# ON TWO FUNCTIONALS INVOLVING THE MAXIMUM OF THE TORSION FUNCTION 

Antoine Henrot ${ }^{1}$, Ilaria Lucardesi ${ }^{1, *}$ and Gérard Philippin ${ }^{2}$


#### Abstract

In this paper we investigate upper and lower bounds of two shape functionals involving the maximum of the torsion function. More precisely, we consider $T(\Omega) /(M(\Omega)|\Omega|)$ and $M(\Omega) \lambda_{1}(\Omega)$, where $\Omega$ is a bounded open set of $\mathbb{R}^{d}$ with finite Lebesgue measure $|\Omega|, M(\Omega)$ denotes the maximum of the torsion function, $T(\Omega)$ the torsion, and $\lambda_{1}(\Omega)$ the first Dirichlet eigenvalue. Particular attention is devoted to the subclass of convex sets.


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## 1. Introduction

The two most classical (and most studied) elliptic PDEs are probably the torsion problem, also known as St-Venant problem, and the Dirichlet eigenvalue problem, see (1.3) and (1.5) below. Many estimates and qualitative properties have been obtained for these classical problems, see for example works by Pólya, Szegö, Schiffer, Payne, Hersch, Bandle and many others. In this paper, following these former works, we are interested in finding bounds (if possible optimal) for quantities involving the maximum of the torsion function. We have been particularly inspired by two recent works in $[25,27]$, where the ratio $T(\Omega) \lambda_{1}(\Omega) /|\Omega|$ has been investigated in a similar way. Here $T(\Omega)$ denotes the torsion, $\lambda_{1}(\Omega)$ the first Dirichlet eigenvalue, and $|\Omega|$ the volume of $\Omega$, see Section 1.2 for the precise definitions.

Le $M(\Omega)$ be the maximum of the torsion function. In this paper we investigate upper and lower bounds for the shape functionals

$$
\begin{align*}
& F(\Omega):=\frac{T(\Omega)}{M(\Omega)|\Omega|}  \tag{1.1}\\
& G(\Omega):=M(\Omega) \lambda_{1}(\Omega) \tag{1.2}
\end{align*}
$$

defined over the bounded open sets $\Omega$ of $\mathbb{R}^{d}$ with finite Lebesgue measure. In Section 2 , we prove that the obvious upper bound $F(\Omega) \leq 1$ is actually sharp. Then, we show that for convex domains we have indeed $F(\Omega) \leq 2 / 3$ and we give more precise lower and upper bounds for regular plane convex domains in terms of the

[^0]curvature of their boundaries. In Section 3, we consider the functional $G$. We prove that the easy lower bound $G(\Omega) \geq 1$ is actually sharp. For convex domains, we recall the lower bound $G(\Omega) \geq \pi^{2} / 8$ obtained by Payne. Finding the optimal upper bound for $G$ seems much more difficult. Using topological derivatives, we prove that no maximizer exists in a wide class of domains. When we restrict to the class of convex domains, we can prove existence of an optimal domain but we cannot identify it. In the plane, we suspect that it is the equilateral triangle (which is definitely better than the disk). At last, we write the shape derivative of $G$ and prove that the equilateral triangle does not cancel this shape derivative, in other words it is not a critical point among all regular open sets.

### 1.1. Notations

We adopt standard notations for Lebesgue and Sobolev spaces on a bounded open set of $\mathbb{R}^{d}$, for example $L^{2}(\Omega)$ and $H^{1}(\Omega)$ (space of functions in $L^{2}$ whose derivative, in the sense of distributions, are still in $L^{2}$ ). The boundary values of a Sobolev function are always intended in the sense of traces. The $(d-1)$-dimensional Hausdorff measure is denoted by $\mathcal{H}^{d-1}$.

Given a bounded open set $\Omega \subset \mathbb{R}^{d}$, we denote by $|\Omega|$ its Lebesgue measure, by $f_{\Omega}$ the average integral over it, and by $\mathcal{D}(\Omega)$ the space of $\mathcal{C}^{\infty}$ functions having compact support contained into $\Omega$. The closure of $\mathcal{D}(\Omega)$ in $H^{1}(\Omega)$ is denoted by $H_{0}^{1}(\Omega)$. If the open set $\Omega$ has Lipschitz boundary, we denote by $n$ the outer unit normal vector to $\partial \Omega$, defined a.e. on the boundary.

Given a point $x \in \mathbb{R}^{d}$ and a positive parameter $r>0$, we denote by $B_{r}(x)$ the ball of radius $r$ centered in $x$, and with $\bar{B}_{r}(x)$ its closure.

We define the minimal width of a set as the minimal distance between two parallel supporting hyperplanes.
We denote by $f^{+}$the positive part of a scalar function $f$, namely $f^{+}(x):=\max \{f(x), 0\}$.
The partial derivative of a scalar function $f$ defined in $\mathbb{R}^{d}$ with respect to the $i$-th variable is denoted either by $\partial f / \partial x_{i}$ or by $f_{, i}$; the same notation is used for higher order partial derivatives.

We adopt the convention of summation over repeated indices.

### 1.2. First properties

Given a bounded open set $\Omega$ of $\mathbb{R}^{d}$ with finite Lebesgue measure, we denote by $u_{\Omega}$ the torsion function of $\Omega$, that is, the solution of

$$
\left\{\begin{array}{l}
-\Delta u=1 \quad \text { in } \Omega  \tag{1.3}\\
u \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

and we set

$$
\begin{equation*}
T(\Omega):=\left\|u_{\Omega}\right\|_{L^{1}(\Omega)}, \quad M(\Omega):=\left\|u_{\Omega}\right\|_{L^{\infty}(\Omega)} . \tag{1.4}
\end{equation*}
$$

It is easy to check that $u_{\Omega}$ is $C^{\infty}$ inside $\Omega$ and non negative in $\bar{\Omega}$, thus

$$
T(\Omega)=\int_{\Omega} u_{\Omega} \mathrm{d} x, \quad M(\Omega)=\max _{\Omega} u_{\Omega}
$$

We denote by $\lambda_{1}(\Omega)$ first eigenvalue of the Dirichlet Laplacian and by $\varphi_{\Omega}$ the corresponding (normalized) eigenfunction, that is, the solution of

$$
\left\{\begin{array}{l}
-\Delta \varphi=\lambda_{1}(\Omega) \varphi \quad \text { in } \Omega  \tag{1.5}\\
\varphi \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

with $\left\|\varphi_{\Omega}\right\|_{L^{2}(\Omega)}=1$.
We recall that the functionals $T$ and $\lambda_{1}$ admit the following variational formulations:

$$
\begin{equation*}
T(\Omega)=\sup _{v \in H_{0}^{1}(\Omega) \backslash\{0\}} \frac{\left(\int_{\Omega} v \mathrm{~d} x\right)^{2}}{\int_{\Omega}|\nabla v|^{2} \mathrm{~d} x}, \quad \lambda_{1}(\Omega)=\inf _{v \in H_{0}^{1}(\Omega) \backslash\{0\}} \frac{\int_{\Omega}|\nabla v|^{2} \mathrm{~d} x}{\int_{\Omega} v^{2} \mathrm{~d} x} . \tag{1.6}
\end{equation*}
$$

