# Regularity in a one-phase free boundary problem for the fractional Laplacian 

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#### Abstract

For a one-phase free boundary problem involving a fractional Laplacian, we prove that "flat free boundaries" are $C^{1, \alpha}$. We recover the regularity results of Caffarelli for viscosity solutions of the classical Bernoulli-type free boundary problem with the standard Laplacian.


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## 1. Introduction

The purpose of this paper is to answer a question left open in [7] on the regularity of free boundaries for the fractional Laplacian of order $\alpha$ - with $0<\alpha<1$, in the particular case $\alpha=1 / 2$. Here is the setting: consider $g$ a viscosity solution (this notion will be defined properly later) of the following free boundary problem in the ball $B_{1} \subset \mathbb{R}^{n+1}=\mathbb{R}^{n} \times \mathbb{R}$,

$$
\begin{cases}\Delta g=0, & \text { in } B_{1}^{+}(g):=B_{1} \backslash\{(x, 0): g(x, 0)=0\},  \tag{1.1}\\ \frac{\partial g}{\partial U}=1, & \text { on } F(g):=\partial_{\mathbb{R}^{n}}\left\{x \in \mathcal{B}_{1}: g(x, 0)>0\right\} \cap \mathcal{B}_{1},\end{cases}
$$

where

$$
\begin{equation*}
\frac{\partial g}{\partial U}\left(x_{0}\right):=\lim _{(t, z) \rightarrow(0,0)} \frac{g\left(x_{0}+t v\left(x_{0}\right), z\right)}{U(t, z)}, \quad x_{0} \in F(g) \tag{1.2}
\end{equation*}
$$

and $\mathcal{B}_{r} \subset \mathbb{R}^{n}$ is the $n$-dimensional ball of radius $r$ (centered at 0 ).
The function $U(t, z)$ is the harmonic extension of $\sqrt{t^{+}}$to the upper half-plane $\mathbb{R}_{+}^{2}=\{(t, z) \in \mathbb{R} \times \mathbb{R}, z>0\}$, reflected evenly across $\{z=0\}$. Precisely, after the polar change of coordinates

$$
t=r \cos \theta, \quad z=r \sin \theta, \quad r \geqslant 0,-\pi \leqslant \theta \leqslant \pi,
$$

[^0]$U$ is given by
\[

$$
\begin{equation*}
U(t, z)=r^{1 / 2} \cos \frac{\theta}{2} . \tag{1.3}
\end{equation*}
$$

\]

One can show that if a function $g \geqslant 0$ is harmonic in $B_{1}^{+}(g)$ and $F(g)$ is smooth around a point $x_{0}$ then $\frac{\partial g}{\partial U}\left(x_{0}\right)$ exists always and is finite. Here $v\left(x_{0}\right)$ denotes as usually the normal to $F(g)$ at $x_{0}$ pointing toward $\{x: g(x, 0)>0\}$.

In this paper, we introduce the notion of viscosity solutions to (1.1) and prove the following result about the regularity of their free boundaries under appropriate flatness assumptions (for all the relevant definitions see Section 2).

Theorem 1.1. There exists a universal constant $\bar{\epsilon}>0$, such that if $g$ is a viscosity solution to (1.1) satisfying

$$
\begin{equation*}
\|g-U\|_{L^{\infty}\left(\bar{B}_{1}\right)} \leqslant \bar{\epsilon}, \tag{1.4}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\{x \in \mathcal{B}_{1}: x_{n} \leqslant-\bar{\epsilon}\right\} \subset\left\{x \in \mathcal{B}_{1}: g(x, 0)=0\right\} \subset\left\{x \in \mathcal{B}_{1}: x_{n} \leqslant \bar{\epsilon}\right\}, \tag{1.5}
\end{equation*}
$$

then $F(g)$ is $C^{1, \alpha}$ in $\mathcal{B}_{1 / 2}$.
Consequently $\frac{\partial g}{\partial U}$ exists and $g$ is a classical solution to (1.1). Moreover, given a point $x_{0}$ on the free boundary $F(g)$ if one knows that a blow-up sequence of $g$ around $x_{0}$ "converges" to the function $U$, then the flatness assumptions (1.4)-(1.5) are satisfied and hence the free boundary is $C^{1, \alpha}$ around that point.

Assumption (1.4) is a (slightly improved) nondegeneracy assumption which is usually true, and certainly satisfied in the framework of [7]. In any case it could be removed, but we keep it for simplicity.

The interest in our free boundary problem (1.1) arises from a natural generalization of the following classical Bernoulli-type one-phase free boundary problem:

$$
\begin{cases}\Delta u=0, & \text { in } \Omega \cap\{u>0\},  \tag{1.6}\\ |\nabla u|=1, & \text { on } \Omega \cap \partial\{u>0\},\end{cases}
$$

with $\Omega$ a domain in $\mathbb{R}^{n}$. A pioneering investigation was that of Alt and Caffarelli [1] (variational context), and then Caffarelli [2-4] (viscosity solutions context). See also [8] for a complete survey.

A special class of viscosity solutions to (1.1) (with the constant 1 replaced by a precise constant $A$ ) is provided by minimizers to the energy functional

$$
J\left(v, B_{1}\right)=\int_{B_{1}}|\nabla v|^{2} d x d z+\mathcal{L}_{\mathbb{R}^{n}}\left(\{v>0\} \cap \mathbb{R}^{n} \cap B_{1}\right) .
$$

Such minimizers have been investigated by Caffarelli, Sire and the second author in [7], where general properties (optimal regularity, nondegeneracy, classification of global solutions), corresponding to those proved by Alt and Caffarelli in [1] for the Bernoulli-type problem (1.6), have been obtained.

As for the next issue, i.e. the regularity of the free boundary, here is what is proved in [7] in the setting of (1.1):
Let $u(x, y, z)$ be a solution of (1.1) in $B_{1} \subset \mathbb{R}^{3}$. Assume that the free boundary of $u$ is a Lipschitz graph in $\mathcal{B}_{1}$. Then it is a $C^{1}$ graph in $\mathcal{B}_{1 / 2}$.

The idea of this result is that (i) one can find two points on each side of 0 where the free boundary is flatter than what is dictated by the Lipschitz constant, (ii) this improvement could be propagated inside a small ball of controlled size. Thus the three-dimensionality of the problem (or, equivalently, the one-dimensionality of the free boundary) is heavily used. Moreover, this argument does not yield the extra Hölder regularity of the derivative - which we believe could itself yield $C^{\infty}$ regularity of the free boundary. What we propose in this paper is to fill the gap between $C^{1}$ and $C^{1, \alpha}$, in arbitrary space dimension.

In view of the results in [7], one knows that the flatness assumptions (1.4), (1.5) in our main Theorem 1.1 are satisfied around each point of the reduced part of the free boundary of a minimizer (see Propositions 4.2 and Theorems 1.2, 1.3 in [7]). We thus obtain the following corollary to Theorem 1.1.

Corollary 1.2. Let $v$ be a local minimizer to

$$
J\left(v, B_{1}\right)=\int_{B_{1}}|\nabla v|^{2} d x d z+\mathcal{L}_{\mathbb{R}^{n}}\left(\{v>0\} \cap \mathbb{R}^{n} \cap B_{1}\right)
$$

Then the reduced part of the free boundary $F^{*}(v)$ is $C^{1, \alpha}$.
Let us now recall how the fractional Laplacian is involved in (1.1). Consider, for $\alpha \in(0,1)$, the model (which generalizes (1.6))

$$
\begin{cases}(-\Delta)^{\alpha} u=0, & \text { in } \Omega \cap\{u>0\},  \tag{1.7}\\ \lim _{t \rightarrow 0^{+}} \frac{u\left(x_{0}+t v\left(x_{0}\right)\right)}{t^{\alpha}}=\text { const., } & \text { on } \Omega \cap \partial\{u>0\},\end{cases}
$$

with $u$ defined on the whole $\mathbb{R}^{n}$ with prescribed values outside of $\Omega$. Recall that, up to a normalization constant

$$
(-\Delta)^{\alpha} u(x)=P V \int_{\mathbb{R}^{n}} \frac{u(x)-u(y)}{|x-y|^{n+2 \alpha}} d y
$$

where $P V$ denotes the Cauchy principal value.
When studying local property of the free boundary in (1.7), the non-locality of the fractional Laplace makes the problem quite delicate. To avoid this "contrast" one can make use of an extension property proved by Caffarelli and Silvestre in [10] (see for example the work of Caffarelli, Savin and the second author [6], the paper [7], and the work of Caffarelli, Salsa and Silvestre [9] where this strategy has been employed). Precisely, let $u \in C^{2}\left(\mathbb{R}^{n}\right)$ and let $v$ solve

$$
\begin{cases}-\operatorname{div}\left(z^{\beta} \nabla v\right)=0, & \text { in } \mathbb{R}_{+}^{n+1}=\left\{(x, z) \in \mathbb{R}^{n} \times \mathbb{R}, z>0\right\}  \tag{1.8}\\ v(x, 0)=u(x), & \text { on } \mathbb{R}^{n},\end{cases}
$$

with $\beta=1-2 \alpha$. Then,

$$
\begin{equation*}
\left(-\Delta^{\alpha}\right) u(x)=-\lim _{z \rightarrow 0}\left(z^{\beta} v_{z}(x, z)\right) . \tag{1.9}
\end{equation*}
$$

After extending $v$ evenly across the hyperplane $\{z=0\}$, the first equation in (1.8) can be thought in the whole $\mathbb{R}^{n+1}$. In view of this formula, the focus shifts on the free boundary problem,

$$
\begin{cases}-\operatorname{div}\left(|z|^{\beta} \nabla v\right)=0, & \text { in } B_{1} \backslash\{(x, 0): v(x, 0)=0\},  \tag{1.10}\\ \lim _{(t, z) \rightarrow(0,0)} \frac{v\left(x_{0}+t v\left(x_{0}\right), z\right)}{U(t, z)}=\text { const., } & \text { on } \partial_{\mathbb{R}^{n}}\{x: v(x, 0)>0\} \cap \mathcal{B}_{1},\end{cases}
$$

where $U(t, z)$ solves (1.8) in $\mathbb{R}^{2}$ with $u(t)=\left(t^{+}\right)^{\alpha}$ and is extended evenly across $\{z=0\}$.
For simplicity of exposition we have focused here on the case when $\alpha=1 / 2$ (in which case the extension formula of [10] is a well-know fact). However our result can probably be extended to the general case $\alpha \in(0,1)$.

Our definition of viscosity solution to (1.1) is similar to the one introduced by Caffarelli in [2,3] to deal with the problem (1.6). Indeed our result generalizes to this non-local setting the "flatness implies regularity" theory developed by Caffarelli in [3]. Let us also mention that Theorem 1.1 is probably optimal. Indeed, quite similarly to what happens for minimal surfaces, singular free boundaries for the Bernoulli-type problem (1.6) were discovered by Jerison and the first author [11].

Let us now describe our strategy to obtain Theorem 1.1. The main idea to prove Theorem 1.1 is to show that $F(g)$ enjoys an "improvement of flatness" property, that is if $F(g)$ oscillates $\epsilon$ away from a hyperplane in $\mathcal{B}_{1}$ ( $\epsilon$ small), then in $\mathcal{B}_{\rho}$ ( $\rho$ universal) it oscillates $\epsilon \rho / 2$ away from possibly a different hyperplane. To obtain this improvement of flatness, we use a compactness argument which goes as follows: assume one cannot do it however flat the free boundary is, then we blow it up in the $x_{n}$-direction, thus linearizing the problem into a limiting one, for which we prove that improvement of flatness holds - thus a contradiction. This scheme was used by Savin [13] to prove regularity of small solutions of fully nonlinear equations - including an elegant proof of the De Giorgi theorem for minimal surfaces. The key tool is a geometric Harnack inequality that localizes the free boundary well, and allows the passage to the limit under rescalings.

Such compactness arguments can also be found in Wang [14] for the regularity of the solutions of $p$-Laplace equations. More recently, the first author followed this strategy in [12] to provide a new proof of the Caffarelli "flat implies smooth" theory. The scheme is the same here, up to the fact that the construction of the sub-solution opening the way to the Harnack inequality is different from that of [12], and that the linear problem obtained eventually is non-standard (and interesting in its own).

The paper is organized as follows. In Section 2 we introduce the notion of viscosity solutions to (1.1) and we prove a basic comparison principle for such solutions. In Section 3 we explain how to interpret our solutions as perturbations of $U$ after a "domain variation" in the $e_{n}$-direction and we present basic facts about such domain variations. Throughout the paper, this will be a convenient way of thinking about our viscosity solutions. In Section 4 we describe the linear problem associated to (1.1) and later in Section 8 we obtain a regularity result for its solutions. Section 5 contains some technical lemmas leading to the proof of Harnack inequality. In Section 6 we exhibit the proof of Harnack inequality using the barrier which we will construct later in Appendix A. Finally in Section 7 we provide the proof of the "improvement of flatness" property.

## 2. Definitions and basic lemmas

In this section we introduce notation and definitions which we will use throughout the paper and we prove a standard basic lemma (Comparison Principle).

A point $X \in \mathbb{R}^{n+1}$ will be denoted by $X=(x, z) \in \mathbb{R}^{n} \times \mathbb{R}$. We will also use the notation $x=\left(x^{\prime}, x_{n}\right)$ with $x^{\prime}=\left(x_{1}, \ldots, x_{n-1}\right)$. A ball in $\mathbb{R}^{n+1}$ with radius $r$ and center $X$ is denoted by $B_{r}(X)$ and for simplicity $B_{r}=B_{r}(0)$. Also we use $\mathcal{B}_{r}$ to denote the $n$-dimensional ball $B_{r} \cap\{z=0\}$.

Let $v(X)$ be a continuous non-negative function in $B_{1}$. We associate to $v$ the following sets:

$$
\begin{aligned}
& B_{1}^{+}(v):=B_{1} \backslash\{(x, 0): v(x, 0)=0\} \subset \mathbb{R}^{n+1} ; \\
& \mathcal{B}_{1}^{+}(v):=B_{1}^{+}(v) \cap \mathcal{B}_{1} \subset \mathbb{R}^{n} ; \\
& F(v):=\partial_{\mathbb{R}^{n}} \mathcal{B}_{1}^{+}(v) \cap \mathcal{B}_{1} \subset \mathbb{R}^{n} ; \\
& \mathcal{B}_{1}^{0}(v):=\operatorname{Int}_{\mathbb{R}^{n}}\left\{x \in \mathbb{R}^{n}: v(x, 0)=0\right\} \subset \mathbb{R}^{n} .
\end{aligned}
$$

Often subsets of $\mathbb{R}^{n}$ are embedded in $\mathbb{R}^{n+1}$, as it will be clear from the context.
We may refer to $\mathcal{B}_{1}^{0}(v)$ as to the zero plate of $v$, while $F(v)$ is called the free boundary of $v$.
We consider the free boundary problem

$$
\begin{cases}\Delta g=0, & \text { in } B_{1}^{+}(g),  \tag{2.1}\\ \frac{\partial g}{\partial U}=1, & \text { on } F(g),\end{cases}
$$

where

$$
\frac{\partial g}{\partial U}\left(x_{0}\right):=\lim _{(t, z) \rightarrow(0,0)} \frac{g\left(x_{0}+t v\left(x_{0}\right), z\right)}{U(t, z)}, \quad X_{0}=\left(x_{0}, 0\right) \in F(g) .
$$

Here $v\left(x_{0}\right)$ denotes the unit normal to $F(g)$ at $x_{0}$ pointing toward $\mathcal{B}_{1}^{+}(g)$ and $U$ is the function defined in (1.3). Also, throughout the paper we call $U(X):=U\left(x_{n}, z\right)$.

We now introduce the notion of viscosity solutions to (2.1). First we need the following standard notion.
Definition 2.1. Given $g, v$ continuous, we say that $v$ touches $g$ by below (resp. above) at $X_{0} \in B_{1}$ if $g\left(X_{0}\right)=v\left(X_{0}\right)$, and

$$
g(X) \geqslant v(X) \quad(\text { resp. } g(X) \leqslant v(X)) \quad \text { in a neighborhood } O \text { of } X_{0} .
$$

If this inequality is strict in $O \backslash\left\{X_{0}\right\}$, we say that $v$ touches $g$ strictly by below (resp. above).
Definition 2.2. We say that $v \in C^{2}\left(B_{1}\right)$ is a (strict) comparison subsolution to (2.1) if $v$ is a non-negative function in $B_{1}$ which is even with respect to $\{z=0\}$ and it satisfies
(i) $\Delta v \geqslant 0$ in $B_{1}^{+}(v)$;
(ii) $F(v)$ is $C^{2}$ and if $x_{0} \in F(v)$ we have

$$
v(x, z)=\alpha U\left(\left(x-x_{0}\right) \cdot v\left(x_{0}\right), z\right)+o\left(\left|\left(x-x_{0}, z\right)\right|^{1 / 2}\right), \quad \text { as }(x, z) \rightarrow\left(x_{0}, 0\right),
$$

with

$$
\alpha \geqslant 1
$$

where $v\left(x_{0}\right)$ denotes the unit normal at $x_{0}$ to $F(v)$ pointing toward $\mathcal{B}_{1}^{+}(v)$;
(iii) Either $v$ is not harmonic in $B_{1}^{+}(v)$ or $\alpha>1$.

