico, Stability of singular horseshoe, Topology 25 (1986) 337-352.
- [8] W. de Mello, S. van Strien, One-Dimensional Dynamics, in: Modern Surveys in Mathematics, Vol. 25, Springer-Verlag, Berlin, 1991.
- [9] C. Morales, M.J. Pacifico, B. San Martín, Lorenz atractor from double homoclinic resonant cycles, Preprint.
- [10] L. Mora, M. Viana, Abundance of strange attractors, Acta Math. 171 (1993) 1-71.
- [11] M.J. Pacífico, A. Rovella, Unfolding contracting singular cycles, Ann. Scient. Éc. Norm. Sup. 4e Série 26 (1993) 691–700.
- [12] J. Palis, On Morse–Smale dynamical systems, Topology 8 (1969).
- [13] J. Palis, S. Smale, Structural Stability Theorems, in: Global Analysis, Proc. Symp. in Pure Math., Vol. XIV, American Mathematical Society, Providence, RI, 1970.
- [14] J. Palis, F. Takens, Hyperbolicity & Sensitive Chaotic Dynamics at Homoclinic Bifurcations, in: Cambridge Studies in Advanced Mathematics, Vol. 35, Cambridge University Press, Cambridge, 1993.
- [15] A. Rovella, The dynamics of perturbations of the contracting Lorenz attractor, Bol. Soc. Bras. Mat. 24 (2) 233–259.
- [16] B. San Martín, Contracting singular cycles, Ann. Inst. Henri Poincaré 15 (5) (1998) 651– 659.
- [17] B. San Martín, J. Vera, Nonhyperbolic persistent attractors near the Morse–Smale boundary II, Informe de matemáticas UCN.
- [18] M. Shub, Global Stability of Dynamical Systems, Springer-Verlag, Berlin, 1987.