It follows from the homogeneity relations

$$
T(t \Omega)=t^{d+2} T(\Omega), \quad M(t \Omega)=t^{2} M(\Omega), \quad \lambda_{1}(t \Omega)=t^{-2} \lambda_{1}(\Omega), \quad t>0
$$

that both $F$ and $G$ are scale invariant.
In the sequel, when no ambiguity may arise, we will denote the torsion function and the first eigenfunction of the Dirichlet Laplacian of a given set $\Omega$ simply by $u$ and $\varphi$, respectively.

Remark 1.1. Note that the functional $G$ can be studied on a larger class of sets, satisfying $\lambda_{1}(\Omega)>0$ and $\left\|u_{\Omega}\right\|_{\infty}<+\infty$. In particular, since these two conditions are equivalent (see, e.g., [26]), $G$ is naturally defined on sets on which a Poincaré inequality holds (e.g., sets bounded in one direction).

## 2. Bounds for the functional $F$

### 2.1. The upper bound

The upper bound $F(\Omega) \leq 1$ is obvious. Actually, we are going to prove that this bound is sharp. This is not so intuitive since the equality $F(\Omega)=1$ is only true for constant functions and clearly a torsion function of any domain $\Omega$ is a priori far to be constant. The idea is to use the theory of homogenization. Indeed, by performing suitable spherical holes (with the appropriate radius) in a domain $\Omega$, we are able to get a sequence of torsion functions which $\gamma$-converges to something which is no longer a torsion function : the "strange term coming from nowhere" in the celebrated paper by D. Cioranescu and F. Murat, [3]. Our theorem is the following.

Theorem 2.1. In any dimension, we can find a sequence of domains $\Omega_{\epsilon}$ such that $F\left(\Omega_{\epsilon}\right) \rightarrow 1$.
Here we recall the construction of a sequence of perforated domains introduced by Cioranescu-Murat in [3], see also [9] for a more precise estimate and convergence result.

Let $\Omega \subset \mathbb{R}^{d}, d \geq 2$ a regular (or a convex) domain, and $C_{0}>0$ be fixed. For every $\epsilon>0$, consider the ball $T_{\epsilon}:=B_{r_{\epsilon}}(0)$ with a radius $r_{\epsilon}$ which satisfies

$$
r_{\epsilon}=\left\{\begin{array}{ccc}
C_{0} \epsilon^{d /(d-2)} & \text { if } \quad d \geq 3  \tag{2.1}\\
\exp \left(-C_{0} / \epsilon^{2}\right) & \text { if } & d=2
\end{array}\right.
$$

and the perforated domain

$$
\begin{equation*}
\Omega_{\epsilon}:=\Omega \backslash \cup_{z \in \mathbb{Z}^{d}}\left(2 \epsilon z+\bar{T}_{\epsilon}\right) \tag{2.2}
\end{equation*}
$$

Note that the removed holes form a periodic set in the plane, with period $2 \epsilon$. Now let $u_{\epsilon}$ denote the torsion function of the perforated domain $\Omega_{\epsilon}$, extended to zero in the holes. It is proved in [3] that the sequence $u_{\epsilon}$ converges weakly in $H_{0}^{1}(\Omega)$ (and strongly in $\left.L^{2}(\Omega)\right)$ to the solution $u^{*}$ of

$$
\left\{\begin{array}{l}
-\Delta u^{*}+a u^{*}=1 \quad \text { in } \quad \Omega \\
u^{*} \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

where the constant $a$ satisfies

$$
a=\left\{\begin{array}{ccc}
\frac{C_{0}^{d-2}}{2^{d}} d(d-2) \omega_{d} & \text { if } & d \geq 3  \tag{2.3}\\
\frac{\pi}{2 C_{0}} & \text { if } & d=2
\end{array}\right.
$$

and $\omega_{d}$ is the volume of the unit ball in $\mathbb{R}^{d}$. As a consequence we have

$$
\begin{equation*}
\int_{\Omega_{\epsilon}} u_{\epsilon} \mathrm{d} x \rightarrow \int_{\Omega} u^{*} \mathrm{~d} x, \quad\left|\Omega_{\epsilon}\right| \rightarrow|\Omega|, \quad \text { as } \epsilon \rightarrow 0 \tag{2.4}
\end{equation*}
$$

Now we want to analyze the asymptotic behavior of the $L^{\infty}$ norm of the functions $u_{\epsilon}$. We cannot hope for uniform convergence of $u_{\epsilon}$ to $u^{*}$, nevertheless we can prove the convergence of the $L^{\infty}$ norms:

Theorem 2.2. Let $u_{\epsilon}$ be the torsion functions of the perforated domains $\Omega_{\epsilon}$ extended to zero in the holes and let $u^{*}$ be their weak limit in $H_{0}^{1}(\Omega)$.

Then, up to a subsequence, $M\left(\Omega_{\epsilon}\right)=\left\|u_{\epsilon}\right\|_{L^{\infty}(\Omega)} \rightarrow\left\|u^{*}\right\|_{L^{\infty}(\Omega)}$ as $\epsilon \rightarrow 0$.
Proof. We are indebted to G. Buttazzo and B. Velichkov of this proof (see the final comment at the end of the paper).