Similarly one can define a (strict) comparison supersolution.
Definition 2.3. We say that $g$ is a viscosity solution to (2.1) if $g$ is a continuous non-negative function in $B_{1}$ which is even with respect to $\{z=0\}$ and it satisfies
(i) $\Delta g=0$ in $B_{1}^{+}(g)$;
(ii) Any (strict) comparison subsolution (resp. supersolution) cannot touch $g$ by below (resp. by above) at a point $X_{0}=\left(x_{0}, 0\right) \in F(g)$.

Remark 2.4. By standard arguments, if $g$ is a viscosity solution to (2.1) and $F(g)$ is $C^{1}$ then $g$ is a classical solution of the free boundary problem (see for example Proposition 4.2 in [7]). Moreover, as remarked in the Introduction one can show that given any continuous function $g$ which is harmonic in $B_{1}^{+}(g)$, then $\frac{\partial g}{\partial U}\left(x_{0}\right)$ exists at each point around which $F(g)$ is $C^{1, \alpha}$. These facts motivate our problem and the definition of viscosity solution.

Remark 2.5. We remark that if $g$ is a viscosity solution to (2.1) in $B_{\rho}$, then

$$
g_{\rho}(X)=\rho^{-1 / 2} g(\rho X), \quad X \in B_{1}
$$

is a viscosity solution to (2.1) in $B_{1}$.
We finish this section by stating and proving a comparison principle for problem (2.1) which will be a key tool in the proof of Harnack inequality in Section 7.

Lemma 2.6 (Comparison Principle). Let $g, v_{t} \in C\left(\bar{B}_{1}\right)$ be respectively a solution and a family of subsolutions to (2.1), $t \in[0,1]$. Assume that
(i) $v_{0} \leqslant g$, in $\bar{B}_{1}$;
(ii) $v_{t} \leqslant g$ on $\partial B_{1}$ for all $t \in[0,1]$;
(iii) $v_{t}<g$ on $\mathcal{F}\left(v_{t}\right)$ which is the boundary in $\partial B_{1}$ of the set $\partial \mathcal{B}_{1}^{+}\left(v_{t}\right) \cap \partial \mathcal{B}_{1}$, for all $t \in[0,1]$;
(iv) $v_{t}(x)$ is continuous in $(x, t) \in \bar{B}_{1} \times[0,1]$ and $\overline{\mathcal{B}_{1}^{+}\left(v_{t}\right)}$ is continuous in the Hausdorff metric.

Then

$$
v_{t} \leqslant g \quad \text { in } \bar{B}_{1}, \text { for all } t \in[0,1] .
$$

Proof. Let

$$
A:=\left\{t \in[0,1]: v_{t}(x) \leqslant g(x) \text { on } \bar{B}_{1}\right\} .
$$

In view of (i) and (iv) $A$ is closed and non-empty. Our claim will follow if we show that $A$ is open. Let $t_{0} \in A$, then $v_{t_{0}} \leqslant g$ on $\bar{B}_{1}$ and by the definition of viscosity solution

$$
F\left(v_{t_{0}}\right) \cap F(g)=\emptyset .
$$

Together with (iii) this implies that

$$
\mathcal{B}_{1}^{+}\left(v_{t_{0}}\right) \subset \mathcal{B}_{1}^{+}(g), \quad F\left(v_{t_{0}}\right) \cup \mathcal{F}\left(v_{t_{0}}\right) \subset\left\{x \in \overline{\mathcal{B}}_{1}: g(x, 0)>0\right\} .
$$

By (iv) this gives that for $t$ close to $t_{0}$

$$
\begin{equation*}
\mathcal{B}_{1}^{+}\left(v_{t}\right) \subset \mathcal{B}_{1}^{+}(g), \quad F\left(v_{t}\right) \cup \mathcal{F}\left(v_{t}\right) \subset\left\{x \in \overline{\mathcal{B}}_{1}: g(x, 0)>0\right\} . \tag{2.2}
\end{equation*}
$$

Call $D:=B_{1} \backslash\left(\mathcal{B}_{1}^{0}\left(v_{t}\right) \cup F\left(v_{t}\right)\right)$. Combining (2.2) with assumption (ii) we get that

$$
v_{t} \leqslant g \quad \text { on } \partial D,
$$

and by the maximum principle the inequality holds also in $D$. Hence

$$
v_{t} \leqslant g \quad \text { in } \bar{B}_{1},
$$

and $t \in A$ which shows that $A$ is open.
Corollary 2.7. Let $g$ be a solution to (2.1) and let $v$ be a subsolution to (2.1) in $B_{2}$ which is strictly monotone increasing in the $e_{n}$-direction in $B_{2}^{+}(v)$. Call

$$
v_{t}(X):=v\left(X+t e_{n}\right), \quad X \in B_{1} .
$$

Assume that for $-1 \leqslant t_{0}<t_{1} \leqslant 1$

$$
v_{t_{0}} \leqslant g, \quad \text { in } \bar{B}_{1},
$$

and

$$
v_{t_{1}} \leqslant g \quad \text { on } \partial B_{1}, \quad v_{t_{1}}<g \quad \text { on } \mathcal{F}\left(v_{t_{1}}\right) .
$$

Then

$$
v_{t_{1}} \leqslant g \quad \text { in } \bar{B}_{1} .
$$

## 3. The function $\tilde{g}$

Let $g$ be a viscosity solution to (2.1). Throughout the paper, it will be convenient to interpret $g$ as a perturbation of $U$ via a domain variation in the $e_{n}$-direction. In this section we explain some basic facts about such domain variations.

Let $\epsilon>0$ and let $g$ be a continuous non-negative function in $\bar{B}_{\rho}$. Here and henceforth we denote by $P$ and $L$ the half-hyperplanes $P:=\left\{X \in \mathbb{R}^{n+1}: x_{n} \leqslant 0, z=0\right\}$ and $L:=\left\{X \in \mathbb{R}^{n+1}: x_{n}=0, z=0\right\}$. To each $X \in \mathbb{R}^{n+1} \backslash P$ we associate $\tilde{g}_{\epsilon}(X) \subset \mathbb{R}$ via the formula

$$
\begin{equation*}
U(X)=g\left(X-\epsilon w e_{n}\right), \quad \forall w \in \tilde{g}_{\epsilon}(X) . \tag{3.1}
\end{equation*}
$$

We sometimes call $\tilde{g}_{\epsilon}$ the $\epsilon$-domain variation associated to $g$. By abuse of notation, from now on we write $\tilde{g}_{\epsilon}(X)$ to denote any of the values in this set.

If $g$ satisfies

$$
\begin{equation*}
U\left(X-\epsilon e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon e_{n}\right) \quad \text { in } B_{\rho}, \tag{3.2}
\end{equation*}
$$

then

$$
\tilde{g}_{\epsilon}(X) \in[-1,1] .
$$

Indeed call

$$
Y=X-\epsilon \tilde{g}_{\epsilon}(X) e_{n}, \quad X \in \mathbb{R}^{n+1} \backslash P .
$$

Then according to (3.2),

$$
U\left(Y-\epsilon e_{n}\right) \leqslant g(Y)=U\left(Y+\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right) \leqslant U\left(Y+\epsilon e_{n}\right) .
$$

Since $U\left(Y+\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right)=U(X)>0$ our claim follows from the strict monotonicity of $U$ in the $e_{n}$-direction (outside of $P$ ).

Moreover, under the assumption (3.2) for each $X \in B_{\rho-\epsilon} \backslash P$ there exists at least one value $\tilde{g}_{\epsilon}(X)$ such that

$$
\begin{equation*}
U(X)=g\left(X-\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right) \tag{3.3}
\end{equation*}
$$

Indeed, it follows from (3.2) that

$$
g\left(X-\epsilon e_{n}\right) \leqslant U(X) \leqslant g\left(X+\epsilon e_{n}\right), \quad X \in B_{\rho-\epsilon}
$$

and our claim follows by the continuity of $g\left(X-\delta \epsilon e_{n}\right), \delta \in[-1,1]$.
Thus if (3.2) holds, for all $\epsilon>0$ we can associate to $g$ a possibly multi-valued function $\tilde{g}_{\epsilon}$ defined at least on $B_{\rho-\epsilon} \backslash P$ and taking values in $[-1,1]$ which satisfies (3.3). Moreover if $g$ is strictly monotone in the $e_{n}$-direction in $B_{\rho}^{+}(g)$, then $\tilde{g}_{\epsilon}$ is single-valued.

The following elementary lemma will be used to obtain a useful comparison principle for the $\epsilon$-domain variations of solutions to (2.1).

Lemma 3.1. Let $g$, $v$ be non-negative continuous functions in $B_{\rho}$. Assume that $g$ satisfies the flatness condition (3.2) in $B_{\rho}$ and that $v$ is strictly increasing in the $e_{n}$-direction in $B_{\rho}^{+}(v)$. Then if

$$
v \leqslant g \quad \text { in } B_{\rho},
$$

and $\tilde{v}_{\epsilon}$ is defined on $B_{\rho-\epsilon} \backslash P$ we have that

$$
\tilde{v}_{\epsilon} \leqslant \tilde{g}_{\epsilon} \quad \text { on } B_{\rho-\epsilon} \backslash P .
$$

Vice versa, if $\tilde{v}_{\epsilon}$ is defined on $B_{s} \backslash P$ and

$$
\tilde{v}_{\epsilon} \leqslant \tilde{g}_{\epsilon} \quad \text { on } B_{s} \backslash P,
$$

then

$$
v \leqslant g \quad \text { on } B_{s-\epsilon} .
$$

Proof. The first implication is obvious. Indeed, assume by contradiction that $v \leqslant g$ in $B_{\rho}$ and there exists $X \in$ $B_{\rho-\epsilon} \backslash P$ such that

$$
\tilde{v}_{\epsilon}(X)>\tilde{g}_{\epsilon}(X) .
$$

By the strict monotonicity of $v$ in the $e_{n}$-direction in $B_{\rho}^{+}(v)$ we have that

$$
0<U(X)=g\left(X-\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right)=v\left(X-\epsilon \tilde{v}_{\epsilon}(X) e_{n}\right)<v\left(X-\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right) .
$$

Thus there exists $Y=X-\epsilon \tilde{g}_{\epsilon}(X) e_{n} \in B_{\rho}$ such that $g(Y)<v(Y)$, a contradiction. Vice versa, suppose that $\tilde{v}_{\epsilon} \leqslant \tilde{g}_{\epsilon}$ in $B_{s} \backslash P$. For a fixed $Y \in B_{s-\epsilon}$ we know by the flatness assumption (3.2) that

$$
U\left(Y-\epsilon e_{n}\right) \leqslant g(Y) \leqslant U\left(Y+\epsilon e_{n}\right) .
$$

Thus, there exists $X \in B_{s}$ with $x_{i}=y_{i}$ for $i \neq n$ and $x_{n} \in\left[y_{n}-\epsilon e_{n}, y_{n}+\epsilon e_{n}\right]$ such that

$$
g(Y)=U(X)
$$

Suppose $g(Y) \neq 0$, then the identity above means that one of the possible values of $\tilde{g}_{\epsilon}(X)=\frac{y_{n}-x_{n}}{\epsilon}$. Again, using that $v$ is increasing in the $e_{n}$-direction we get:

$$
g(Y)=U(X)=v\left(X-\epsilon \tilde{v}_{\epsilon}(X) e_{n}\right) \geqslant v\left(X-\epsilon \tilde{g}_{\epsilon}(X) e_{n}\right)=v(Y), \quad Y \in B_{s}^{+}(g) .
$$

Thus the desired inequality holds in $B_{s}^{+}(g)$ and hence by continuity it holds in the full ball $B_{s}$.
We now state and prove the desired comparison principle, which will follow immediately from the lemma above and Corollary 2.7.

Lemma 3.2. Let $g, v$ be respectively a solution and a subsolution to (2.1) in $B_{2}$, with $v$ strictly increasing in the $e_{n}$-direction in $B_{2}^{+}(v)$. Assume that $g$ satisfies the flatness assumption (3.2) in $B_{2}$ for $\epsilon>0$ small and that $\tilde{v}_{\epsilon}$ is defined in $B_{2-\epsilon} \backslash P$ and satisfies

$$
\left|\tilde{v}_{\epsilon}\right| \leqslant C .
$$

If,

$$
\begin{equation*}
\tilde{v}_{\epsilon}+c \leqslant \tilde{g}_{\epsilon} \quad \text { in }\left(B_{3 / 2} \backslash \bar{B}_{1 / 2}\right) \backslash P, \tag{3.4}
\end{equation*}
$$

then

$$
\begin{equation*}
\tilde{v}_{\epsilon}+c \leqslant \tilde{g}_{\epsilon} \quad \text { in } B_{3 / 2} \backslash P \tag{3.5}
\end{equation*}
$$

Proof. We wish to apply Corollary 2.7 to the functions $g$ and

$$
v_{\epsilon t}=v\left(X+\epsilon t e_{n}\right) .
$$

We need to verify that for some $t_{0}<t_{1}=c$

$$
\begin{equation*}
v_{\epsilon t_{0}} \leqslant g \quad \text { in } \bar{B}_{1}, \tag{3.6}
\end{equation*}
$$

and for all $\delta>0$ and small

$$
\begin{equation*}
v_{\epsilon t_{1}} \leqslant g \quad \text { on } \partial B_{1}, \quad v_{\epsilon t_{1}}<g \quad \text { on } \mathcal{F}\left(v_{\epsilon\left(t_{1}-\delta\right)}\right) . \tag{3.7}
\end{equation*}
$$

Then our corollary implies

$$
v_{\epsilon\left(t_{1}-\delta\right)} \leqslant g \quad \text { in } \bar{B}_{1} .
$$

By letting $\delta$ go to 0 , we obtain that

$$
v_{\epsilon t_{1}} \leqslant g \quad \text { in } \bar{B}_{1},
$$

which in view of Lemma 3.1 gives

$$
\widetilde{\left(v_{\epsilon t_{1}}\right)_{\epsilon}} \leqslant \tilde{g}_{\epsilon} \quad \text { in } B_{1-\epsilon} \backslash P,
$$

assuming that the $\epsilon$-domain variation on the left-hand side exists on $B_{1-\epsilon} \backslash P$. On the other hand, it is easy to verify that on such set

$$
\begin{equation*}
\left(\widetilde{v_{\epsilon t}}\right)_{\epsilon}(X)=\tilde{v}_{\epsilon}(X)+t \tag{3.8}
\end{equation*}
$$

and hence we have

$$
\tilde{v}_{\epsilon}+c=\tilde{v}_{\epsilon}+t_{1} \leqslant \tilde{g}_{\epsilon} \quad \text { in } B_{1-\epsilon} \backslash P
$$

which gives the desired conclusion. We are left with the proof of (3.6)-(3.7).
In view of Lemma 3.1, in order to obtain (3.6) it suffices to show that

$$
\widetilde{\left(v_{\epsilon t_{0}}\right)_{\epsilon}} \leqslant \tilde{g}_{\epsilon}, \quad \text { in } B_{1+\epsilon} \backslash P
$$

which by (3.8) becomes

$$
\tilde{v}_{\epsilon}+t_{0} \leqslant \tilde{g}_{\epsilon}, \quad \text { in } B_{1+\epsilon} \backslash P .
$$

This last inequality holds trivially since $\tilde{g}_{\epsilon}$ and $\tilde{v}_{\epsilon}$ are bounded.
For (3.7), notice that the first inequality follows easily from our assumption (3.4) together with (3.8) and Lemma 3.1. More precisely we have that

$$
v_{\epsilon t_{1}} \leqslant g \quad \text { in } B_{\frac{3}{2}-\epsilon} \backslash B_{\frac{1}{2}+\epsilon} .
$$

In particular, from the strict monotonicity of $v$ in the $e_{n}$-direction in $B_{2}^{+}(v)$ we have that

$$
v_{\epsilon t_{1}}>0 \quad \text { on } \mathcal{F}\left(v_{\epsilon\left(t_{1}-\delta\right)}\right),
$$

which combined with the previous inequality gives that

$$
g>0 \quad \text { on } \mathcal{F}\left(v_{\epsilon\left(t_{1}-\delta\right)}\right),
$$

that is the second condition in (3.7).
Finally, given $\epsilon>0$ small and a Lipschitz function $\tilde{\varphi}$ defined on $B_{\rho}(\bar{X})$, with values in $[-1,1]$, then there exists a unique function $\varphi_{\epsilon}$ defined at least on $B_{\rho-\epsilon}(\bar{X})$ such that

$$
\begin{equation*}
U(X)=\varphi_{\epsilon}\left(X-\epsilon \tilde{\varphi}(X) e_{n}\right), \quad X \in B_{\rho}(\bar{X}) . \tag{3.9}
\end{equation*}
$$

Moreover such function $\varphi_{\epsilon}$ is increasing in the $e_{n}$-direction.
With a similar argument as in Lemma 3.1 we can conclude that if $g$ satisfies the flatness assumption (3.2) in $B_{1}$ and $\tilde{\varphi}$ is as above then (say $\rho, \epsilon<1 / 4, \bar{X} \in B_{1 / 2}$,)

$$
\begin{equation*}
\tilde{\varphi} \leqslant \tilde{g}_{\epsilon} \quad \text { in } B_{\rho}(\bar{X}) \backslash P \quad \Rightarrow \quad \varphi_{\epsilon} \leqslant g \quad \text { in } B_{\rho-\epsilon}(\bar{X}) . \tag{3.10}
\end{equation*}
$$

We will use this fact in the proof of our improvement of flatness theorem.