First of all, up to a subsequence, we can assume that $u_{\epsilon}$ converges pointwise almost everywhere to $u^{*}$ :

$$
\text { for a.e. } x \in \Omega, \quad u_{\epsilon}(x) \rightarrow u^{*}(x)
$$

Applying this to a ball centered at a point where $u^{*}$ is maximum, we infer that

$$
\begin{equation*}
\left\|u^{*}\right\|_{L^{\infty}(\Omega)} \leq \liminf _{\epsilon}\left\|u_{\epsilon}\right\|_{L^{\infty}(\Omega)} \tag{2.5}
\end{equation*}
$$

Now let us assume that the inequality in (2.5) is strict. Then we could find two positive numbers $b_{1}<b_{2}$ such that

$$
\left\|u^{*}\right\|_{L^{\infty}(\Omega)} \leq b_{1}<b_{2} \leq \liminf _{\epsilon}\left\|u_{\epsilon}\right\|_{L^{\infty}(\Omega)}
$$

It is proved in [28], (Prop. 3.2.34) that for any non negative function $v$ satisfying $\Delta v+1 \geq 0$ in $\Omega$, the following inequality holds:

$$
\begin{equation*}
\|v\|_{L^{\infty}(\Omega)} \leq C\left(\int_{\Omega} v(x) \mathrm{d} x\right)^{d /(d+2)} \tag{2.6}
\end{equation*}
$$

where $C$ is a positive constant which only depends on $\Omega$. By taking in (2.6) the subsolution $v:=\left(u_{\epsilon}-b\right)^{+}$to the torsion equation in $\Omega$, with $b:=\left(b_{1}+b_{2}\right) / 2$, we obtain

$$
\begin{equation*}
\left\|\left(u_{\epsilon}-b\right)^{+}\right\|_{L^{\infty}(\Omega)} \leq C\left(\int_{\Omega}\left(u_{\epsilon}-b\right)^{+} \mathrm{d} x\right)^{d /(d+2)} \tag{2.7}
\end{equation*}
$$

It comes on the one hand

$$
\underset{\epsilon}{\liminf }\left\|\left(u_{\epsilon}-b\right)^{+}\right\|_{L^{\infty}(\Omega)} \geq b_{2}-b=\left(b_{2}-b_{1}\right) / 2>0
$$

while on the other hand, by $L^{2}$ convergence,

$$
\int_{\Omega}\left(u_{\epsilon}-b\right)^{+} \mathrm{d} x \rightarrow \int_{\Omega}\left(u^{*}-b\right)^{+} \mathrm{d} x=0
$$

contradicting inequality (2.7).
Now we are in position to prove Theorem 2.1. Let us introduce $v^{*}:=a u^{*}$. In the following Lemma, we list some properties of this function.

Lemma 2.3. Let $a \in \mathbb{R}^{+}$and $v^{*}$ be the solution of

$$
\begin{equation*}
-a^{-1} \Delta v^{*}+v^{*}=1 \tag{2.8}
\end{equation*}
$$

in $H_{0}^{1}(\Omega)$. Then $0<v^{*} \leq 1$ in $\Omega$ and, in the limit as $a \rightarrow+\infty$, $v^{*} \rightharpoonup 1$ weakly in $L^{2}(\Omega)$.

Proof. The positivity of $v^{*}$ in $\Omega$ is a simple consequence of the maximum principle. For $x_{0} \in \Omega$ maximum point for $v^{*}$ it holds $\Delta v^{*}\left(x_{0}\right) \leq 0$. In particular, for every $x \in \Omega$ we have

$$
v^{*}(x) \leq v^{*}\left(x_{0}\right) \leq-\Delta v^{*}\left(x_{0}\right)+v^{*}\left(x_{0}\right)=1
$$

which proves the upper bound. Exploiting the optimality of $v^{*}$ for the functional

$$
\begin{equation*}
H_{0}^{1}(\Omega) \ni v \mapsto \frac{1}{2} \int_{\Omega} a^{-1}|\nabla v|^{2} \mathrm{~d} x+\int_{\Omega}\left(v^{2}-v\right) \mathrm{d} x \tag{2.9}
\end{equation*}
$$

it is easy to see that $v^{*}$ and $a^{-1 / 2} \nabla v^{*}$ are uniformly (with respect to $a$ ) bounded in $L^{2}(\Omega)$ and $L^{2}\left(\Omega ; \mathbb{R}^{2}\right)$, respectively. The former bound implies that, in the limit as $a \rightarrow+\infty$, up to subsequences, $v^{*}$ weakly converges in $L^{2}(\Omega)$ to some $\bar{v}^{*}$. The latter bound, combined with (2.8), implies that the weak limit $\bar{v}^{*}$ is 1 .

An immediate consequence of the previous lemma is

$$
\begin{equation*}
\liminf _{\epsilon}\left\|u_{\epsilon}\right\|_{L^{\infty}(\Omega)}=\left\|u^{*}\right\|_{L^{\infty}(\Omega)} \leq \frac{1}{a} \tag{2.10}
\end{equation*}
$$

Therefore, using (2.4) and (2.10), we get, for a subsequence:

$$
\limsup _{\epsilon} F\left(\Omega_{\epsilon}\right) \geq \frac{\int_{\Omega} u^{*} \mathrm{~d} x}{\left(\max u^{*}\right)|\Omega|}=\frac{\int_{\Omega} v^{*} \mathrm{~d} x}{\left(\max v^{*}\right)|\Omega|} \geq \frac{\int_{\Omega} v^{*} \mathrm{~d} x}{|\Omega|} \rightarrow 1
$$

as $a \rightarrow+\infty$. This finishes the proof of Theorem 2.1.

### 2.2. The upper bound for convex sets

The maximizing sequence used in the previous section is very specific, thus we can expect that in the convex case we can significantly improve the upper bound. Indeed, let us prove the following

Theorem 2.4. Let $\Omega$ be any bounded convex domain in $\mathbb{R}^{d}$, then

$$
\begin{equation*}
F(\Omega) \leq \frac{2}{3} \tag{2.11}
\end{equation*}
$$

Moreover, inequality (2.11) is sharp.
Proof. We use the maximum principle for $P$-functions. Following [15] (for regular convex domains) or [21] (for general convex domains), it is known that the function $\psi:=|\nabla u|^{2}+2 u$ takes its maximum at a critical point of $u$, namely at the point where $u$ is maximum. Note that $\sqrt{u}$ has only one critical point since it is strictly concave (see, e.g., [10] for the concavity property and [11], (Introduction) for the strict one). Therefore, for every $x \in \Omega$, we have

$$
\begin{equation*}
|\nabla u(x)|^{2}+2 u(x) \leq 2 M(\Omega) \tag{2.12}
\end{equation*}
$$

in particular, integrating (2.12) over $\Omega$ yields $3 T(\Omega) \leq 2 M(\Omega)|\Omega|$.
In order to prove the sharpness of the inequality, let us consider the sequence of rectangles in the plane $\Omega_{n}:=(-n, n) \times(0,1)$. The same construction holds in any dimension $d$, using the sequence of parallelepipeds $\Omega_{n}:=(-n, n)^{d-1} \times(0,1)$. Let us denote by $u_{n}$ the torsion function of $\Omega_{n}$. By the maximum principle, we have

$$
\begin{equation*}
u_{n}(x, y) \leq \frac{1}{2} y(1-y) \tag{2.13}
\end{equation*}
$$

The function $\frac{1}{2} y(1-y)$ can be seen as the torsion function of the unbounded strip $\{0<y<1\}$. Therefore (2.13) implies that $M\left(\Omega_{n}\right) \leq \frac{1}{8}$.

Now, in view of (1.3), it is easy to check that the torsion admits the variational formulation

$$
\begin{equation*}
-\frac{1}{2} T\left(\Omega_{n}\right)=\min _{v \in H_{0}^{1}\left(\Omega_{n}\right)}\left\{\frac{1}{2} \int_{\Omega_{n}}|\nabla v|^{2} \mathrm{~d} x-\int_{\Omega_{n}} v \mathrm{~d} x\right\} \tag{2.14}
\end{equation*}
$$

Let us introduce the function $\psi_{n}(x)$ defined as

$$
\psi_{n}(x):=\left\{\begin{array}{ccl}
1 & \text { if } & x \in[-n+1, n-1]  \tag{2.15}\\
n-x & \text { if } & x \in[n-1, n] \\
x-n & \text { if } & x \in[-n,-n+1] \\
0 & \text { if } & |x|>n
\end{array}\right.
$$

and let us choose as a test function in (2.14) the function $v(x, y):=\psi_{n}(x) \frac{1}{2} y(1-y)$, which is an element of $H_{0}^{1}\left(\Omega_{n}\right)$. We immediately get

$$
\int_{\Omega_{n}} v \mathrm{~d} x=\int_{-n}^{n} \psi_{n}(x) \mathrm{d} x \int_{0}^{1} \frac{1}{2} y(1-y) \mathrm{d} y=\frac{n}{6}-\frac{1}{12} \geq \frac{n}{6}-1
$$

Since $\left|\psi_{n}^{\prime}(x)\right|=1$ if $x \in(-n,-n+1) \cup(n-1, n)$ and it is 0 otherwise, and $y(1-y) / 2<1$ for every $y \in(0,1)$, we have

$$
\int_{\Omega_{n}}|\nabla v|^{2} \mathrm{~d} x \leq 2+\int_{-n}^{n} \psi_{n}^{2}(x) \mathrm{d} x \int_{0}^{1}\left(\frac{1}{2}-y\right)^{2} \mathrm{~d} y=2+\frac{n}{6}-\frac{1}{9} \leq \frac{n}{6}+2
$$

Thus

$$
-\frac{1}{2} T\left(\Omega_{n}\right) \leq-\frac{n}{12}+2
$$

which implies that

$$
F\left(\Omega_{n}\right) \geq \frac{2}{3}-\frac{16}{n} \rightarrow 2 / 3 \quad \text { when } n \rightarrow+\infty
$$

For strictly convex and regular domains in the plane, one can improve this upper bound by the following:
Theorem 2.5. Let $\Omega$ be a strictly convex bounded domain of class $C^{2}$ in $\mathbb{R}^{2}$. Let us introduce the quantity $\beta$ which depends only on the geometry of $\Omega$ (actually its curvature $k$ ):

$$
\begin{equation*}
\beta=2-\frac{1}{4}\left(\frac{\min _{\partial \Omega} k}{\max _{\partial \Omega} k}\right)^{3} \leq 2 \tag{2.16}
\end{equation*}
$$

Then we have

$$
\begin{equation*}
F(\Omega) \leq \frac{\beta}{\beta+1} \leq \frac{2}{3} \tag{2.17}
\end{equation*}
$$

We postpone the proof of this Theorem to Section 2.4, where the proof for a similar lower bound will also be given at the same time.

### 2.3. The lower bound

Clearly, by the positivity of $T, M$, and Lebesgue measure, the infimum of $F$ is greater than or equal to zero. It is easy to show that the lower bound 0 is optimal: consider the sequence of sets

$$
\Omega_{n}:=B_{1}(0) \bigcup_{i=1}^{n} B_{r_{n}}\left(x_{i}\right), \quad n \in \mathbb{N}
$$

with $x_{1}, \ldots, x_{n} \neq 0$ distinct points in a compact set and $r_{n}>0$ a small parameter (whose precise value will be chosen later). In this case, $u_{n}:=u_{\Omega_{n}}$ is the sum of the torsion functions associated to every single connected component of $\Omega_{n}$, namely $u_{n}=\sum_{i=0}^{n} u_{i}$ with $u_{0}:=u_{B_{1}(0)}$ and $u_{i}:=u_{B_{n^{-1 / 4}\left(x_{i}\right)}}, i=1 \ldots, n$. Since

$$
u_{0}(x)=\frac{1-|x|^{2}}{2 d} \quad \text { and } \quad u_{i}=\frac{r_{n}^{2}-\left|x-x_{i}\right|^{2}}{2 d}
$$

it is easy to see that

$$
T\left(\Omega_{n}\right)=\frac{\omega_{d}}{d^{2}(d+2)}\left(1+n r_{n}^{d+2}\right), \quad M\left(\Omega_{n}\right)=\frac{1}{2 d}, \quad\left|\Omega_{n}\right|=\frac{\omega_{d}}{d}\left(1+n r_{n}^{d}\right)
$$

where $\omega_{d}$ is the volume of the unit ball in $\mathbb{R}^{d}$. By taking $r_{n}=n^{-1 /(2 d)}$, we infer that

$$
F\left(\Omega_{n}\right)=\frac{2}{d+2} \frac{n^{\frac{1}{2}-\frac{1}{d}}+1}{n^{\frac{1}{2}}+1} \sim n^{-\frac{1}{d}} \rightarrow 0
$$

implying that $\inf F=0$.

### 2.4. The lower bound for convex sets

By strict concavity of $\sqrt{u}$ when $\Omega$ is convex, see e.g. [11], it is easy to get a lower bound for convex sets:
Theorem 2.6. Let $\Omega$ be any bounded convex set in $\mathbb{R}^{d}$, then

$$
\begin{equation*}
F(\Omega) \geq \frac{1}{(d+1)^{2}} \tag{2.18}
\end{equation*}
$$

Proof. Since $\sqrt{u}$ is concave, its graph is above the cone of basis $\Omega$ and vertex $M_{0}$ the maximum point of $\sqrt{u}$. Therefore, by comparison of the volumes:

$$
\int_{\Omega} \sqrt{u(x)} \mathrm{d} x \geq \frac{\sqrt{M(\Omega)}|\Omega|}{d+1}
$$

By taking the square of the previous inequality and using Cauchy-Schwarz inequality for the left-hand side

$$
\left(\int_{\Omega} \sqrt{u(x)} \mathrm{d} x\right)^{2} \leq|\Omega| \int_{\Omega} u(x) \mathrm{d} x
$$

we get the desired inequality.
We believe that inequality (2.18) is not optimal. For example, in the plane, we conjecture:
Conjecture: For any plane convex domain, the following lower bound holds:

$$
\begin{equation*}
F(\Omega) \geq \frac{1}{3} \tag{2.19}
\end{equation*}
$$

Moreover, this inequality should be optimal, and a minimizing sequence could be a sequence of isosceles triangles degenerating to a segment. Let us remark that when $u$ is concave, for instance when $\Omega$ is an ellipse, we obtain exactly in the same way

$$
\begin{equation*}
\frac{T(\Omega)}{|\Omega| M(\Omega)} \geq \frac{1}{3} \tag{2.20}
\end{equation*}
$$

Sufficient conditions on the geometry of $\Omega$ to insure the concavity of $u$ have been established by Kosmodem'yanskii in [12].