## 4. The linearized problem

We introduce here the linearized problem associated to (2.1). Here and later $U_{n}$ denotes the $x_{n}$-derivative of the function $U$ defined in (1.3). Recall also that we denote by $P$ and $L$ the half-hyperplanes $P:=\left\{X \in \mathbb{R}^{n+1}: x_{n} \leqslant 0\right.$, $z=0\}$ and $L:=\left\{X \in \mathbb{R}^{n+1}: x_{n}=0, z=0\right\}$.

Given $w \in C\left(B_{1}\right)$ and $X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1} \cap L$, we call

$$
\left|\nabla_{r} w\right|\left(X_{0}\right):=\lim _{\left(x_{n}, z\right) \rightarrow(0,0)} \frac{w\left(x_{0}^{\prime}, x_{n}, z\right)-w\left(x_{0}^{\prime}, 0,0\right)}{r}, \quad r^{2}=x_{n}^{2}+z^{2}
$$

Once the change of unknowns (3.1) has been done, the linearized problem associated to (2.1) is

$$
\begin{cases}\Delta\left(U_{n} w\right)=0, & \text { in } B_{1} \backslash P  \tag{4.1}\\ \left|\nabla_{r} w\right|=0, & \text { on } B_{1} \cap L\end{cases}
$$

As we will show later in Section 8, if $w \in C\left(B_{1}\right)$ satisfies

$$
\Delta\left(U_{n} w\right)=0 \quad \text { in } B_{1} \backslash P,
$$

$w$ is even with respect to $\{z=0\}$, and $w$ is smooth in the $x^{\prime}$-direction, then given $X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1} \cap L$,

$$
\begin{equation*}
w(X)=w\left(X_{0}\right)+a \cdot\left(x^{\prime}-x_{0}^{\prime}\right)+b r+O\left(\left|x^{\prime}-x_{0}^{\prime}\right|^{2}+r^{3 / 2}\right), \tag{4.2}
\end{equation*}
$$

with $a \in \mathbb{R}^{n-1}, b \in \mathbb{R}$ depending on $X_{0}$.
This motivates our notion of viscosity solution for this problem which we define below.
Definition 4.1. We say that $w$ is a solution to (4.1) if $w \in C\left(B_{1}\right), w$ is even with respect to $\{z=0\}$ and it satisfies
(i) $\Delta\left(U_{n} w\right)=0$ in $B_{1} \backslash P$;
(ii) Let $\phi$ be continuous around $X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1} \cap L$ and satisfy

$$
\phi(X)=\phi\left(X_{0}\right)+a\left(X_{0}\right) \cdot\left(x^{\prime}-x_{0}^{\prime}\right)+b\left(X_{0}\right) r+O\left(\left|x^{\prime}-x_{0}^{\prime}\right|^{2}+r^{3 / 2}\right)
$$

with

$$
b\left(X_{0}\right) \neq 0 .
$$

If $b\left(X_{0}\right)>0$ then $\phi$ cannot touch $w$ by below at $X_{0}$, and if $b\left(X_{0}\right)<0$ then $\phi$ cannot touch $w$ by above at $X_{0}$.
In Section 8, we will show the following main regularity result about viscosity solutions to (4.1).

Theorem 4.2 (Improvement of flatness). There exists a universal constant $C$ such that if $w$ is a viscosity solution to (8.1) in $B_{1}$ with

$$
-1 \leqslant w(X) \leqslant 1 \quad \text { in } B_{1},
$$

then

$$
a_{0} \cdot x^{\prime}-C|X|^{3 / 2} \leqslant w(X)-w(0) \leqslant a_{0} \cdot x^{\prime}+C|X|^{3 / 2}
$$

for some vector $a_{0} \in \mathbb{R}^{n-1}$.
We conclude this short section with a remark which we will use in the proof of the theorem above.
Lemma 4.3. Let $w_{1}, w_{2} \in C\left(B_{1}\right)$ satisfy

$$
\Delta\left(U_{n} w_{i}\right)=0, \quad \text { in } B_{1} \backslash P, i=1,2
$$

Then $w_{1}$ and $w_{2}$ cannot touch (either by above or below) on $P \backslash L$, unless they coincide.
Proof. Assume by contradiction that

$$
w_{1}\left(X_{0}\right)=w_{2}\left(X_{0}\right), \quad X_{0} \in P \backslash L,
$$

and

$$
w_{1} \geqslant w_{2}, \quad \text { in } B_{\rho}\left(X_{0}\right)
$$

Then $U_{n}\left(w_{1}-w_{2}\right)$ is a non-negative harmonic function in $B_{1} \backslash P$ which vanishes continuously on $P \backslash L$. Hence unless $w_{1}=w_{2}$, by the boundary Harnack inequality (in the appropriate domain),

$$
U_{n}\left(w_{1}-w_{2}\right) \geqslant \delta U_{n} \quad \text { in } B_{\rho / 2}\left(X_{0}\right) \cap\{z>0\},
$$

for some small positive constant $\delta$. Thus

$$
w_{1}-w_{2} \geqslant \delta \quad \text { in } B_{\rho / 2}\left(X_{0}\right) \cap\{z>0\}
$$

and by continuity

$$
\left(w_{1}-w_{2}\right)\left(X_{0}\right)>0
$$

a contradiction.

## 5. Properties of $\boldsymbol{U}$

The first two lemmas in this section describe properties of $U$ which will be used in the proof of Harnack inequality, and in particular when constructing the barriers which are used in that proof.

The third lemma, which is incorporated here since its proof uses similar arguments to the proof of the first two lemmas, allows us to replace the assumptions in our main Theorem 1.1 with a more standard "flatness" assumption of the form

$$
U\left(X-\epsilon e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon e_{n}\right), \quad \text { in } B_{1} .
$$

Lemma 5.1. Let $g \in C\left(B_{2}\right), g \geqslant 0$ be a harmonic function in $B_{2}^{+}(g)$ and let $\bar{X}=\frac{3}{2} e_{n}$. Assume that

$$
g \geqslant U \quad \text { in } B_{2}, \quad g(\bar{X})-U(\bar{X}) \geqslant \delta_{0}
$$

for some $\delta_{0}>0$, then

$$
\begin{equation*}
g \geqslant\left(1+c \delta_{0}\right) U \quad \text { in } B_{1} \tag{5.1}
\end{equation*}
$$

for a small universal constant c. In particular, for any $0<\epsilon<2$

$$
\begin{equation*}
U\left(X+\epsilon e_{n}\right) \geqslant(1+c \epsilon) U(X) \quad \text { in } B_{1}, \tag{5.2}
\end{equation*}
$$

with c small universal.

Proof. Call $g^{*}$ the harmonic function in

$$
D=B_{3 / 2} \backslash\left\{x \in \mathcal{B}_{3 / 2}: x_{n} \leqslant 0\right\},
$$

such that

$$
g^{*}=g \quad \text { on } \partial B_{3 / 2}, \quad g^{*}=0 \quad \text { on }\left\{x \in \mathcal{B}_{3 / 2}: x_{n} \leqslant 0\right\} .
$$

Then by the maximum principle

$$
g \geqslant g^{*} \quad \text { on } \bar{B}_{3 / 2},
$$

and it suffices to show that (5.1) holds with $g^{*}$ on the left-hand side.
Since $g \geqslant U$ in $B_{2}$ we have

$$
g^{*}-U=g-U \geqslant 0 \quad \text { on } \partial B_{3 / 2}, \quad g^{*}-U=0 \quad \text { on }\left\{x \in \mathcal{B}_{3 / 2}: x_{n} \leqslant 0\right\},
$$

and hence $g^{*}-U \geqslant 0$ in $D$ where it is also harmonic. Moreover, from the assumption $g(\bar{X})-U(\bar{X}) \geqslant \delta_{0}$ we get by Harnack inequality that

$$
g^{*}-U=g-U \geqslant c_{0} \delta_{0} \quad \text { on } \partial B_{3 / 2} \cap B_{1 / 4}(\bar{X}) .
$$

Thus

$$
g^{*}(\widetilde{X})-U(\widetilde{X}) \geqslant c_{1} \delta_{0}, \quad \text { at some } \widetilde{X} \in B_{1} \cap D .
$$

Thus, by the boundary Harnack inequality we get that for $c>0$ universal,

$$
g^{*}-U \geqslant c_{2} \frac{g^{*}(\widetilde{X})-U(\widetilde{X})}{U(\widetilde{X})} U \geqslant c \delta_{0} U \quad \text { in } B_{1}
$$

as desired.
In particular, if $g(X)=U\left(X+\epsilon e_{n}\right)$ the assumptions of the lemma are satisfied. Indeed $U$ is monotone increasing in the $e_{n}$-direction thus $U\left(X+\epsilon e_{n}\right) \geqslant U(X)$ in $B_{2}$. Moreover,

$$
U\left(\bar{X}+\epsilon e_{n}\right)-U(\bar{X})=U_{t}\left(\bar{X}+\lambda e_{n}\right) \epsilon \geqslant c^{\prime} \epsilon, \quad \lambda \in(0, \epsilon),
$$

with $c^{\prime}$ universal.
Lemma 5.2. For any $\epsilon>0$ small, given $2 \epsilon<\bar{\delta}<1$, there exists a constant $C>0$ depending on $\bar{\delta}$ such that

$$
U(t+\epsilon, z) \leqslant(1+C \epsilon) U(t, z) \quad \text { in } \bar{B}_{1} \backslash B_{\bar{\delta}} \subset \mathbb{R}^{2} .
$$

Proof. In this lemma $B_{\rho}$ denotes a ball of radius $\rho$ in $\mathbb{R}^{2}$.
Since $U$ is monotone increasing in the $t$-direction, $U(t+\epsilon, z)-U(t, z)$ is non-negative and harmonic in the set $D_{\bar{\delta}}:=\left(B_{2} \backslash\{t \in(-2,2): t \leqslant 0\}\right) \backslash \bar{B}_{\bar{\delta} / 2}$. Moreover,

$$
U(3 / 2+\epsilon, 0)-U(3 / 2,0)=U_{t}(t, 0) \epsilon \leqslant C_{0} \epsilon, \quad t \in(3 / 2,3 / 2+\epsilon),
$$

with $C_{0}$ universal. By the boundary Harnack inequality in $D_{\bar{\delta}}$,

$$
U(t+\epsilon, z)-U(t, z) \leqslant C_{1} \frac{U(3 / 2+\epsilon, 0)-U(3 / 2,0)}{U(3 / 2,0)} U(t, z) \leqslant C \epsilon U(t, z) \quad \text { in } \bar{B}_{1} \backslash B_{\bar{\delta}},
$$

as desired.
Lemma 5.3. Let $g \in C\left(\bar{B}_{2}\right), g \geqslant 0$ be a harmonic function in $B_{2}^{+}(g)$ satisfying

$$
\begin{equation*}
\|g-U\|_{L^{\infty}\left(\bar{B}_{2}\right)} \leqslant \delta \tag{5.3}
\end{equation*}
$$

and

$$
\left\{x \in \mathcal{B}_{2}: x_{n} \leqslant-\delta\right\} \subset\left\{x \in \mathcal{B}_{2}: g(x, 0)=0\right\} \subset\left\{x \in \mathcal{B}_{2}: x_{n} \leqslant \delta\right\},
$$

with $\delta>0$ small universal. Then

$$
\begin{equation*}
U\left(X-\epsilon e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon e_{n}\right) \quad \text { in } B_{1}, \tag{5.4}
\end{equation*}
$$

for some $\epsilon=K \delta, K$ universal.

Proof. Let $\bar{g}$ be the harmonic function in

$$
B_{2}^{\delta}:=B_{2} \backslash\left\{x \in \mathcal{B}_{2}: x_{n} \leqslant-\delta\right\},
$$

such that

$$
\bar{g}=g \quad \text { on } \partial B_{2}, \quad \bar{g}=0 \quad \text { on }\left\{x \in \mathcal{B}_{2}: x_{n} \leqslant-\delta\right\} .
$$

Since $g$ is subharmonic in $B_{2}^{\delta}$, the maximum principle gives us

$$
g \leqslant \bar{g} \quad \text { on } \bar{B}_{2} .
$$

We need to show that for $K>0$ universal,

$$
\begin{equation*}
\bar{g} \leqslant U\left(X+K \delta e_{n}\right) \quad \text { in } B_{1} . \tag{5.5}
\end{equation*}
$$

(The lower bound follows from a similar argument.) Since $U$ is monotone increasing in the $e_{n}$-direction and it satisfies (5.3) we get that

$$
U\left(X+\delta e_{n}\right) \geqslant U(X) \geqslant g(X)-\delta \quad \text { on } \bar{B}_{2},
$$

and hence

$$
U\left(X+\delta e_{n}\right) \geqslant \bar{g}(X)-\delta \quad \text { on } \partial B_{2} .
$$

By the maximum principle in the domain $B_{2}^{\delta}$ we get that this inequality holds in $B_{2}$ and hence

$$
\begin{equation*}
\bar{g}(X)-U\left(X+\delta e_{n}\right) \leqslant \delta \quad \text { in } B_{2} . \tag{5.6}
\end{equation*}
$$

Let $g^{*}$ be the harmonic function in $B_{3 / 2} \backslash\left\{x \in \mathcal{B}_{3 / 2}: x_{n} \leqslant-\delta\right\}$ such that

$$
g^{*}=\delta \quad \text { on } \partial B_{3 / 2}, \quad g^{*}=0 \quad \text { on }\left\{x \in \mathcal{B}_{3 / 2}: x_{n} \leqslant-\delta\right\} .
$$

Clearly

$$
0 \leqslant g^{*} \leqslant \delta
$$

Then by the boundary Harnack inequality, say for $\bar{X}=e_{n}$

$$
\begin{equation*}
g^{*}(X) \leqslant \bar{C} \frac{g^{*}(\bar{X})}{U\left(\bar{X}+\delta e_{n}\right)} U\left(X+\delta e_{n}\right) \leqslant C \delta U\left(X+\delta e_{n}\right) \quad \text { in } B_{1}, \tag{5.7}
\end{equation*}
$$

with $C>0$ universal. Moreover, in view of (5.6) again by the maximum principle we have

$$
\bar{g}(X)-U\left(X+\delta e_{n}\right) \leqslant g^{*}(X) \quad \text { in } B_{3 / 2} .
$$

This inequality together with (5.7) gives that

$$
\bar{g}(X) \leqslant(1+C \delta) U\left(X+\delta e_{n}\right) \quad \text { in } B_{1} .
$$

By (5.2) (applied to a translate of $U$ ) we have that for $K>1$

$$
(1+C \delta) U\left(X+\delta e_{n}\right) \leqslant \frac{1+C \delta}{1+c K \delta} U\left(X+K \delta e_{n}\right) \leqslant U\left(X+K \delta e_{n}\right) \quad \text { in } B_{1},
$$

as long as $K$ is large enough. Combining these two last inequalities we obtain the desired claim (5.5).