Let us conclude this section with a theorem in the spirit of Theorem 2.5.

Theorem 2.7. Let $\Omega$ be a strictly convex bounded domain of class $C^{2}$ in $\mathbb{R}^{2}$. Then we have

$$
\begin{equation*}
F(\Omega) \geq \frac{1}{4}\left(\frac{\min _{\partial \Omega} k}{\max _{\partial \Omega} k}\right)^{3} \tag{2.21}
\end{equation*}
$$

where $k$ is the curvature of $\Omega$.
Set for brevity $\alpha:=\left(\min _{\partial \Omega} k / \max _{\partial \Omega} k\right)^{3} / 4$. Note that inequality (2.21) is better than the general inequality eqrefineqconv when $\alpha>1 / 8$, which occurs when $\min _{\partial \Omega} k>\max _{\partial \Omega} k / \sqrt[3]{2}$. In [19], Payne and Philippin have derived sharp upper bounds for $|\nabla u|$. The goal of this section is to derive new lower bounds for these quantities by using the same approach as in [19]. For the torsion problem one can associate an auxiliary function involving the curvature $k$ of the level lines $\{u=$ const. $\}$. Properly chosen, the auxiliary function turns out to satisfy some minimum principles, implying the convexity of the level sets of $u$, under suitable convexity assumptions on $\Omega$. These results have been derived by Makar-Limanov in [13] for the torsion problem with the associate function

$$
\begin{equation*}
P(x):=k|\Delta u|^{3}+u\left[(\Delta u)^{2}-u_{, i j} u_{, i j}\right] . \tag{2.22}
\end{equation*}
$$

Proof of Theorems 2.5 and 2.7. Making use of normal coordinates with respect to the level lines $\{u=$ const. $\}$ we have

$$
\begin{gather*}
|\nabla u|^{2}=u_{, i} u_{, i}=u_{n}^{2}  \tag{2.23}\\
\Delta u=u_{n n}+k u_{n}  \tag{2.24}\\
u_{, i j} u_{, i j}=u_{n n}^{2}+k^{2} u_{n}^{2}+2 u_{n s}^{2} \tag{2.25}
\end{gather*}
$$

where an index $n$ stands for the outward normal derivative and an index $s$ stands for the derivative along the level lines $\{u=$ const. $\}$, and $k$ is the curvature of the level lines defined as

$$
\begin{equation*}
k:=-\left(\frac{u_{, i}}{|\nabla u|}\right)_{, i}=\frac{u_{, i j} u_{, i} u_{, j}-|\nabla u|^{2} \Delta u}{|\nabla u|^{3}} . \tag{2.26}
\end{equation*}
$$

The Makar-Limanov function $P$ introduced in (2.22) may be rewritten in terms of normal coordinates as

$$
\begin{equation*}
P=|\nabla u|^{3} k-2\left[k u_{n}+k^{2} u_{n}^{2}+u_{n s}^{2}\right] . \tag{2.27}
\end{equation*}
$$

Makar-Limanov's result is based on the fact that $P$ is super-harmonic. It then follows that $P$ takes its minimum value $P_{\min }$ on $\partial \Omega$, so that the following quadratic inequality for $k$ holds:

$$
\begin{equation*}
|\nabla u|^{3} k-2 u\left[k u_{n}+k^{2} u_{n}^{2}+u_{n s}^{2}\right] \geq P_{\min }, \quad x \in \Omega . \tag{2.28}
\end{equation*}
$$

Omitting the term containing $u_{n s}^{2}$ and solving (2.28) for $k$, we obtain

$$
\begin{equation*}
\left.\frac{\Phi}{4 u}\{1-\sqrt{1-z}\} \leq k|\nabla u| \leq \frac{\Phi}{4 u}\{1+\sqrt{1-z}\}\right\} \tag{2.29}
\end{equation*}
$$

with

$$
\begin{align*}
\Phi & :=|\nabla u|^{2}+2 u  \tag{2.30}\\
z & :=\frac{8 P_{\min } u}{\Phi^{2}} \tag{2.31}
\end{align*}
$$

We note that $z \leq 1$ in view of the inequality

$$
\begin{equation*}
P \leq \frac{1}{8 u} \Phi^{2} \quad \text { in } \Omega, \tag{2.32}
\end{equation*}
$$

derived in [20]. Multiplying (2.29) by $-2|\nabla u| \sqrt{u}$, we obtain

$$
\begin{equation*}
-\frac{\Phi|\nabla u|}{2 \sqrt{u}}\{1+\sqrt{1-z}\} \leq \sqrt{u} \frac{\partial \Phi}{\partial n} \leq-\frac{\Phi|\nabla u|}{2 \sqrt{u}}\{1-\sqrt{1-z}\} \tag{2.33}
\end{equation*}
$$

in view of

$$
\begin{equation*}
\frac{\partial \Phi}{\partial n}=-2 k|\nabla u|^{2} \tag{2.34}
\end{equation*}
$$

For convenience we set

$$
\begin{equation*}
\theta:=\frac{\Phi}{\sqrt{u}}=\frac{|\nabla u|^{2}}{\sqrt{u}}+2 \sqrt{u} \tag{2.35}
\end{equation*}
$$

Replacing (2.35) in (2.33), these inequalities reduce to

$$
\begin{equation*}
\frac{\partial u}{\partial n} \sqrt{\left(\frac{\theta}{2}\right)^{2}-2 P_{\min }} \leq u \frac{\partial \theta}{\partial n} \leq-\frac{\partial u}{\partial n} \sqrt{\left(\frac{\theta}{2}\right)^{2}-2 P_{\min }} \tag{2.36}
\end{equation*}
$$

which are equivalent to

$$
\begin{equation*}
-\frac{1}{2} \frac{\mathrm{~d} u}{u} \leq-\frac{\mathrm{d} \theta}{\sqrt{\theta^{2}-8 P_{\min }}} \leq \frac{1}{2} \frac{\mathrm{~d} u}{u} \tag{2.37}
\end{equation*}
$$

These inequalities link the functions $u$ and $\theta$ to their differentials along the orthogonal trajectories of the level lines (also called fall lines of $u$ ). Rewriting (2.37) in the form

$$
\begin{equation*}
-\frac{1}{2} \mathrm{~d}(\log u) \leq-\mathrm{d}\left(\log \left[\theta+\sqrt{\theta^{2}-8 P_{\min }}\right]\right) \leq \frac{1}{2} \mathrm{~d}(\log u) \tag{2.38}
\end{equation*}
$$

and integrating from a point $x \in \Omega$ to the maximum point $x_{0}$ of $u$ along the fall line joining these points, we obtain

$$
\begin{equation*}
\sqrt{\frac{u(x)}{u\left(x_{0}\right)}} \leq \frac{\theta(x)+\sqrt{\theta^{2}(x)-8 P_{\min }}}{\theta_{0}+\sqrt{\theta_{0}^{2}-8 P_{\min }}} \leq \sqrt{\frac{u\left(x_{0}\right)}{u(x)}} \tag{2.39}
\end{equation*}
$$

with

$$
\begin{equation*}
\theta_{0}:=\theta\left(x_{0}\right)=2 \sqrt{u\left(x_{0}\right)} \tag{2.40}
\end{equation*}
$$

Multiplying (2.39) by $\left(\theta_{0}+\sqrt{\theta_{0}^{2}-8 P_{\min }}\right) \sqrt{u}$, replacing back $\theta$ by $\Phi$, and recalling that $u\left(x_{0}\right)=M(\Omega)$, we obtain

$$
\begin{align*}
2 u\left(1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}\right)-\Phi & \leq \sqrt{\Phi^{2}-8 P_{\min } u}  \tag{2.41}\\
& \leq 2 M(\Omega)\left(1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}\right)-\Phi
\end{align*}
$$

Squaring (2.41) and solving for $\Phi$, we obtain

$$
\begin{align*}
\frac{2 P_{\min }}{1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}} & +u\left(1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}\right) \leq \Phi  \tag{2.42}\\
& \leq \frac{2 P_{\min } u}{M(\Omega)\left(1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}\right.}+M(\Omega)\left(1+\sqrt{1-\frac{2 P_{\min }}{M(\Omega)}}\right) \tag{2.43}
\end{align*}
$$

Replacing (2.30) in (2.42), after some reduction we obtain the basic inequalities

$$
\begin{equation*}
\tilde{\alpha}(M(\Omega)-u) \leq|\nabla u|^{2} \leq \tilde{\beta}(M(\Omega)-u) \quad \text { in } \Omega \tag{2.44}
\end{equation*}
$$

with

$$
\begin{align*}
& \tilde{\alpha}:=1-\sqrt{1-\frac{2 P_{\mathrm{min}}}{M(\Omega)}}  \tag{2.45}\\
& \tilde{\beta}:=1+\sqrt{1-\frac{2 P_{\mathrm{min}}}{M(\Omega)}} \tag{2.46}
\end{align*}
$$

The upper bound for $|\nabla u|^{2}$ in (2.44) was already derived in [19]. The lower bound is nontrivial only for strictly convex $\Omega$, whereas the upper bound makes sense even for nonconvex $\Omega$. However in this case $P_{\min }$ is negative, and $\tilde{\beta}$ is greater than two. We note that inequalities (2.44) are exact when $\tilde{\alpha}=\tilde{\beta}$, i.e. when $2 P_{\min }(M(\Omega))^{-1}=1$. This is the case if and only if $\Omega$ is a disk. For practical use of (2.44) a computable positive lower bound for the quantity $2 P_{\min }(M(\Omega))^{-1}$ is needed. To this end, we write

$$
\begin{equation*}
P_{\min }=\min _{\partial \Omega}\left(|\nabla u|^{3} k\right) \geq\left(\min _{\partial \Omega}|\nabla u|\right)^{3}\left(\min _{\partial \Omega} k\right) \tag{2.47}
\end{equation*}
$$

and make use of the inequalities

$$
\begin{gather*}
\min _{\partial \Omega}|\nabla u| \geq \frac{1}{2 \max _{\partial \Omega} k}  \tag{2.48}\\
M(\Omega)=\max _{\Omega} u \leq \frac{1}{2} \rho^{2} \leq \frac{1}{2}\left(\min _{\partial \Omega} k\right)^{-2} \tag{2.49}
\end{gather*}
$$

derived in $[17,18]$, where $\rho$ is the inradius of $\Omega$. Using (2.47), (2.48), and (2.49), for strictly convex $\Omega$ we have

$$
\begin{equation*}
\frac{2 P_{\min }}{M(\Omega)} \geq \frac{2\left(\min _{\partial \Omega}|\nabla u|\right)^{3}\left(\min _{\partial \Omega} k\right)}{M(\Omega)} \geq \frac{1}{2}\left(\frac{\min _{\partial \Omega} k}{\max _{\partial \Omega} k}\right)^{3}=2 \alpha \tag{2.50}
\end{equation*}
$$

Replacing (2.50) in (2.45) and in (2.46), we obtain the bounds

$$
\begin{equation*}
\tilde{\alpha} \geq \alpha, \quad \tilde{\beta} \leq \beta:=2-\alpha \tag{2.51}
\end{equation*}
$$

in particular, the quantities $\alpha$ and $\beta$ may be used in (2.44) instead of $\tilde{\alpha}$ and $\tilde{\beta}$, respectively. Integrating (2.44) over $\Omega$ and exploiting the estimates (2.51), we obtain the following bounds for $\frac{T(\Omega)}{|\Omega| M(\Omega)}$ :

$$
\begin{equation*}
\frac{\alpha}{\alpha+1} \leq \frac{T(\Omega)}{|\Omega| M(\Omega)} \leq \frac{\beta}{\beta+1} \leq \frac{2}{3} \tag{2.52}
\end{equation*}
$$

The upper bound proves Theorem 2.5 .
A better lower bound for $\frac{T(\Omega)}{|\Omega| M(\Omega)}$ may be derived by integrating the inequality

$$
\begin{equation*}
P(x)=u_{, i j} u_{. i} u_{. j}-|\nabla u|^{2} \Delta u+u\left[(\Delta u)^{2}-u_{. i j} u_{, i j}\right] \geq P_{\min } \tag{2.53}
\end{equation*}
$$

over $\Omega$. Making use of

$$
\begin{equation*}
\int_{\Omega} u_{, i j} u_{, i} u_{, j} \mathrm{~d} x=-\int_{\Omega} u u_{, i j} u_{, i j} \mathrm{~d} x \tag{2.54}
\end{equation*}
$$

we obtain

$$
\begin{equation*}
\int_{\Omega} P(x) \mathrm{d} x=\int_{\Omega} u\left[(\Delta u)^{2}-2 u_{, i j} u_{, i j}\right] \mathrm{d} x+T(\Omega) \geq P_{\min }|\Omega| \tag{2.55}
\end{equation*}
$$

Since $(\Delta u)^{2}-2 u, i j u_{, i j} \leq 0$, it follows that

$$
\begin{equation*}
T(\Omega) \geq P_{\min }|\Omega| \tag{2.56}
\end{equation*}
$$

This inequality, together with (2.50), gives the lower bound $F(\Omega) \geq \alpha$, concluding the proof of Theorem 2.7.