## 6. Harnack inequality

In this section we state and prove a Harnack type inequality for solutions to our free boundary problem (2.1).
Theorem 6.1 (Harnack inequality). There exists $\bar{\epsilon}>0$ such that if $g$ solves (2.1) and it satisfies

$$
\begin{equation*}
U\left(X+\epsilon a_{0} e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon b_{0} e_{n}\right) \quad \text { in } B_{\rho}\left(X^{*}\right) \tag{6.1}
\end{equation*}
$$

with

$$
\epsilon\left(b_{0}-a_{0}\right) \leqslant \bar{\epsilon} \rho,
$$

then

$$
\begin{equation*}
U\left(X+\epsilon a_{1} e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon b_{1} e_{n}\right) \quad \text { in } B_{\eta \rho}\left(X^{*}\right), \tag{6.2}
\end{equation*}
$$

with

$$
a_{0} \leqslant a_{1} \leqslant b_{1} \leqslant b_{0}, \quad\left(b_{1}-a_{1}\right) \leqslant(1-\eta)\left(b_{0}-a_{0}\right),
$$

for a small universal constant $\eta$.
From this statement we get the desired corollary to be used in the proof of our main result. Precisely, if $g$ satisfies (6.1) with say $\rho=1 / 2$, then we can apply Harnack inequality repeatedly and obtain

$$
U\left(X+\epsilon a_{m} e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon b_{m} e_{n}\right) \quad \text { in } B_{\frac{1}{2} \eta^{m}}\left(X^{*}\right),
$$

with

$$
\begin{equation*}
b_{m}-a_{m} \leqslant\left(b_{0}-a_{0}\right)(1-\eta)^{m}, \tag{6.3}
\end{equation*}
$$

for all $m$ 's such that

$$
\begin{equation*}
2 \epsilon(1-\eta)^{m} \eta^{-m}\left(b_{0}-a_{0}\right) \leqslant \bar{\epsilon} . \tag{6.4}
\end{equation*}
$$

This implies that for all such $m$ 's, the function $\tilde{g}_{\epsilon}$ defined in Section 2.3 satisfies

$$
\begin{equation*}
a_{m} \leqslant \tilde{g}_{\epsilon}(X) \leqslant b_{m}, \quad \text { in } B_{\frac{1}{2} \eta^{m}-\epsilon}\left(X^{*}\right) \backslash P, \tag{6.5}
\end{equation*}
$$

with $a_{m}, b_{m}$ as in (6.3). Let $A_{\epsilon}$ be the following set

$$
\begin{equation*}
A_{\epsilon}:=\left\{\left(X, \tilde{g}_{\epsilon}(X)\right): X \in B_{1-\epsilon} \backslash P\right\} \subset \mathbb{R}^{n+1} \times\left[a_{0}, b_{0}\right] . \tag{6.6}
\end{equation*}
$$

Since $\tilde{g}_{\epsilon}$ may be multi-valued, we mean that given $X$ all pairs ( $X, \tilde{g}_{\epsilon}(X)$ ) belong to $A_{\epsilon}$ for all possible values of $\tilde{g}_{\epsilon}(X)$. In view of (6.5) we then get

$$
\begin{equation*}
A_{\epsilon} \cap\left(B_{\frac{1}{2} \eta^{m}-\epsilon}\left(X^{*}\right) \times\left[a_{0}, b_{0}\right]\right) \subset B_{\frac{1}{2} \eta^{m}-\epsilon}\left(X^{*}\right) \times\left[a_{m}, b_{m}\right], \tag{6.7}
\end{equation*}
$$

with $a_{m}, b_{m}$ as in (6.3) for all $m$ 's such that (6.4) holds.
Thus we get the following corollary.

## Corollary 6.2. If

$$
U\left(X-\epsilon e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon e_{n}\right) \quad \text { in } B_{1},
$$

with $\epsilon \leqslant \bar{\epsilon} / 2$, given $m_{0}>0$ such that

$$
2 \epsilon(1-\eta)^{m_{0}} \eta^{-m_{0}} \leqslant \bar{\epsilon},
$$

then the set $A_{\epsilon} \cap\left(B_{1 / 2} \times[-1,1]\right)$ is above the graph of a function $y=a_{\epsilon}(X)$ and it is below the graph of a function $y=b_{\epsilon}(X)$ with

$$
b_{\epsilon}-a_{\epsilon} \leqslant 2(1-\eta)^{m_{0}-1},
$$

and $a_{\epsilon}, b_{\epsilon}$ having a modulus of continuity bounded by the Hölder function $\alpha t^{\beta}$ for $\alpha, \beta$ depending only on $\eta$.
The proof of Harnack inequality will easily follow from the lemma below.
Lemma 6.3. There exists $\bar{\epsilon}>0$ such that for all $0<\epsilon \leqslant \bar{\epsilon}$ if $g$ is a solution to (2.1) in $B_{1}$ such that

$$
\begin{equation*}
g(X) \geqslant U(X) \quad \text { in } B_{1 / 2} \tag{6.8}
\end{equation*}
$$

and at $\bar{X} \in B_{1 / 8}\left(\frac{1}{4} e_{n}\right)$

$$
\begin{equation*}
g(\bar{X}) \geqslant U\left(\bar{X}+\epsilon e_{n}\right), \tag{6.9}
\end{equation*}
$$

then

$$
\begin{equation*}
g(X) \geqslant U\left(X+\tau \epsilon e_{n}\right) \quad \text { in } B_{\delta} \tag{6.10}
\end{equation*}
$$

for universal constants $\tau, \delta$. Similarly, if

$$
g(X) \leqslant U(X) \quad \text { in } B_{1 / 2},
$$

and

$$
g(\bar{X}) \leqslant U\left(\bar{X}-\epsilon e_{n}\right),
$$

then

$$
g(X) \leqslant U\left(X-\tau \epsilon e_{n}\right) \quad \text { in } B_{\delta} .
$$

The main tool in the proof of Lemma 6.3 will be the following family of radial subsolutions. Let $R>0$ and denote

$$
V_{R}(t, z)=U(t, z)\left((n-1) \frac{t}{R}+1\right)
$$

Then set

$$
\begin{equation*}
v_{R}(X)=V_{R}\left(R-\sqrt{\left|x^{\prime}\right|^{2}+\left(x_{n}-R\right)^{2}}, z\right) \tag{6.11}
\end{equation*}
$$

that is we obtain the $(n+1)$-dimensional function $v_{R}$ by rotating the 2 -dimensional function $V_{R}$ around $(0, R, z)$.
Proposition 6.4. If $R$ is large enough, the function $v_{R}(X)$ is a comparison subsolution to (2.1) in $B_{2}$ which is strictly monotone increasing in the $e_{n}$-direction in $B_{2}^{+}\left(v_{R}\right)$. Moreover, there exists a function $\tilde{v}_{R}$ such that

$$
U(X)=v_{R}\left(X-\tilde{v}_{R}(X) e_{n}\right) \quad \text { in } B_{1} \backslash P,
$$

and

$$
\left|\tilde{v}_{R}(X)-\gamma_{R}(X)\right| \leqslant \frac{C}{R^{2}}|X|^{2}, \quad \gamma_{R}(X)=-\frac{\left|x^{\prime}\right|^{2}}{2 R}+2(n-1) \frac{x_{n} r}{R}
$$

with $r=\sqrt{x_{n}^{2}+z^{2}}$ and $C$ universal.
The proof of Proposition 6.4 follows from long and tedious computations and we postpone it till Appendix A. Using the estimate for $\tilde{v}_{R}$ in Proposition 6.4 and Lemma 3.1, we also obtain the following corollary which will be crucial for the proof of Lemma 6.3. Its proof is again presented in Appendix A.

Corollary 6.5. There exist $\delta, c_{0}, C_{0}, C_{1}$ universal constants, such that

$$
\begin{equation*}
v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \leqslant\left(1+\frac{C_{0}}{R}\right) U(X), \quad \text { in } \bar{B}_{1} \backslash B_{1 / 4}, \tag{6.12}
\end{equation*}
$$

with strict inequality on $F\left(v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right)\right) \cap \bar{B}_{1} \backslash B_{1 / 4}$,

$$
\begin{align*}
& v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \geqslant U\left(X+\frac{c_{0}}{2 R} e_{n}\right), \quad \text { in } B_{\delta},  \tag{6.13}\\
& v_{R}\left(X-\frac{C_{1}}{R} e_{n}\right) \leqslant U(X), \quad \text { in } \bar{B}_{1} . \tag{6.14}
\end{align*}
$$

We are now ready to present the proof of Lemma 6.3.
Proof of Lemma 6.3. We prove the first statement. In view of (6.9)

$$
g(\bar{X})-U(\bar{X}) \geqslant U\left(\bar{X}+\epsilon e_{n}\right)-U(\bar{X})=\partial_{t} U\left(\bar{X}+\lambda e_{n}\right) \epsilon \geqslant c \epsilon, \quad \lambda \in(0, \epsilon) .
$$

As in Lemma 5.1 we then get

$$
\begin{equation*}
g(X) \geqslant\left(1+c^{\prime} \epsilon\right) U(X) \quad \text { in } \bar{B}_{1 / 4} \tag{6.15}
\end{equation*}
$$

Now let

$$
R=\frac{C_{0}}{c^{\prime} \epsilon}
$$

where from now on the $C_{i}, c_{i}$ are the constants in Corollary 6.5. Then, for $\epsilon$ small enough $v_{R}$ is a subsolution to (2.1) in $B_{2}$ which is monotone increasing in the $e_{n}$-direction and it also satisfies (6.12)-(6.14). We now wish to apply the comparison principle as stated in Corollary 2.7. Let

$$
v_{R}^{t}(X)=v_{R}\left(X+t e_{n}\right), \quad X \in B_{1}
$$

then according to (6.14),

$$
v_{R}^{t_{0}} \leqslant U \leqslant g \quad \text { in } \bar{B}_{1 / 4}, \text { with } t_{0}=-C_{1} / R
$$

Moreover, from (6.12) and (6.15) we get that for our choice of $R$,

$$
v_{R}^{t_{1}} \leqslant\left(1+c^{\prime} \epsilon\right) U \leqslant g \quad \text { on } \partial B_{1 / 4}, \text { with } t_{1}=c_{0} / R
$$

with strict inequality on $F\left(v_{R}^{t_{1}}\right) \cap \partial B_{1 / 4}$. In particular

$$
g>0 \quad \text { on } \mathcal{F}\left(v_{R}^{t_{1}}\right) \text { in } \partial B_{1 / 4}
$$

Thus we can apply Corollary 2.7 in the ball $B_{1 / 4}$ to obtain

$$
v_{R}^{t_{1}} \leqslant g \quad \text { in } B_{1 / 4}
$$

From (6.13) we have that

$$
U\left(X+\frac{c_{1}}{R} e_{n}\right) \leqslant v_{R}^{t_{1}}(X) \leqslant g(X) \quad \text { on } B_{\delta}
$$

which is the desired claim (6.10) with $\tau=\frac{c_{1} c^{\prime}}{C_{0}}$.
We now present the proof of the Harnack inequality.
Proof of Theorem 6.1. Without loss of generality, we can assume $a_{0}=-1, b_{0}=1$. Also, in view of Remark 2.5 we can take $\rho=1$ (thus $2 \epsilon \leqslant \bar{\epsilon}$ ).

We distinguish several cases. In what follows $\bar{\epsilon}$ and $\delta$ denote the universal constants in Lemma 6.3.
Case 1. If

$$
d\left(X^{*},\left\{x_{n} \geqslant \epsilon, z=0\right\}\right)>\delta / 16
$$

then $U\left(X-\epsilon e_{n}\right)>0$ in $B_{\delta / 16}\left(X^{*}\right) \subset B_{1}\left(X^{*}\right)$. Thus the functions $U\left(X-\epsilon e_{n}\right), U\left(X+\epsilon e_{n}\right)$ and $g(X)$ are positive and harmonic in $B_{\delta / 16}\left(X^{*}\right)$. Assume that (the other case is treated similarly)

$$
g\left(X^{*}\right) \geqslant U\left(X^{*}\right)
$$

Then,

$$
g\left(X^{*}\right) \geqslant U\left(X^{*}\right)=U\left(X^{*}-\epsilon e_{n}\right)+U_{n}\left(X^{*}-\lambda \epsilon e_{n}\right) \epsilon, \quad \lambda \in(0,1)
$$

Since $U_{n}$ is positive and harmonic in $B_{\delta / 16}\left(X^{*}\right)$ and for $\bar{\epsilon}<\delta / 16$

$$
X^{*}-\lambda \epsilon e_{n} \in B_{\delta / 32}\left(X^{*}\right)
$$

we can apply Harnack inequality to conclude that

$$
g\left(X^{*}\right) \geqslant U\left(X^{*}-\epsilon e_{n}\right)+c U_{n}\left(X^{*}\right) \epsilon
$$

for $c$ small universal.

Then again by Harnack inequality in $B_{\delta / 16}\left(X^{*}\right)$ for $g(X)-U\left(X-\epsilon e_{n}\right) \geqslant 0$ we get that for $c^{\prime}$ universal

$$
g(X) \geqslant U\left(X-\epsilon e_{n}\right)+c^{\prime} U_{n}\left(X^{*}\right) \epsilon, \quad \text { in } B_{\delta / 32}\left(X^{*}\right) .
$$

By a similar argument, for $\bar{\epsilon}<\delta / 32$

$$
U\left(X-(1-\eta) \epsilon e_{n}\right)-U\left(X-\epsilon e_{n}\right) \leqslant C U_{n}\left(X^{*}\right) \eta \epsilon, \quad \text { in } B_{\delta / 64}\left(X^{*}\right),
$$

with $C$ universal.
Thus, combining these two last inequalities we obtain that for $\eta=\min \left\{c^{\prime} / C, \delta / 64\right\}$

$$
g(X) \geqslant U\left(X-(1-\eta) \epsilon e_{n}\right), \quad \text { in } B_{\eta}\left(X^{*}\right),
$$

as desired.

## Case 2. If

$$
d\left(X^{*},\left\{x_{n}=-\epsilon, z=0\right\}\right) \leqslant \delta / 2,
$$

we wish to apply Lemma 6.3. Then (for $\bar{\epsilon}<\delta / 4$ )

$$
g(X) \geqslant U\left(X-\epsilon e_{n}\right) \quad \text { in } B_{1 / 2}\left(\epsilon e_{n}\right) \subset B_{1}\left(X^{*}\right) .
$$

Let $\bar{X}=\frac{1}{4} e_{n}$ and assume that (the other case follows similarly)

$$
g(\bar{X}) \geqslant U(\bar{X})
$$

Since (for $\bar{\epsilon}$ small)

$$
\bar{X} \in B_{1 / 8}\left(\left(\frac{1}{4}+\epsilon\right) e_{n}\right),
$$

we can apply Lemma 6.3 and conclude that

$$
g(X) \geqslant U\left(X-(1-\eta) \epsilon e_{n}\right), \quad \text { in } B_{\delta}\left(\epsilon e_{n}\right)
$$

Thus the desired improvement holds by choosing $\eta<\delta / 4$. Indeed for such $\eta$ and $\bar{\epsilon}<\delta / 4$ we have that $d\left(X^{*}\right.$, $\left.\left\{x_{n}=\epsilon, z=0\right\}\right) \leqslant 3 \delta / 4$ and hence

$$
B_{\eta}\left(X^{*}\right) \subset B_{\delta}\left(\epsilon e_{n}\right) .
$$

## Case 3. If

$$
d\left(X^{*},\left\{x_{n}=-\epsilon, z=0\right\}\right)>\delta / 2 \quad \text { and } \quad d\left(X^{*},\left\{x_{n} \geqslant \epsilon, z=0\right\}\right) \leqslant \delta / 16,
$$

then the functions $U\left(X-\epsilon e_{n}\right), U\left(X+\epsilon e_{n}\right)$ and $g(X)$ are positive and harmonic in the half-ball $B_{\delta / 4}(\widetilde{X}) \cap\{z>0\}$ for some $\widetilde{X} \in\left\{x_{n} \leqslant-\epsilon, z=0\right\}$ and they all vanish continuously on $B_{\delta / 4}(\widetilde{X}) \cap\{z=0\}$. Thus we can repeat a similar argument as in Case 1, by using the boundary Harnack inequality. Precisely, let $\bar{X}=\widetilde{X}+\frac{\delta}{6} e_{n+1}$ and assume that (the other case is treated similarly)

$$
g(\bar{X}) \geqslant U(\bar{X})
$$

Then,

$$
g(\bar{X}) \geqslant U(\bar{X})=U\left(\bar{X}-\epsilon e_{n}\right)+U_{n}\left(\bar{X}-\lambda \epsilon e_{n}\right) \epsilon, \quad \lambda \in(0,1) .
$$

By Harnack inequality for $U_{n}$ in the ball $B_{2 \epsilon}(\bar{X}) \subset B_{\delta / 4}(\widetilde{X}) \cap\{z>0\}$ (with $\bar{\epsilon}<\delta / 12$ ) we conclude that

$$
\begin{equation*}
g(\bar{X}) \geqslant U\left(\bar{X}-\epsilon e_{n}\right)+c U_{n}(\bar{X}) \epsilon, \tag{6.16}
\end{equation*}
$$

for $c$ small universal.
Then by boundary Harnack inequality in $B_{\delta / 4}(\widetilde{X}) \cap\{z>0\}$, for the functions $g(X)-U\left(X-\epsilon e_{n}\right)$ and $U_{n}(X)$ we get that for $c^{\prime}$ universal

$$
\begin{equation*}
g(X) \geqslant U\left(X-\epsilon e_{n}\right)+c^{\prime} U_{n}(X) \epsilon, \quad \text { in } B_{\delta / 8}(\widetilde{X}) \cap\{z \geqslant 0\} . \tag{6.17}
\end{equation*}
$$

Thus to obtain the desired claim it is enough to choose $\eta$ small such that for $X \in B_{\delta / 8}(\tilde{X}) \cap\{z \geqslant 0\}$

$$
U\left(X-\epsilon e_{n}\right)+c^{\prime} U_{n}(X) \epsilon \geqslant U\left(X-(1-\eta) \epsilon e_{n}\right) .
$$

By a similar argument as above

$$
U\left(\bar{X}-(1-\eta) \epsilon e_{n}\right)-U\left(\bar{X}-\epsilon e_{n}\right) \leqslant C U_{n}(\bar{X}) \eta \epsilon,
$$

and hence by boundary Harnack inequality,

$$
U\left(X-(1-\eta) \epsilon e_{n}\right)-U\left(X-\epsilon e_{n}\right) \leqslant C^{\prime} U_{n}(X) \eta \epsilon, \quad \text { in } B_{\delta / 8}(\widetilde{X}) \cap\{z \geqslant 0\} .
$$

Combining this inequality with (6.17) we obtain that for $\eta=c^{\prime} / C^{\prime}$

$$
g(X) \geqslant U\left(X-(1-\eta) \epsilon e_{n}\right), \quad \text { in } B_{\delta / 8}(\widetilde{X}) \cap\{z \geqslant 0\} .
$$

Since all the functions involved are even with respect to $\{z=0\}$ and for $\eta<\delta / 16$

$$
B_{\eta}\left(X^{*}\right) \subset B_{\delta / 8}(\tilde{X}),
$$

our proof is complete.