## 3. Bounds for the functional $G$

### 3.1. The upper bound

The best available upper bound for $G$, still not optimal, is very recent and is due to Vogt: in ([29], Thm 1.5), exploiting semigroups techniques, the author proved that

$$
G(\Omega) \leq \frac{d}{8}+\frac{1}{4} \sqrt{5(1+\ln 2 / 4)} \sqrt{d}+1 .
$$

Among the upper bounds obtained before Vogt, let us mention the one obtained in Theorem 1 of [26], stating that

$$
G(\Omega) \leq 3 d \ln 2+4
$$

and the one in [24], based on a variational technique, simple and interesting at te same time.
Finding the optimal upper bound suggests to look at the shape optimization problem:

$$
\begin{equation*}
\mathcal{P}_{G} \quad \sup \left\{G(\Omega), \Omega \subset \mathbb{R}^{d}\right\} . \tag{3.1}
\end{equation*}
$$

Even if it looks like as a standard shape optimization problem, the existence of a solution is not clear for us. We believe that a maximizer does not exist and a partial result in this direction is given by the following

Proposition 3.1. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open set of class $C^{2}$. Assume that $G$ is differentiable at $\Omega$ (i.e., the shape derivative of $G$ at $\Omega$ exists and is given by (3.13)). Then $\Omega$ is not a maximizer for $G$.

The proof is postponed to Section 3.5 and is based on a topological derivative argument: under suitable regularity assumptions on $\Omega$, removing a small hole near the boundary makes $G$ increase.

In order to further investigate $\mathcal{P}_{G}$, other useful tools are represented by numerical tests and the theory of shape derivatives. The former technique suggests that, in the case of polygons in the plane, the optimum should be non convex. The latter, that we detail in Section 3.5, provides a necessary condition for critical shapes; in particular, it turns out that the equilateral triangle, even if strictly better than the disk, is not optimal for $\sup G$ in dimension $d=2$ (see Cor. 3.7 below).

Let us now consider the restricted class of convex domains, for which the equilateral triangle could be a maximizer.

### 3.2. The maximization problem in the convex framework

Unlike what happens in the general case, if we add the convex constraint, the existence of a maximizer for $G$ is guaranteed.

Theorem 3.2. The shape functional $G$ admits a maximizer in the class of bounded convex sets of $\mathbb{R}^{d}$.
Proof. Let $\Omega_{n}$ be a maximizing sequence of convex subsets of $\mathbb{R}^{d}$. By the scale invariance of $G$, without loss of generality, we may assume that the elements of the sequence are all contained in a fixed bounded set $K$. By the Blaschke selection theorem, there exists a subsequence (not relabeled) converging to some convex set $\Omega \subset K$ in the Hausdorff metric.

We claim that the minimal width $w_{n}$ of $\Omega_{n}$ does not vanish as $n \rightarrow+\infty$, so that the limit set $\Omega$ has nonempty interior: choosing a suitable reference frame in $\mathbb{R}^{d}$, we may assume that $\Omega_{n}$ is contained in the strip
$\left\{x \in \mathbb{R}^{d}: 0 \leq x_{d} \leq w_{n}\right\}$; by the maximum principle, it is easy to see that the torsion function $u_{n}$ of $\Omega_{n}$ satisfies $u_{n}(x) \leq x_{d}\left(w_{n}-x_{d}\right) / 2$ in $\bar{\Omega}_{n}$, in particular

$$
\begin{equation*}
M\left(\Omega_{n}\right) \leq \frac{w_{n}^{2}}{8} ; \tag{3.2}
\end{equation*}
$$

on the other hand, in equation (4.8) from [27], the authors provide the following upper bound for $\lambda_{1}$ in terms of the minimal width:

$$
\begin{equation*}
\lambda_{1}\left(\Omega_{n}\right) \leq \frac{\pi^{2}}{w_{n}^{2}}\left(1+c_{n}\left(3 / 2+3 / 2^{1 / 3}+2^{1 / 3}\right)\right) \tag{3.3}
\end{equation*}
$$

with $c_{n}=c_{n}\left(\Omega_{n}\right)$ a positive constant vanishing as $w_{n} \rightarrow 0$ (see (19) in [27]). By combining (3.2), (3.3), and the lower bound (3.5), we conclude that $G\left(\Omega_{n}\right) \rightarrow \pi^{2} / 8$. This gives a contradiction, since $\pi^{2} / 8$ is clearly not the maximum of $G$.

In order to conclude the proof, we need to show that the limit set $\Omega$ is a maximizer for $G$. To this aim, we prove that $M, \lambda_{1}$, and hence $G$, are continuous in the class of bounded open convex sets, with respect to the Hausdorff metric.

Let $\left\{\Omega_{h}\right\}_{h}$ be a sequence of bounded convex sets converging to some convex set $\Omega$, with respect to the Hausdorff metric. Under these assumptions, there exists a sequence of positive numbers $t_{h} \rightarrow 1$ such that $t_{h} \Omega_{h} \subset \Omega$ for every $h$. By comparison, such inclusion readily implies $M\left(t_{h} \Omega_{h}\right) \leq M(\Omega)$. On the other hand, by applying [8], Proposition 3.6 .1 (see also the comments below its proof), we get $M(\Omega) \leq M\left(t_{h} \Omega_{h}\right)+d_{\mathcal{H}}\left(t_{h} \Omega_{h}, \Omega\right)$, being $d_{\mathcal{H}}$ the Hausdorff distance. Thus $M\left(t_{h} \Omega_{h}\right) \rightarrow M(\Omega)$ as $h \rightarrow+\infty$. Finally, since $M\left(t_{h} \Omega_{h}\right)=t_{h}^{2} M\left(\Omega_{h}\right)$ and $t_{h} \rightarrow 1$, we conclude that $M\left(\Omega_{h}\right) \rightarrow M(\Omega)$. The continuity of $\lambda_{1}$ is more classical and can be found, e.g., in [2] or [8].