## 7. Improvement of flatness

In this section we state the improvement of flatness property for solutions to (2.1) and we provide its proof. Our main Theorem 1.1 follows from Theorem 7.1 and Lemma 5.3.

Theorem 7.1 (Improvement of flatness). There exist $\bar{\epsilon}>0$ and $\rho>0$ universal constants such that for all $0<\epsilon \leqslant \bar{\epsilon}$ if $g$ solves (2.1) with $0 \in F(g)$ and it satisfies

$$
\begin{equation*}
U\left(X-\epsilon e_{n}\right) \leqslant g(X) \leqslant U\left(X+\epsilon e_{n}\right) \quad \text { in } B_{1}, \tag{7.1}
\end{equation*}
$$

then

$$
\begin{equation*}
U\left(x \cdot v-\frac{\epsilon}{2} \rho, z\right) \leqslant g(X) \leqslant U\left(x \cdot v+\frac{\epsilon}{2} \rho, z\right) \text { in } B_{\rho}, \tag{7.2}
\end{equation*}
$$

for some direction $v \in \mathbb{R}^{n},|\nu|=1$.
The proof of Theorem 7.1 will easily follow from the next four lemmas.
Lemma 7.2. Let $g$ be a solution to (2.1) with $0 \in F(g)$ and satisfying (7.1). Assume that the corresponding $\tilde{g}_{\epsilon}$ satisfies

$$
\begin{equation*}
a_{0} \cdot x^{\prime}-\frac{1}{4} \rho \leqslant \tilde{g}_{\epsilon}(X) \leqslant a_{0} \cdot x^{\prime}+\frac{1}{4} \rho \quad \text { in } B_{2 \rho} \backslash P, \tag{7.3}
\end{equation*}
$$

for some $a_{0} \in \mathbb{R}^{n-1}$. Then if $\epsilon \leqslant \bar{\epsilon}\left(a_{0}, \rho\right)$, $g$ satisfies (7.2) in $B_{\rho}$.
Proof. We prove that the lower bound holds (the upper bound can be proved similarly).
Let,

$$
v=\left(v^{\prime}, v_{n}\right):=\frac{(0,1)+\epsilon\left(a_{0}, 0\right)}{\sqrt{1+\epsilon^{2} a_{0}^{2}}}
$$

and call

$$
u(X)=U\left(x \cdot v-\frac{\epsilon}{2} \rho, z\right) .
$$

Notice that since $\nu_{n}>0, u$ is strictly monotone increasing in the $e_{n}$-direction say in $B_{2 \rho}^{+}(u)$. Also, we can easily compute $\tilde{u}_{\epsilon}$ by its definition. Indeed, the identity

$$
u\left(X-\epsilon \tilde{u}_{\epsilon}(X) e_{n}\right)=U(X), \quad X \in \mathbb{R}^{n+1} \backslash P
$$

reads as

$$
U\left(x^{\prime} \cdot v^{\prime}+x_{n} v_{n}-\epsilon \tilde{u}_{\epsilon}(X) v_{n}-\frac{\epsilon}{2} \rho, z\right)=U\left(x_{n}, z\right)
$$

and hence

$$
\begin{equation*}
\tilde{u}_{\epsilon}(X)=\frac{x^{\prime} \cdot v^{\prime}+\left(v_{n}-1\right) x_{n}}{\epsilon v_{n}}-\frac{\rho}{2 v_{n}} . \tag{7.4}
\end{equation*}
$$

Thus, according to Lemma 3.1 it suffices to show that

$$
\tilde{u}_{\epsilon} \leqslant \tilde{g}_{\epsilon} \quad \text { in } B_{\rho+\epsilon} \backslash P,
$$

and hence in view of (7.3) we must show that

$$
\tilde{u}_{\epsilon}(X) \leqslant a_{0} \cdot x^{\prime}-\frac{1}{4} \rho \quad \text { in } B_{\rho+\epsilon} \backslash P .
$$

From the choice of $v$ we see that

$$
\frac{\nu^{\prime}}{\epsilon v_{n}}=a_{0},
$$

and

$$
\frac{\left|v_{n}-1\right|}{\epsilon v_{n}}=\frac{1-v_{n}}{\epsilon v_{n}} \leqslant \epsilon a_{0}^{2} .
$$

Thus, in view of the formula (7.4) the desired inequality reduces to

$$
x^{\prime} \cdot a_{0}+2 \rho \epsilon a_{0}^{2}-\frac{\rho}{2} \leqslant x^{\prime} \cdot a_{0}-\frac{\rho}{4},
$$

which is trivially satisfied for $\epsilon$ small enough (depending on $a_{0}, \rho$ ).
The next lemma follows immediately from Corollary 6.2 to Harnack inequality.
Lemma 7.3. Let $\epsilon_{k} \rightarrow 0$ and let $g_{k}$ be a sequence of solutions to (2.1) with $0 \in F\left(g_{k}\right)$ satisfying

$$
\begin{equation*}
U\left(X-\epsilon_{k} e_{n}\right) \leqslant g_{k}(X) \leqslant U\left(X+\epsilon_{k} e_{n}\right) \quad \text { in } B_{1} . \tag{7.5}
\end{equation*}
$$

Denote by $\tilde{g}_{k}$ the $\epsilon_{k}$-domain variation of $g_{k}$. Then the sequence of sets

$$
A_{k}:=\left\{\left(X, \tilde{g}_{k}(X)\right): X \in B_{1-\epsilon_{k}} \backslash P\right\}
$$

has a subsequence that converges uniformly (in Hausdorff distance) in $B_{1 / 2} \backslash P$ to the graph

$$
A_{\infty}:=\left\{\left(X, \tilde{g}_{\infty}(X)\right): X \in B_{1 / 2} \backslash P\right\}
$$

where $\tilde{g}_{\infty}$ is a Holder continuous function.
From here on $\tilde{g}_{\infty}$ will denote the function from Lemma 7.3.
Lemma 7.4. The function $\tilde{g}_{\infty}$ satisfies the linearized problem (4.1) in $B_{1 / 2}$.
Proof. We start by showing that $U_{n} \tilde{g}_{\infty}$ is harmonic in $B_{1 / 2} \backslash P$.
Let $\tilde{\varphi}$ be a $C^{2}$ function which touches $\tilde{g}_{\infty}$ strictly by below at $X_{0} \in B_{1 / 2} \backslash P$. We need to show that

$$
\begin{equation*}
\Delta\left(U_{n} \tilde{\varphi}\right)\left(X_{0}\right) \leqslant 0 . \tag{7.6}
\end{equation*}
$$

Since by the previous lemma, the sequence $A_{k}$ converges uniformly to $A_{\infty}$ in $B_{1 / 2} \backslash P$ we conclude that there exist a sequence of constants $c_{k} \rightarrow 0$ and a sequence of points $X_{k} \in B_{1 / 2} \backslash P, X_{k} \rightarrow X_{0}$ such that $\tilde{\varphi}_{k}:=\tilde{\varphi}+c_{k}$ touches $\tilde{g}_{k}$ by below at $X_{k}$ for all $k$ large enough.

Define the function $\varphi_{k}$ by the following identity

$$
\begin{equation*}
\varphi_{k}\left(X-\epsilon_{k} \tilde{\varphi}_{k}(X) e_{n}\right)=U(X) \tag{7.7}
\end{equation*}
$$

Then according to (3.10) $\varphi_{k}$ touches $g_{k}$ by below at $Y_{k}=X_{k}-\epsilon_{k} \tilde{\varphi}_{k}\left(X_{k}\right) e_{n} \in B_{1}^{+}\left(g_{k}\right)$, for $k$ large enough. Thus, since $g_{k}$ satisfies (2.1) in $B_{1}$ it follows that

$$
\begin{equation*}
\Delta \varphi_{k}\left(Y_{k}\right) \leqslant 0 \tag{7.8}
\end{equation*}
$$

Let us compute $\Delta \varphi_{k}\left(Y_{k}\right)$. Since $\tilde{\varphi}$ is smooth, for any $Y$ in a neighborhood of $Y_{k}$ we can find a unique $X=X(Y)$ such that

$$
\begin{equation*}
Y=X-\epsilon_{k} \tilde{\varphi}_{k}(X) e_{n} \tag{7.9}
\end{equation*}
$$

Thus (7.7) reads

$$
\varphi_{k}(Y)=U(X(Y))
$$

with $Y_{i}=X_{i}$ if $i \neq n$ and

$$
\frac{\partial X_{j}}{\partial Y_{i}}=\delta_{i j}, \quad \text { when } j \neq n
$$

Using these identities we can compute that

$$
\begin{equation*}
\Delta \varphi_{k}(Y)=U_{n}(X) \Delta X_{n}(Y)+\sum_{j \neq n}\left(U_{j j}(X)+2 U_{j n}(X) \frac{\partial X_{n}}{\partial Y_{j}}\right)+U_{n n}(X)\left|\nabla X_{n}\right|^{2}(Y) . \tag{7.10}
\end{equation*}
$$

From (7.9) we have that

$$
D_{X} Y=I-\epsilon_{k} D_{X}\left(\tilde{\varphi}_{k} e_{n}\right)
$$

Thus, since $\tilde{\varphi}_{k}=\tilde{\varphi}+c_{k}$

$$
D_{Y} X=I+\epsilon_{k} D_{X}\left(\tilde{\varphi} e_{n}\right)+O\left(\epsilon_{k}^{2}\right),
$$

with a constant depending only on the $C^{2}$-norm of $\tilde{\varphi}$.
It follows that

$$
\begin{equation*}
\frac{\partial X_{n}}{\partial Y_{j}}=\delta_{j n}+\epsilon_{k} \partial_{j} \tilde{\varphi}(X)+O\left(\epsilon_{k}^{2}\right) \tag{7.11}
\end{equation*}
$$

Hence

$$
\begin{equation*}
\left|\nabla X_{n}\right|^{2}(Y)=1+2 \epsilon_{k} \partial_{n} \tilde{\varphi}(X)+O\left(\epsilon_{k}^{2}\right) \tag{7.12}
\end{equation*}
$$

and also,

$$
\frac{\partial^{2} X_{n}}{\partial Y_{j}^{2}}=\epsilon_{k} \sum_{i} \partial_{j i} \tilde{\varphi} \frac{\partial X_{i}}{\partial Y_{j}}+O\left(\epsilon_{k}^{2}\right)=\epsilon_{k} \sum_{i \neq n} \partial_{j i} \tilde{\varphi} \delta_{i j}+\epsilon_{k} \partial_{j n} \tilde{\varphi} \frac{\partial X_{n}}{\partial Y_{j}}+O\left(\epsilon_{k}^{2}\right),
$$

from which we obtain that

$$
\begin{equation*}
\Delta X_{n}=\epsilon_{k} \Delta \tilde{\varphi}+O\left(\epsilon_{k}^{2}\right) \tag{7.13}
\end{equation*}
$$

Combining (7.10) with (7.12) and (7.13) we get that

$$
\Delta \varphi_{k}(Y)=\Delta U(X)+\epsilon_{k} U_{n} \Delta \tilde{\varphi}+2 \epsilon_{k} \nabla \tilde{\varphi} \cdot \nabla U_{n}+O\left(\epsilon_{k}^{2}\right)\left(U_{n}(X)+U_{n n}(X)\right)
$$

Using (7.8) together with the fact that $U$ is harmonic at $X_{k}$ we conclude that

$$
0 \geqslant \Delta\left(U_{n} \tilde{\varphi}\right)\left(X_{k}\right)+O\left(\epsilon_{k}^{2}\right)\left(U_{n}\left(X_{k}\right)+U_{n n}\left(X_{k}\right)\right)
$$

The desired inequality (7.6) follows by letting $k \rightarrow+\infty$.
Next we need to show that

$$
\left|\nabla_{r} \tilde{g}_{\infty}\right|\left(X_{0}\right)=0, \quad X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1 / 2} \cap L,
$$

in the viscosity sense of Definition 4.1.

Assume by contradiction that there exists a function $\phi$ which touches $\tilde{g}_{\infty}$ by below at $X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1 / 2} \cap L$ and such that

$$
\phi(X)=\phi\left(X_{0}\right)+a\left(X_{0}\right) \cdot\left(x^{\prime}-x_{0}^{\prime}\right)+b\left(X_{0}\right) r+O\left(\left|x^{\prime}-x_{0}^{\prime}\right|^{2}+r^{3 / 2}\right)
$$

with

$$
b\left(X_{0}\right)>0 .
$$

Then we can find constants $\alpha, \delta, \bar{r}$ and a point $Y^{\prime}=\left(y_{0}^{\prime}, 0,0\right) \in B_{2}$ depending on $\phi$ such that the polynomial

$$
q(X)=\phi\left(X_{0}\right)-\frac{\alpha}{2}\left|x^{\prime}-y_{0}^{\prime}\right|^{2}+2 \alpha(n-1) x_{n} r
$$

touches $\phi$ by below at $X_{0}$ in a tubular neighborhood $N_{\bar{r}}=\left\{\left|x^{\prime}-x_{0}^{\prime}\right| \leqslant \bar{r}, r \leqslant \bar{r}\right\}$ of $X_{0}$, with

$$
\phi-q \geqslant \delta>0, \quad \text { on } N_{\bar{r}} \backslash N_{\bar{r} / 2}
$$

This implies that

$$
\begin{equation*}
\tilde{g}_{\infty}-q \geqslant \delta>0, \quad \text { on } N_{\bar{r}} \backslash N_{\bar{r} / 2}, \tag{7.14}
\end{equation*}
$$

and

$$
\begin{equation*}
\tilde{g}_{\infty}\left(X_{0}\right)-q\left(X_{0}\right)=0 . \tag{7.15}
\end{equation*}
$$

In particular,

$$
\begin{equation*}
\left|\tilde{g}_{\infty}\left(X_{k}\right)-q\left(X_{k}\right)\right| \rightarrow 0, \quad X_{k} \in N_{\bar{r}} \backslash P, X_{k} \rightarrow X_{0} \tag{7.16}
\end{equation*}
$$

Now, let us choose $R_{k}=1 /\left(\alpha \epsilon_{k}\right)$ and let us define

$$
w_{k}(X)=v_{R_{k}}\left(X-Y^{\prime}+\epsilon_{k} \phi\left(X_{0}\right) e_{n}\right), \quad Y^{\prime}=\left(y_{0}^{\prime}, 0,0\right)
$$

with $v_{R}$ the function defined in Proposition 6.4. Then the $\epsilon_{k}$-domain variation of $w_{k}$, which we call $\tilde{w}_{k}$, can be easily computed from the definition

$$
w_{k}\left(X-\epsilon_{k} \tilde{w}_{k}(X) e_{n}\right)=U(X)
$$

Indeed, since $U$ is constant in the $x^{\prime}$-direction, this identity is equivalent to

$$
v_{R_{k}}\left(X-Y^{\prime}+\epsilon_{k} \phi\left(X_{0}\right) e_{n}-\epsilon_{k} \tilde{w}_{k}(X) e_{n}\right)=U\left(X-Y^{\prime}\right)
$$

which in view of Proposition 6.4 gives us

$$
\tilde{v}_{R_{k}}\left(X-Y^{\prime}\right)=\epsilon_{k}\left(\tilde{w}_{k}(X)-\phi\left(X_{0}\right)\right) .
$$

From the choice of $R_{k}$, the formula for $q$ and (A.2), we then conclude that

$$
\tilde{w}_{k}(X)=q(X)+\alpha^{2} \epsilon_{k} O\left(\left|X-Y^{\prime}\right|^{2}\right),
$$

and hence

$$
\begin{equation*}
\left|\tilde{w}_{k}-q\right| \leqslant C \epsilon_{k} \quad \text { in } N_{\bar{r}} \backslash P . \tag{7.17}
\end{equation*}
$$

Thus, from the uniform convergence of $A_{k}$ to $A_{\infty}$ and (7.14)-(7.17) we get that for all $k$ large enough

$$
\begin{equation*}
\tilde{g}_{k}-\tilde{w}_{k} \geqslant \frac{\delta}{2} \quad \text { in }\left(N_{\bar{r}} \backslash N_{\bar{r} / 2}\right) \backslash P . \tag{7.18}
\end{equation*}
$$

Similarly, from the uniform convergence of $A_{k}$ to $A_{\infty}$ and (7.17)-(7.16) we get that for $k$ large

$$
\begin{equation*}
\tilde{g}_{k}\left(X_{k}\right)-\tilde{w}_{k}\left(X_{k}\right) \leqslant \frac{\delta}{4}, \quad \text { for some sequence } X_{k} \in N_{\bar{r}} \backslash P, X_{k} \rightarrow X_{0} . \tag{7.19}
\end{equation*}
$$

On the other hand, it follows from Lemma 3.2 and (7.18) that

$$
\tilde{g}_{k}-\tilde{w}_{k} \geqslant \frac{\delta}{2} \quad \text { in } N_{\bar{r}} \backslash P,
$$

which contradicts (7.19).

The lemmas above allow us to reduce the regularity question for our free boundary problem (2.1) to the regularity of the linear problem (4.1). We will analyze such question in the next section and we will consequently obtain the following lemma, which we use here to conclude the proof of our improvement of flatness theorem.

Lemma 7.5. There exists a universal constant $\rho>0$ such that $\tilde{g}_{\infty}$ satisfies

$$
\begin{equation*}
a_{0} \cdot x^{\prime}-\frac{1}{8} \rho \leqslant \tilde{g}_{\infty}(X) \leqslant a_{0} \cdot x^{\prime}+\frac{1}{8} \rho \quad \text { in } B_{2 \rho} \text {, } \tag{7.20}
\end{equation*}
$$

for a vector $a_{0} \in \mathbb{R}^{n-1}$.
We are now ready to prove our main theorem, by combining all the lemmas above.
Proof of Theorem 7.1. Let $\rho$ be the universal constant from Lemma 7.5 and assume by contradiction that we can find a sequence $\epsilon_{k} \rightarrow 0$ and a sequence $g_{k}$ of solutions to (2.1) in $B_{1}$ such that $g_{k}$ satisfies (7.1), i.e.

$$
\begin{equation*}
U\left(X-\epsilon_{k} e_{n}\right) \leqslant g_{k}(X) \leqslant U\left(X+\epsilon_{k} e_{n}\right) \quad \text { in } B_{1}, \tag{7.21}
\end{equation*}
$$

but it does not satisfy the conclusion of the theorem.
Denote by $\tilde{g}_{k}$ the $\epsilon_{k}$-domain variation of $g_{k}$. Then by Lemma 7.3 the sequence of sets

$$
A_{k}:=\left\{\left(X, \tilde{g}_{k}(X)\right): X \in B_{1-\epsilon_{k}} \backslash P\right\}
$$

converges uniformly (up to extracting a subsequence) in $B_{1 / 2} \backslash P$ to the graph

$$
A_{\infty}:=\left\{\left(X, \tilde{g}_{\infty}(X)\right): X \in B_{1 / 2} \backslash P\right\}
$$

where $\tilde{g}_{\infty}$ is a Holder continuous function in $B_{1 / 2}$. By Lemma 7.4, the function $\tilde{g}_{\infty}$ solves the linearized problem (4.1) and hence by Lemma $7.5 \tilde{g}_{\infty}$ satisfies

$$
\begin{equation*}
a_{0} \cdot x^{\prime}-\frac{1}{8} \rho \leqslant \tilde{g}_{\infty}(X) \leqslant a_{0} \cdot x^{\prime}+\frac{1}{8} \rho \quad \text { in } B_{2 \rho}, \tag{7.22}
\end{equation*}
$$

with $a_{0} \in \mathbb{R}^{n-1}$.
From the uniform convergence of $A_{k}$ to $A_{\infty}$, we get that for all $k$ large enough

$$
\begin{equation*}
a_{0} \cdot x^{\prime}-\frac{1}{4} \rho \leqslant \tilde{g}_{k}(X) \leqslant a_{0} \cdot x^{\prime}+\frac{1}{4} \rho \quad \text { in } B_{2 \rho} \backslash P \tag{7.23}
\end{equation*}
$$

and hence from Lemma 7.2, the $g_{k}$ satisfy the conclusion of our theorem (for $k$ large). We have thus reached a contradiction.

## 8. The regularity of the linearized problem

The purpose of this section is to prove an improvement of flatness result for viscosity solutions to the linearized problem associated to (2.1), that is

$$
\begin{cases}\Delta\left(U_{n} w\right)=0, & \text { in } B_{1} \backslash P  \tag{8.1}\\ \left|\nabla_{r} w\right|=0, & \text { on } B_{1} \cap L\end{cases}
$$

where we recall that for $X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1} \cap L$, we set

$$
\left|\nabla_{r} w\right|\left(X_{0}\right):=\lim _{\left(x_{n}, z\right) \rightarrow(0,0)} \frac{w\left(x_{0}^{\prime}, x_{n}, z\right)-w\left(x_{0}^{\prime}, 0,0\right)}{r}, \quad r^{2}=x_{n}^{2}+z^{2} .
$$

We remark that if we restrict this linear problem to the class of functions $w(X)=\tilde{w}\left(x^{\prime}, r\right)$ that depend only on $\left(x^{\prime}, r\right)$ then the problem reduces to the classical Neumann problem

$$
\begin{cases}\Delta \tilde{w}=0, & \text { in } B_{1}^{+} \\ \tilde{w}_{r}=0, & \text { on }\{r=0\}\end{cases}
$$

The following is our main theorem.

Theorem 8.1. Given a boundary data $\bar{h} \in C\left(\partial B_{1}\right),|\bar{h}| \leqslant 1$, which is even with respect to $\{z=0\}$, there exists a unique classical solution $h$ to (8.1) such that $h \in C\left(\bar{B}_{1}\right), h=\bar{h}$ on $\partial B_{1}, h$ is even with respect to $\{z=0\}$ and it satisfies

$$
\begin{equation*}
\left|h(X)-h\left(X_{0}\right)-a^{\prime} \cdot\left(x^{\prime}-x_{0}^{\prime}\right)\right| \leqslant C\left(\left|x^{\prime}-x_{0}^{\prime}\right|^{2}+r^{3 / 2}\right), \quad X_{0} \in B_{1 / 2} \cap L \tag{8.2}
\end{equation*}
$$

for a universal constant $C$ and a vector $a^{\prime} \in \mathbb{R}^{n-1}$ depending on $X_{0}$.
As a corollary of the theorem above we obtain the desired regularity result, as stated also in Section 3.
Theorem 8.2 (Improvement of flatness). There exists a universal constant $C$ such that if $w$ is a viscosity solution to (8.1) in $B_{1}$ with

$$
-1 \leqslant w(X) \leqslant 1 \quad \text { in } B_{1}
$$

then

$$
\begin{equation*}
a_{0} \cdot x^{\prime}-C|X|^{3 / 2} \leqslant w(X)-w(0) \leqslant a_{0} \cdot x^{\prime}+C|X|^{3 / 2} \tag{8.3}
\end{equation*}
$$

for some vector $a_{0} \in \mathbb{R}^{n-1}$.
Proof. Let $h$ be the unique solution to (8.1) in $B_{1 / 2}$ with boundary data $w$. We will prove that $w=h$ in $B_{1 / 2}$ and hence it satisfies the desired estimate in view of (8.2). Denote by

$$
\bar{h}_{\epsilon}:=h-\epsilon+\epsilon^{2} r .
$$

Then, for $\epsilon$ small

$$
\bar{h}_{\epsilon}<w \quad \text { on } \partial B_{1 / 2}
$$

We wish to prove that

$$
\begin{equation*}
\bar{h}_{\epsilon} \leqslant w \quad \text { in } B_{1 / 2} \tag{8.4}
\end{equation*}
$$

Now, notice that $\bar{h}_{\epsilon}$ (and all its translations) is a classical strict subsolution to (8.1) that is

$$
\begin{cases}\Delta\left(U_{n} \bar{h}_{\epsilon}\right)=0, & \text { in } B_{1 / 2} \backslash P  \tag{8.5}\\ \left|\nabla_{r} \bar{h}_{\epsilon}\right|>0, & \text { on } B_{1 / 2} \cap L\end{cases}
$$

Since $w$ is bounded, for $t$ large enough $\bar{h}_{\epsilon}-t$ lies strictly below $w$. We let $t \rightarrow 0$ and show that the first contact point cannot occur for $t \geqslant 0$. Indeed since $\bar{h}_{\epsilon}-t$ is a strict subsolution which is strictly below $w$ on $\partial B_{1 / 2}$ then no touching can occur either in $B_{1 / 2} \backslash P$ or on $B_{1 / 2} \cap L$. We only need to check that no touching occurs on $P \backslash L$. This follows from Lemma 4.3.

Thus (8.4) holds. Passing to the limit as $\epsilon \rightarrow 0$ we get that

$$
h \leqslant w \quad \text { in } B_{1 / 2}
$$

Similarly we also obtain that

$$
h \geqslant w \quad \text { in } B_{1 / 2}
$$

and the desired equality holds.
The existence of the classical solution of Theorem 8.1 will be achieved via a variational approach in the appropriate weighted Sobolev space.

We say that $h \in H^{1}\left(U_{n}^{2} d X, B_{1}\right)$ is a minimizer to the energy functional

$$
J(h):=\int_{B_{1}} U_{n}^{2}|\nabla h|^{2} d X
$$

if

$$
J(h) \leqslant J(h+\phi), \quad \forall \phi \in C_{0}^{\infty}\left(B_{1}\right)
$$

Since $J$ is strictly convex this is equivalent to

$$
\lim _{\epsilon \rightarrow 0} \frac{J(h)-J(h+\epsilon \phi)}{\epsilon}=0, \quad \forall \phi \in C_{0}^{\infty}\left(B_{1}\right),
$$

which is satisfied if and only if

$$
\int_{B_{1}} U_{n}^{2} \nabla h \cdot \nabla \phi d X=0, \quad \forall \phi \in C_{0}^{\infty}\left(B_{1}\right) .
$$

As remarked above, if we restrict to the space of functions $h$ which are axisymmetric with respect to $L$ then the energy above reduces to the Dirichlet energy.

We start with a few standard facts about minimizers of $J$. First, $h$ solves the equation

$$
\operatorname{div}\left(U_{n}^{2} \nabla h\right)=0 \quad \text { in } B_{1},
$$

which is uniformly elliptic in any compact subset of $B_{1} \backslash P$ where $U_{n}$ is bounded. In particular $h \in C^{\infty}\left(B_{1} \backslash P\right)$, and we easily obtain the following lemma.

Lemma 8.3. Let h be a minimizer to $J$ in $B_{1}$, then

$$
\Delta\left(U_{n} h\right)=0 \quad \text { in } B_{1} \backslash P .
$$

Proof. Since $h$ is smooth in $B_{1} \backslash P$, from

$$
\operatorname{div}\left(U_{n}^{2} \nabla h\right)=0 \quad \text { in } B_{1},
$$

we obtain that

$$
U_{n}^{2} \Delta h+2 \sum_{i=1}^{n+1} U_{n} U_{n i} h_{i}=0 \quad \text { in } B_{1} \backslash P .
$$

Since $U_{n}>0$ and $\Delta U=0$ in $B_{1} \backslash P$ the identity above is equivalent to

$$
\Delta\left(U_{n} h\right)=U_{n} \Delta h+2 \nabla U_{n} \cdot \nabla h=0 \quad \text { in } B_{1} \backslash P,
$$

as desired.
The next lemma contains a characterization of minimizer, which will be useful later in this section.
Lemma 8.4. Let $h \in C\left(B_{1}\right)$ be a solution to

$$
\begin{equation*}
\Delta\left(U_{n} h\right)=0 \quad \text { in } B_{1} \backslash P, \tag{8.6}
\end{equation*}
$$

and assume that

$$
\lim _{r \rightarrow 0} h_{r}\left(x^{\prime}, x_{n}, z\right)=b\left(x^{\prime}\right) \text {, }
$$

with $b\left(x^{\prime}\right)$ a continuous function. Then $h$ is a minimizer to $J$ in $B_{1}$ if and only if $b \equiv 0$.
Proof. By integration by parts and the computation in Lemma 8.3 the identity

$$
\int_{B_{1}} U_{n}^{2} \nabla h \cdot \nabla \phi d X=0, \quad \forall \phi \in C_{0}^{\infty}\left(B_{1}\right),
$$

is equivalent to the following two conditions

$$
\begin{equation*}
\Delta\left(U_{n} h\right)=0 \quad \text { in } B_{1} \backslash P, \tag{8.7}
\end{equation*}
$$

and

$$
\begin{equation*}
\lim _{\delta \rightarrow 0} \int_{\partial C_{\delta} \cap B_{1}} U_{n}^{2} \phi \nabla h \cdot v d \sigma=0, \tag{8.8}
\end{equation*}
$$

where $C_{\delta}$ is the cylinder $\{r \leqslant \delta\}$ and $\nu$ the inward unit normal to $C_{\delta}$.

Here we use that

$$
\lim _{\epsilon \rightarrow 0} \int_{\{|z|=\epsilon\} \cap\left(B_{1} \backslash C_{\delta}\right)} U_{n}^{2} \phi h_{\nu} d \sigma=0 .
$$

Indeed, in the set $\{|z|=\epsilon\} \cap\left(B_{1} \backslash C_{\delta}\right)$ we have

$$
U_{n} \leqslant C \epsilon,
$$

and

$$
\left|\nabla\left(U_{n} h\right)\right|,\left|\nabla U_{n}\right| \leqslant C,
$$

from which it follows that

$$
U_{n}|\nabla h| \leqslant C .
$$

In conclusion we need to show that (8.8) is equivalent to $b\left(x^{\prime}\right)=0$.
Indeed,

$$
\begin{aligned}
\int_{\partial C_{\delta} \cap B_{1}} U_{n}^{2} \phi \nabla h \cdot v d \sigma & =\frac{1}{\delta} \int_{\partial C_{\delta} \cap B_{1}} \cos ^{2}\left(\frac{\theta}{2}\right) h_{r} \phi d \sigma \\
& =\int_{\partial C_{1} \cap B_{1}} \cos ^{2}\left(\frac{\theta}{2}\right)\left(h_{r} \phi\right)\left(X^{\prime}, \delta \cos \theta, \delta \sin \theta\right) d x^{\prime} d \theta
\end{aligned}
$$

hence

$$
\lim _{\delta \rightarrow 0} \int_{\partial C_{\delta} \cap B_{1}} U_{n}^{2} \phi \nabla h \cdot v d \sigma=\pi \int_{L} b\left(x^{\prime}\right) \phi\left(x^{\prime}, 0,0\right) d x^{\prime}
$$

and our claim clearly follows.
The next lemma follows by standard arguments, hence we omit its proof.
Lemma 8.5 (Comparison Principle). Let $h_{1}, h_{2}$ be minimizers to $J$ in $B_{1}$. If

$$
h_{1} \geqslant h_{2} \quad \text { a.e in } B_{1} \backslash B_{\rho},
$$

then

$$
h_{1} \geqslant h_{2} \quad \text { a.e. in } B_{1} .
$$

Finally one of the main ingredients in the proof of Theorem 8.1 is the following Harnack inequality.
Lemma 8.6 (Harnack inequality). Let $h$ be a minimizer to $J$ in $B_{1}$ which is even with respect to $\{z=0\}$. Then $h \in C^{\alpha}\left(B_{1 / 2}\right)$ and
$[h]_{C^{\alpha}\left(B_{1 / 2}\right)} \leqslant C$,
with $C$ universal.
The proof of this lemma follows the same lines as the proof of Harnack inequality (Theorem 6.1) for the free boundary problem (2.1). We briefly sketch it in what follows.

Sketch of the proof of Lemma 8.6. The key step consists in proving the following claim, which plays the same role as Lemma 6.3 in the proof of Theorem 6.1. The remaining ingredients are the standard Harnack inequality and boundary Harnack inequality for harmonic functions.

Claim. There exist universal constants $\delta, c$ such that if $h \geqslant 0$ a.e. in $B_{1}$ and

$$
h\left(\frac{1}{4} e_{n}\right) \geqslant 1
$$

then

$$
h \geqslant c \quad \text { a.e. in } B_{\delta} .
$$

As in the proof of Lemma 6.3 , since minimizers satisfy the comparison principle Lemma 8.5 , the claim will follow if we provide the right family of comparison minimizers. This family plays the same role as the $v_{R}$ 's in Lemma 6.3 and it is obtained by translations and multiplication by constants of the following function

$$
v(X):=-\frac{\left|x^{\prime}\right|^{2}}{n-1}+2 x_{n} r .
$$

We need to show that $v$ is a minimizer to $J$ in $B_{1}$. To do so we prove that $v$ satisfies Lemma 8.4.
To prove that

$$
\Delta\left(U_{n} v\right)=0 \quad \text { in } B_{1} \backslash P,
$$

we use that $2 r U_{n}=U$ and that $U, U_{n}$ are harmonic outside of $P$ and do not depend on $x^{\prime}$. Thus

$$
\Delta\left(U_{n} v\right)=-\Delta\left(\frac{\left|x^{\prime}\right|^{2}}{n-1} U_{n}\right)+\Delta\left(x_{n} U\right)=-2 U_{n}+2 U_{n}=0
$$

Finally the fact that

$$
\lim _{r \rightarrow 0} v_{r}\left(x^{\prime}, x_{n}, z\right)=0
$$

follows immediately from the definition of $v$.
Since our linear problem is invariant under translations in the $x^{\prime}$-direction, we see that discrete differences of the form

$$
h(X+\tau)-h(X),
$$

with $\tau$ in the $x^{\prime}$-direction are also minimizers. Now by standard arguments (see [5]) we obtain the following corollary.
Corollary 8.7. Let h be a minimizer to $J$ in $B_{1}$ which is even with respect to $\{z=0\}$. Then $D_{x^{\prime}}^{\beta} h \in C^{\alpha}\left(B_{1 / 2}\right)$ and

$$
\left[D_{x^{\prime}}^{\beta},\right]_{C^{\alpha}\left(B_{1 / 2}\right)} \leqslant C
$$

with $C$ depending on $\beta$.
We are now ready to prove our main theorem.
Proof of Theorem 8.1. We divide our proof in several steps.
Step 1. In this step, we show the existence of a classical solution to our problem, which achieves the boundary data continuously.

Assume without loss of generality that $\bar{h} \in C^{\infty}\left(\partial B_{1}\right)$. The general case when $\bar{h} \in C\left(\partial B_{1}\right)$ follows by approximation and the comparison principle.

We minimize $J(\cdot)$ among all functions $h$ with boundary data $\bar{h}$, which are even with respect to $\{z=0\}$. From Lemma 8.3 we have that

$$
\Delta\left(U_{n} h\right)=0 \quad \text { in } B_{1} \backslash P .
$$

In Steps 2-3 we will show that $h$ satisfies the estimate (8.2) and in particular

$$
h\left(x_{0}^{\prime}, x_{n}, x\right)-h\left(x_{0}^{\prime}, 0,0\right)=O\left(r^{3 / 2}\right), \quad X_{0}=\left(x_{0}^{\prime}, 0,0\right) \in B_{1} \cap L
$$

which gives

$$
\left|\nabla_{r} h\right|=0 \quad \text { on } B_{1} \cap L
$$

where we recall that

$$
\left|\nabla_{r} h\right|\left(X_{0}\right)=\lim _{\left(x_{n}, z\right) \rightarrow(0,0)} \frac{h\left(x_{0}^{\prime}, x_{n}, z\right)-h\left(x_{0}^{\prime}, 0,0\right)}{r}, \quad r^{2}=x_{n}^{2}+z^{2} .
$$

Notice that since $|\bar{h}| \leqslant 1$, also $|h| \leqslant 1$ and $h, D_{x^{\prime}} h \in C^{0, \alpha}\left(B_{1}\right)$ in view of Lemma 8.6 and its corollary.
We now show that $h$ achieves the boundary data $\bar{h}$ continuously. Indeed this follows by classical elliptic theory if we restrict to $\partial B_{1} \backslash P$.

If $X_{0} \in \partial B_{1} \cap(P \backslash L)$ then in a small neighborhood of $X_{0}$ intersected with $B_{1} \cap\{z>0\}$ the function $U_{n} h$ is harmonic continuous up to the boundary and vanishes continuously on $\{z=0\}$ (since $h$ is bounded). The continuity of $h$ at $X_{0}$ then follows from standard boundary regularity for the harmonic function $U_{n} h$.

Finally, on the set $\partial B_{1} \cap L$ as in the case of Laplace equation, it suffices to construct at each point $X_{0}$ a local barrier (minimizer) for $h$ which is zero at $X_{0}$ and strictly negative in a neighborhood of $X_{0}$. Such a barrier is given by (see Lemma 8.4)

$$
\left(x^{\prime}-x_{0}^{\prime}\right) \cdot x_{0}^{\prime} .
$$

Step 2. In this step we wish to prove that

$$
\begin{align*}
& \left|h\left(x^{\prime}, x_{n}, x\right)-h\left(x^{\prime}, 0,0\right)-b\left(x^{\prime}\right) r\right| \leqslant C r^{3 / 2}, \quad\left(x^{\prime}, 0,0\right) \in B_{1 / 2} \cap L  \tag{8.9}\\
& \left|h_{r}\left(x^{\prime}, x_{n}, z\right)-b\left(x^{\prime}\right)\right| \leqslant C r^{1 / 2}, \quad\left(x^{\prime}, 0,0\right) \in B_{1 / 2} \cap L \tag{8.10}
\end{align*}
$$

with $C$ universal and $b\left(x^{\prime}\right)$ a Lipschitz function.
Indeed, $h$ solves

$$
\Delta\left(U_{n} h\right)=0 \quad \text { in } B_{1} \backslash P .
$$

Since $U_{n}$ is independent on $x^{\prime}$ we can rewrite this equation as

$$
\begin{equation*}
\Delta_{x_{n}, z}\left(U_{n} h\right)=-U_{n} \Delta_{x^{\prime}} h, \tag{8.11}
\end{equation*}
$$

and according to Lemma 8.6 we have that

$$
\Delta_{x^{\prime}} h \in C^{\alpha}\left(B_{1 / 2}\right),
$$

with universal bound. Thus, for each fixed $x^{\prime}$, we need to investigate the 2-dimensional problem

$$
\Delta\left(U_{t} h\right)=U_{t} f, \quad \text { in } B_{1 / 2} \backslash\{t \leqslant 0, z=0\}
$$

with

$$
f \in C^{\alpha}\left(B_{1 / 2}\right),
$$

and $h$ bounded. Without loss of generality, for a fixed $x^{\prime}$ we may assume $h\left(x^{\prime}, 0,0\right)=0$.
Let $H(t, z)$ be the solution to the problem

$$
\Delta H=U_{t} f, \quad \text { in } B_{1 / 2} \backslash\{t \leqslant 0, z=0\},
$$

such that

$$
H=U_{t} h \quad \text { on } \partial B_{1 / 2}, \quad H=0 \quad \text { on } B_{1 / 2} \cap\{t \leqslant 0, z=0\} .
$$

We wish to prove that

$$
\begin{equation*}
U_{t} h=H \tag{8.12}
\end{equation*}
$$

First notice that

$$
\Delta\left(H-U_{t} h\right)=0 \quad \text { in } B_{1 / 2} \backslash\{t \leqslant 0, z=0\},
$$

and

$$
H=0 \quad \text { on } \partial B_{1 / 2} \cup\left(B_{1 / 2} \cap\{t<0, z=0\}\right) .
$$

We claim that

$$
\begin{equation*}
\lim _{(t, z) \rightarrow(0,0)} \frac{H-U_{t} h}{U_{t}}=\lim _{(t, z) \rightarrow(0,0)} \frac{H}{U_{t}}=0 \tag{8.13}
\end{equation*}
$$

If the claim holds, then given any $\epsilon>0$

$$
-\epsilon U_{t} \leqslant H-U_{t} h \leqslant \epsilon U_{t}, \quad \text { in } B_{\delta}
$$

with $\delta=\delta(\epsilon)$. Then by the maximum principle the inequality above holds in the whole $B_{1 / 2}$ and by letting $\epsilon \rightarrow 0$ we obtain (8.12).

To prove the claim (8.13) we show that $H$ satisfies the following

$$
\begin{equation*}
|H(t, z)-a U(t, z)| \leqslant C_{0} r^{1 / 2} U(t, z), \quad r^{2}=t^{2}+z^{2}, a \in \mathbb{R} \tag{8.14}
\end{equation*}
$$

with $C_{0}$ universal.
To do so, we consider the holomorphic transformation

$$
\Phi:(s, y) \rightarrow(t, z)=\left(\frac{1}{2}\left(s^{2}-y^{2}\right), s y\right)
$$

which maps $B_{1} \cap\{s>0\}$ into $B_{1 / 2} \backslash\{t \leqslant 0, z=0\}$ and call

$$
\widetilde{H}(s, y)=H(t, z), \quad \tilde{f}(s, y)=f(t, z) .
$$

Then, easy computations show that

$$
\Delta \widetilde{H}=s \tilde{f} \quad \text { in } B_{1} \cap\{s>0\}, \quad \widetilde{H}=0 \quad \text { on } B_{1} \cap\{s=0\} .
$$

Since the right-hand side is $C^{\alpha}$ we conclude that $\tilde{H} \in C^{2, \alpha}$. In particular $\widetilde{H}_{s}$ satisfies

$$
\left|\widetilde{H}_{s}(s, y)-a\right| \leqslant C_{0}|(s, y)|, \quad a=\widetilde{H}_{s}(0,0)
$$

with $C_{0}$ universal. Integrating this inequality between 0 and $s$ and using that $\widetilde{H}=0$ on $B_{1} \cap\{s=0\}$ we get that

$$
|\widetilde{H}(s, y)-a s| \leqslant C_{0} s|(s, y)| .
$$

In terms of $H$, this equation gives us

$$
\begin{equation*}
|H-a U| \leqslant C_{0} r^{1 / 2} U \tag{8.15}
\end{equation*}
$$

as desired.
Thus (8.12) and (8.14) hold and by combining them and using that $U / U_{t}=2 r$ we get that

$$
|h-2 a r| \leqslant 2 C_{0} r^{3 / 2},
$$

which is the desired estimate (8.9) i.e. (recall that above we assumed $h\left(x^{\prime}, 0,0\right)=0$ )

$$
\left|h\left(x^{\prime}, x_{n}, x\right)-h\left(x^{\prime}, 0,0\right)-b\left(x^{\prime}\right) r\right| \leqslant 2 C_{0} r^{3 / 2}
$$

with

$$
b\left(x^{\prime}\right)=2 \widetilde{H}_{s}\left(x^{\prime}, 0,0\right)
$$

We remark that $b\left(x^{\prime}\right)$ is Lipschitz. Indeed, notice that the derivatives $h_{i}, i=1, \ldots, n-1$, still satisfy the same Eq. (8.11) as $h$, where the $C^{\alpha}$-norm of the right-hand side has a universal bound. Thus, we can argue as above to conclude that

$$
\left|\partial_{i} \widetilde{H}_{s}\left(x^{\prime}, 0,0\right)\right| \leqslant C,
$$

which together with the formula for $b\left(x^{\prime}\right)$ shows that $b\left(x^{\prime}\right)$ is a Lipschitz function.
Finally, to obtain the second of our estimates (8.10) we proceed similarly as above.

Since $U_{t} h=H$ one can compute easily that

$$
\begin{equation*}
U_{t} h_{r}=H_{r}+\frac{1}{2} \frac{H}{r} . \tag{8.16}
\end{equation*}
$$

Moreover, after our holomorphic transformation

$$
\begin{equation*}
2 r H_{r}(t, z)=s \widetilde{H}_{s}(s, y)+y \widetilde{H}_{y}(s, y) . \tag{8.17}
\end{equation*}
$$

As observed above,

$$
\left|\widetilde{H}_{s}-a\left(x^{\prime}\right)\right| \leqslant C|(s, y)|,
$$

and similarly since $\widetilde{H}=0$ on $B_{1} \cap\{s=0\}$

$$
\left|\widetilde{H}_{y}\right| \leqslant C s
$$

These two inequalities combined with (8.17) give us

$$
\begin{equation*}
\left|2 r H_{r}-a\left(x^{\prime}\right) U\right| \leqslant C r^{1 / 2} U . \tag{8.18}
\end{equation*}
$$

Combining (8.16) with (8.15)-(8.18) we obtain (8.10) as desired.
Step 3. In this step we show that $h$ satisfies (8.2).
In view of Lemma 8.4 and estimate (8.10) we obtain that $b\left(x^{\prime}\right)=0$. Since $b\left(x^{\prime}\right)=0$ then (8.9) reads

$$
\left|h\left(x^{\prime}, x_{n}, z\right)-h\left(x^{\prime}, 0,0\right)\right| \leqslant C r^{3 / 2} .
$$

Since $h$ is $C^{\infty}$ in the $x^{\prime}$-direction, we have that for a given $X_{0} \in B_{1 / 4} \cap L$

$$
\left|h\left(x^{\prime}, 0,0\right)-h\left(x_{0}^{\prime}, 0,0\right)-a^{\prime} \cdot\left(x^{\prime}-x_{0}^{\prime}\right)\right| \leqslant C\left|x^{\prime}-x_{0}^{\prime}\right|^{2}
$$

which combined with the previous inequality gives us the desired bound (8.2).

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## Appendix A

The purpose of this section is to provide the proof of Proposition 6.4 and Corollary 6.5 which have been used in the proof of Harnack inequality in Section 6. For the reader's convenience we recall their statements.

Let $R>0$ and denote

$$
V_{R}(t, z)=U(t, z)\left((n-1) \frac{t}{R}+1\right) .
$$

Then set

$$
v_{R}(X)=V_{R}\left(R-\sqrt{\left|x^{\prime}\right|^{2}+\left(x_{n}-R\right)^{2}}, z\right)
$$

Proposition 6.4. If $R$ is large enough, the function $v_{R}(X)$ is a comparison subsolution to (2.1) in $B_{2}$ which is strictly monotone increasing in the $e_{n}$-direction in $B_{2}^{+}\left(v_{R}\right)$. Moreover, there exists a function $\tilde{v}_{R}$ such that

$$
\begin{equation*}
U(X)=v_{R}\left(X-\tilde{v}_{R}(X) e_{n}\right), \quad \text { in } B_{1} \backslash P \tag{A.1}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\tilde{v}_{R}(X)-\gamma_{R}(X)\right| \leqslant \frac{C}{R^{2}}|X|^{2}, \quad \gamma_{R}(X)=-\frac{\left|x^{\prime}\right|^{2}}{2 R}+2(n-1) \frac{x_{n} r}{R} \tag{A.2}
\end{equation*}
$$

with $r=\sqrt{x_{n}^{2}+z^{2}}$ and $C$ universal.

Proof. We divide the proof of this proposition in two steps.
Step 1. In this step we show that $v_{R}$ is a comparison subsolution in $B_{2}$ which is monotone in the $e_{n}$-direction.
First we need to show that $v_{R}$ is subharmonic (but not harmonic) in $B_{2}^{+}\left(v_{R}\right)$. From the formula for $v_{R}$ we see immediately that $(R>2)$

$$
B_{2}^{+}\left(v_{R}\right)=B_{2} \backslash\left(\mathcal{B}_{2} \backslash \overline{\mathcal{B}}_{R}\left(R e_{n}\right)\right) .
$$

One can easily compute that on such set,

$$
\Delta v_{R}(X)=\left(\left(\partial_{t t}+\partial_{z z}\right) V_{R}\right)(R-\rho, z)-\frac{n-1}{\rho} \partial_{t} V_{R}(R-\rho, z),
$$

where for simplicity we call

$$
\rho:=\sqrt{\left|x^{\prime}\right|^{2}+\left(x_{n}-R\right)^{2}} .
$$

Also for $(t, z)$ outside the set $\{(t, 0): t \leqslant 0\}$

$$
\begin{aligned}
\Delta_{t, z} V_{R}(t, z) & =\left(\partial_{t t}+\partial_{z z}\right) V_{R}(t, z)=\frac{2(n-1)}{R} \partial_{t} U(t, z)+\left(1+(n-1) \frac{t}{R}\right) \Delta_{t, z} U(t, z) \\
& =\frac{2(n-1)}{R} \partial_{t} U(t, z)
\end{aligned}
$$

and

$$
\begin{equation*}
\partial_{t} V_{R}(t, z)=\left(1+(n-1) \frac{t}{R}\right) \partial_{t} U(t, z)+\frac{n-1}{R} U(t, z) . \tag{A.3}
\end{equation*}
$$

Thus to show that $\Delta v_{R}$ is subharmonic in $B_{2}^{+}\left(v_{R}\right)$ we need to prove that in such set

$$
\frac{2(n-1)}{R} \partial_{t} U-\frac{n-1}{\rho}\left[\left(1+(n-1) \frac{R-\rho}{R}\right) \partial_{t} U+\frac{n-1}{R} U\right] \geqslant 0,
$$

where $U$ and $\partial_{t} U$ are evaluated at $(R-\rho, z)$.
Set $t=R-\rho$, then straightforward computations reduce the inequality above to

$$
(n-1)\left[2(R-t)-R-(n-1)^{2} t\right] \partial_{t} U(t, z)-(n-1)^{2} U(t, z) \geqslant 0 .
$$

Using that $\partial_{t} U(t, z)=U(t, z) /(2 r)$ with $r^{2}=t^{2}+z^{2}$, this inequality becomes

$$
R \geqslant 2 t+(n-1)^{2} t+2(n-1) r .
$$

This last inequality is easily satisfied for $R$ large enough, since $t, r \leqslant 3$.
Now we prove that $v_{R}$ satisfies the free boundary condition in Definition 2.2. First observe that

$$
F\left(v_{R}\right)=\partial \mathcal{B}_{R}\left(R e_{n}, 0\right) \cap \mathcal{B}_{2},
$$

and hence it is smooth. By the radial symmetry it is enough to show that the free boundary condition is satisfied at $0 \in F\left(v_{R}\right)$ that is

$$
\begin{equation*}
v_{R}(x, z)=\alpha U\left(x_{n}, z\right)+o\left(|(x, z)|^{1 / 2}\right), \quad \text { as }(x, z) \rightarrow(0,0) \tag{A.4}
\end{equation*}
$$

with $\alpha \geqslant 1$.
First notice since $U$ is Holder continuous with exponent $1 / 2$, it follows from the formula for $V_{R}$ that

$$
\left|V_{R}(t, z)-V_{R}\left(t_{0}, z\right)\right| \leqslant C\left|t-t_{0}\right|^{1 / 2} \quad \text { for }\left|t-t_{0}\right| \leqslant 1
$$

Thus for $(x, z) \in B_{s}, s$ small

$$
\left|v_{R}(x, z)-V_{R}\left(x_{n}, z\right)\right|=\left|V_{R}(R-\rho, z)-V_{R}\left(x_{n}, z\right)\right| \leqslant C\left|R-\rho-x_{n}\right|^{1 / 2} \leqslant C s,
$$

where we have used that (recall that $\rho:=\sqrt{\left|x^{\prime}\right|^{2}+\left(x_{n}-R\right)^{2}}$ )

$$
\begin{equation*}
R-\rho-x_{n}=-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\rho} \tag{A.5}
\end{equation*}
$$

It follows that for $(x, z) \in B_{s}$

$$
\begin{aligned}
\left|v_{R}(x, z)-U\left(x_{n}, z\right)\right| & \leqslant\left|v_{R}(x, z)-V_{R}\left(x_{n}, z\right)\right|+\left|V_{R}\left(x_{n}, z\right)-U\left(x_{n}, z\right)\right| \\
& \leqslant C s+\left|V_{R}\left(x_{n}, z\right)-U\left(x_{n}, z\right)\right| .
\end{aligned}
$$

Thus from the formula for $V_{R}$

$$
\left|v_{R}(x, z)-U\left(x_{n}, z\right)\right| \leqslant C s+(n-1) \frac{\left|x_{n}\right|}{R} U\left(x_{n}, z\right) \leqslant C^{\prime} s, \quad(x, z) \in B_{s}
$$

which gives the desired expansion (A.4) with $\alpha=1$.
Now, we show that $v_{R}$ is strictly monotone increasing in the $e_{n}$-direction in $B_{2}^{+}\left(v_{R}\right)$. Outside of its zero plate,

$$
\partial_{x_{n}} v_{R}(x)=-\frac{x_{n}-R}{\rho} \partial_{t} V_{R}(R-\rho, z) .
$$

Thus we only need to show that $V_{R}(t, z)$ is strictly monotone increasing in $t$ outside $\{(t, 0): t \leqslant 0\}$. This follows immediately from (A.3) and the formula for $U$.

Step 2. In this step we show the existence of $\tilde{v}_{R}$ satisfying (A.1) and (A.2). Since we have a precise formula for $v_{R}$ in terms of $U$, this is only a matter of straightforward (though tedious) computations which we present here for completeness.

First we show that there exists a unique $\tilde{t}$ such that (here $B_{1}$ is the 2 -dimensional ball)

$$
\begin{equation*}
U(t, z)=V_{R}(t+\tilde{t}, z), \quad \text { in } B_{1} \backslash\{(t, 0): t \leqslant 0\}, \tag{A.6}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\tilde{t}+\frac{2(n-1) t r}{R}\right| \leqslant \frac{\widetilde{C}}{R^{2}} r^{3}, \quad r^{2}=t^{2}+z^{2} \tag{A.7}
\end{equation*}
$$

with $\widetilde{C}$ universal. Since $V_{R}$ is strictly increasing in the $t$-direction in $B_{1} \backslash\{(t, 0): t \leqslant 0\}$ it suffices to show that

$$
\begin{equation*}
V_{R}\left(t-\frac{2(n-1) t r}{R}-\frac{\widetilde{C}}{R^{2}} r^{3}\right)<U(t, z)<V_{R}\left(t-\frac{2(n-1) t r}{R}+\frac{\widetilde{C}}{R^{2}} r^{3}\right) \tag{A.8}
\end{equation*}
$$

Let us prove the lower bound. We call

$$
\bar{t}=-\frac{2(n-1) t r}{R}-\frac{\widetilde{C}}{R^{2}} r^{3},
$$

and we use Taylor's formula to compute

$$
\begin{equation*}
V_{R}(t+\bar{t}, z)=V_{R}(t, z)+\partial_{t} V_{R}(t, z) \bar{t}+\frac{1}{2} E \bar{t}^{2}, \quad|E| \leqslant\left|\partial_{t t} V_{R}(s, z)\right|, \tag{A.9}
\end{equation*}
$$

with $s$ between $t$ and $t+\bar{t}$. We claim that

$$
\begin{equation*}
\left|\partial_{t t} V_{R}(s, z)\right| \leqslant \frac{C^{\prime}}{r^{2}} U(t, z) \tag{A.10}
\end{equation*}
$$

Indeed one can compute that

$$
\begin{align*}
\partial_{t t} V_{R}(s, z) & =\left(1+(n-1) \frac{s}{R}\right) \partial_{t t} U(s, z)+\frac{2(n-1)}{R} \partial_{t} U(s, z) \\
& =\left(1+(n-1) \frac{s}{R}\right) r^{-3 / 2} \partial_{t t} U\left(\frac{s}{r}, \frac{z}{r}\right)+r^{-1 / 2} \frac{2(n-1)}{R} \partial_{t} U\left(\frac{s}{r}, \frac{z}{r}\right) \tag{A.11}
\end{align*}
$$

where we have used that $U$ is homogeneous of degree $1 / 2$.
Since $s$ lies between $t$ and $t+\bar{t}$ we get that $(s / r, z / r) \in B_{3 / 2} \backslash B_{1 / 2}$, thus by boundary Harnack inequality in this annulus we get

$$
\left|\partial_{t t} U\left(\frac{s}{r}, \frac{z}{r}\right)\right| \leqslant K_{1} U\left(\frac{s}{r}, \frac{z}{r}\right), \quad \partial_{t} U\left(\frac{s}{r}, \frac{z}{r}\right) \leqslant K_{2} U\left(\frac{s}{r}, \frac{z}{r}\right),
$$

with $K_{1}, K_{2}$ universal.

Combining the above inequalities with (A.11) we obtain that

$$
\left|\partial_{t t} V_{R}(s, z)\right| \leqslant \bar{C} r^{-3 / 2} U\left(\frac{s}{r}, \frac{z}{r}\right)
$$

Thus the claim in (A.10) will follow if we show that for some $K$ universal

$$
U\left(\frac{s}{r}, \frac{z}{r}\right) \leqslant K U\left(\frac{t}{r}, \frac{z}{r}\right)
$$

Again, this follows by the boundary Harnack inequality in the annulus $B_{3 / 2} \backslash B_{1 / 2}$ between $U\left(\frac{t}{r}, \frac{z}{r}\right)$ and its translation $U\left(\frac{t+\tau}{r}, \frac{z}{r}\right)$, for $\tau$ small. Our claim is thus proved.

Thus, using (A.9) together with this claim, the lower bound in (A.8) will be proved if we show that

$$
U(t, z)>V_{R}(t, z)+\partial_{t} V_{R}(t, z) \bar{t}+\frac{C^{\prime}}{2 r^{2}} U(t, z) \bar{t}^{2}
$$

From the definition of $V_{R}$ this is equivalent to showing that

$$
(n-1) \frac{t}{R} U(t, z)+\left[\left(1+(n-1) \frac{t}{R}\right) \partial_{t} U(t, z)+\frac{n-1}{R} U(t, z)\right] \bar{t}+\frac{C^{\prime}}{2 r^{2}} U(t, z) \bar{t}^{2}<0 .
$$

Dividing everything by $\partial_{t} U(t, z)=\frac{1}{2 r} U(t, z)$ we get

$$
\frac{2(n-1) r t}{R}+\left[\left(1+(n-1) \frac{t}{R}\right)+\frac{2 r(n-1)}{R}\right] \bar{t}+\frac{C^{\prime}}{r} \bar{t}^{2}<0,
$$

and using the definition of $\bar{t}$ we finally need to show that

$$
(n-1) \frac{t+2 r}{R} \bar{t}+\frac{C^{\prime}}{r} \bar{t}^{2}<\frac{\widetilde{C}}{R^{2}} r^{3} .
$$

Using that for $R$ large

$$
|\overline{\bar{t}}| \leqslant K r^{2} / R
$$

for $K$ universal, we easily obtain that the inequality above holds for the appropriate $\widetilde{C}$ (universal) and $R$ large.
To conclude our proof, we use (A.5) to write

$$
v_{R}\left(X-\tilde{v}_{R} e_{n}\right)=V_{R}\left(R-\rho\left(\tilde{v}_{R}\right), z\right)=V_{R}\left(x_{n}-\tilde{v}_{R}-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\tilde{v}_{R}+\rho\left(\tilde{v}_{R}\right)}, z\right),
$$

with

$$
\rho(\eta):=\sqrt{\left|x^{\prime}\right|^{2}+\left(x_{n}-\eta-R\right)^{2}} .
$$

In view of (A.6) if there exists $\tilde{v}_{R}=\tilde{v}_{R}(X)$ such that

$$
\begin{equation*}
-\tilde{v}_{R}-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\tilde{v}_{R}+\rho\left(\tilde{v}_{R}\right)}=\tilde{t} \tag{A.12}
\end{equation*}
$$

then

$$
v_{R}\left(X-\tilde{v}_{R} e_{n}\right)=U(X), \quad \text { in } B_{1}^{+}(U)
$$

By the strict monotonicity of $v_{R}$ in the $e_{n}$-direction in $B_{1}^{+}\left(v_{R}\right)$, in such set $\tilde{v}_{R}$ must be unique.
Thus our claim will follow if we show that there exists $\tilde{v}_{R}$ satisfying (A.12) and such that

$$
\left|\tilde{v}_{R}(X)-\gamma_{R}(X)\right| \leqslant C \frac{|X|^{2}}{R^{2}}
$$

To do so, we call

$$
f(\eta)=-\eta-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\eta+\rho(\eta)}, \quad-1 \leqslant \eta \leqslant 1,
$$

and we show that

$$
f\left(\gamma_{R}(X)+C \frac{|X|^{2}}{R^{2}}\right) \leqslant \tilde{t} \leqslant f\left(\gamma_{R}(X)-C \frac{|x|^{2}}{R^{2}}\right) .
$$

In view of (A.7) we need to prove that

$$
f\left(\gamma_{R}(X)+C \frac{|X|^{2}}{R^{2}}\right) \leqslant-\frac{2(n-1) x_{n} r}{R}-\widetilde{C} \frac{r^{3}}{R^{2}},
$$

and

$$
f\left(\gamma_{R}(X)-C \frac{|X|^{2}}{R^{2}}\right) \geqslant-\frac{2(n-1) x_{n} r}{R}+\widetilde{C} \frac{r^{3}}{R^{2}} .
$$

Let us prove the first inequality (the second one follows similarly). Call

$$
\bar{\eta}=\gamma_{R}(X)+C \frac{|X|^{2}}{R^{2}} .
$$

From the definition of $f$ and $\gamma_{R}$ the desired inequality is equivalent to

$$
\frac{\left|x^{\prime}\right|^{2}}{2 R}-C \frac{|X|^{2}}{R^{2}}-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\bar{\eta}+\rho(\bar{\eta})} \leqslant-\widetilde{C} \frac{r^{3}}{R^{2}} .
$$

Clearly $-1 \leqslant \bar{\eta} \leqslant 1$, and one can easily verify that

$$
R-x_{n}+\bar{\eta}+\rho(\bar{\eta}) \leqslant 2 R+5 .
$$

Thus

$$
\frac{\left|x^{\prime}\right|^{2}}{2 R}-\frac{\left|x^{\prime}\right|^{2}}{R-x_{n}+\bar{\eta}+\rho(\bar{\eta})} \leqslant\left|x^{\prime}\right|^{2}\left(\frac{1}{2 R}-\frac{1}{2 R+5}\right) \leqslant \frac{\left|x^{\prime}\right|^{2}}{R^{2}},
$$

and the desired inequality follows if we show that

$$
\frac{\left|x^{\prime}\right|^{2}}{R^{2}}-C \frac{|X|^{2}}{R^{2}} \leqslant-\widetilde{C} \frac{r^{3}}{R^{2}} .
$$

This inequality is trivially satisfied as long as $C-\widetilde{C} \geqslant 1$.
We now recall the statement of Corollary 6.5 and sketch its proof.
Corollary 6.5. There exist $\delta, c_{0}, C_{0}, C_{1}$ universal constants, such that

$$
\begin{equation*}
v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \leqslant\left(1+\frac{C_{0}}{R}\right) U(X), \quad \text { in } \bar{B}_{1} \backslash B_{1 / 4}, \tag{A.13}
\end{equation*}
$$

with strict inequality on $F\left(v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right)\right) \cap \bar{B}_{1} \backslash B_{1 / 4}$,

$$
\begin{align*}
& v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \geqslant U\left(X+\frac{c_{0}}{2 R} e_{n}\right), \quad \text { in } B_{\delta},  \tag{A.14}\\
& v_{R}\left(X-\frac{C_{1}}{R} e_{n}\right) \leqslant U(X), \quad \text { in } \bar{B}_{1} . \tag{A.15}
\end{align*}
$$

Proof. Estimates (A.14) and (A.15) are immediate consequences of (A.2) and Lemma 3.1.
To obtain (A.13), notice that in view of (A.2) and Lemma 3.1,

$$
v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \leqslant U(X) \quad \text { in }\left\{X \in B_{1}:\left|x^{\prime}\right| \geqslant 1 / 8,\left|x_{n}\right| \leqslant \bar{\delta}\right\},
$$

for some $c_{0}, \bar{\delta}$ small universal and $R$ large (with strict inequality on $F\left(v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right)\right)$ ). Hence the estimate (A.13) holds on the set $\left\{X \in B_{1} \backslash B_{1 / 4}: \sqrt{x_{n}^{2}+z^{2}} \leqslant \bar{\delta}\right\}$ and we only need to prove it on the complement of this set.

Again, from (A.2) and Lemma 3.1 we get that

$$
\begin{equation*}
v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \leqslant U\left(X+\frac{\bar{C}}{R} e_{n}\right) \quad \text { in } B_{1}, \tag{A.16}
\end{equation*}
$$

for $\bar{C}$ large universal. From Lemma 5.2 we know that

$$
U\left(x_{n}+\frac{\bar{C}}{R}, z\right) \leqslant\left(1+C \frac{\bar{C}}{R}\right) U\left(x_{n}, z\right),
$$

as long as $\sqrt{x_{n}^{2}+z^{2}}>\bar{\delta}$, with $C=C(\bar{\delta})$ (and $R$ large). Combining this fact with (A.16) we get

$$
v_{R}\left(X+\frac{c_{0}}{R} e_{n}\right) \leqslant\left(1+\frac{C_{0}}{R}\right) U\left(x_{n}, z\right)
$$

on the desired set.

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