The problem of finding an optimal set is still open. We conjecture that in dimension $d=2$ the maximizer of $G$ among the convex sets is the equilateral triangle $T$, namely, for every $\Omega$ convex, $G(\Omega) \leq \frac{4}{27} \pi^{2}=G(T)$ (see the computations in the proof of Cor. 3.7 below).

### 3.3. The lower bound

The lower bound $G(\Omega) \geq 1$ is obvious: indeed, making use of (1.3) and (1.5), we get

$$
\int_{\Omega} \varphi \mathrm{d} x=\int_{\Omega} \nabla u \cdot \nabla \varphi \mathrm{~d} x=\lambda_{1}(\Omega) \int_{\Omega} u \varphi \mathrm{~d} x \leq G(\Omega) \int_{\Omega} \varphi \mathrm{d} x
$$

An alternative proof can also be found in Proposition 6 of [1], and Theorem 5 of [22].
Exploiting the same strategy used for the upper bound of $F$, we show that the constant 1 is sharp.
Theorem 3.3. In any dimension, we can find a sequence of domains $\Omega_{\epsilon}$ such that $G\left(\Omega_{\epsilon}\right) \rightarrow 1$.

Proof. Let $\epsilon>0$ be fixed and consider the perforated domain $\Omega_{\epsilon}$ defined in (2.2), obtained by removing to a given regular set $\Omega$ periodic spherical holes of period $2 \epsilon$ and radius $r_{\epsilon}$. The value of the constant $C_{0}$ appearing in (2.1) will be chosen later. Let $A_{\epsilon}: L^{2}(\Omega) \rightarrow L^{2}(\Omega)$ be the resolvent operator of the Dirichlet Laplacian on $\Omega_{\epsilon}$, which associates to $f \in L^{2}(\Omega)$ the unique solution $u \in H_{0}^{1}\left(\Omega_{\epsilon}\right)$ to $-\Delta u=f$, extended by zero outside $\Omega_{\epsilon}$.

By applying Theorem 2.5 in [9], we infer that, for every $f \in L^{2}(\Omega), A_{\epsilon}(f)$ strongly converges to $A(f)$ in $L^{2}(\Omega)$, where $A$ is the resolvent operator of $-\Delta+a$ in $H^{1}(\Omega)$ with Dirichlet boundary conditions, being $a$ (function of $C_{0}$ ) defined in (2.3). In particular, in view of Theorem 2.3.2 from [7], the eigenvalues of $A_{\epsilon}$ converge to the corresponding eigenvalue of $A$; in other words, we have

$$
\lambda_{1}\left(\Omega_{\epsilon}\right) \rightarrow \lambda_{1}(\Omega)+a
$$

as $\epsilon \rightarrow 0$. On the other hand, as already noticed in (2.10), we have $\liminf _{\epsilon} M\left(\Omega_{\epsilon}\right) \leq \frac{1}{a}$. Thus

$$
\begin{equation*}
\liminf _{\epsilon} G\left(\Omega_{\epsilon}\right)=\liminf _{\epsilon} \lambda_{1}\left(\Omega_{\epsilon}\right) M\left(\Omega_{\epsilon}\right) \leq \frac{1}{a}\left(\lambda_{1}(\Omega)+a\right)=1+\frac{\lambda_{1}(\Omega)}{a} . \tag{3.4}
\end{equation*}
$$

By choosing a suitable $C_{0}$ (small enough in the case of $d=2$ and large enough in the case of $d \geq 3$ ), the parameter $a$ can be taken arbitrarily large, so that the right-hand side of (3.4) is arbitrarily close to 1 . This fact, together with the trivial lower bound $G \geq 1$, concludes the proof.

### 3.4. The lower bound for convex sets

In the convex setting, the optimal lower bound for $G$ was provided by Payne in 1981: for every bounded convex domain $\Omega$ of $\mathbb{R}^{d}$, we have

$$
\begin{equation*}
G(\Omega)<\frac{\pi^{2}}{8} \tag{3.5}
\end{equation*}
$$

and the bound is sharp (see [16], Thm. and Eq. (3.12)). The optimality of the constant can be checked, e.g., by considering the sequence of parallelepipeds $\Omega_{n}:=(-n, n)^{d-1} \times(0,1)$. Indeed, as already seen in the proof of Theorem 2.4, by comparing the torsion function of $\Omega_{n}$ with the function $x_{d}\left(1-x_{d}\right) / 2$ we get

$$
\begin{equation*}
M\left(\Omega_{n}\right) \leq 1 / 8 \tag{3.6}
\end{equation*}
$$

on the other hand, recalling the definition (2.15) of $\psi_{n}$ and taking

$$
v\left(x_{1}, \ldots, x_{d}\right):=\sin \left(\pi x_{d}\right) \Pi_{j=1}^{d-1} \psi_{n}\left(x_{j}\right) \in H_{0}^{1}\left(\Omega_{n}\right)
$$

as test function in the variational formulation (1.6) of $\lambda_{1}\left(\Omega_{n}\right)$, we get

$$
\begin{equation*}
\lambda_{1}\left(\Omega_{n}\right) \leq \frac{\int_{\Omega_{n}}|\nabla v|^{2} \mathrm{~d} x}{\int_{\Omega_{n}} v^{2} \mathrm{~d} x}=\frac{d-1}{n-2 / 3}+\pi^{2} . \tag{3.7}
\end{equation*}
$$

From (3.6) and (3.7) we obtain the inequality

$$
G\left(\Omega_{n}\right) \leq \frac{\pi^{2}}{8}+\frac{d-1}{8(n-2 / 3)}
$$

whose right-hand side is arbitrarily close to $\pi^{2} / 8$ as $n \rightarrow+\infty$.

### 3.5. Optimality conditions via shape derivatives

In this section we derive optimality conditions by computing the first order shape derivative of $G$. Namely, given $\Omega \subset \mathbb{R}^{d}$ bounded, open, regular or convex, connected set, we study the limit (when the latter exists)

$$
G^{\prime}(\Omega, V):=\lim _{t \rightarrow 0} \frac{G\left(\Omega_{t}\right)-G(\Omega)}{t},
$$

with $\Omega_{t}:=(I+t V)(\Omega), I$ being the identity map and $V: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ an arbitrary $C^{1}$ vector field.
Recalling that $G(\Omega)=M(\Omega) \lambda_{1}(\Omega)$, if the shape derivative exists, it reads

$$
\begin{equation*}
G^{\prime}(\Omega, V)=M^{\prime}(\Omega, V) \lambda_{1}(\Omega)+M(\Omega) \lambda_{1}^{\prime}(\Omega, V) . \tag{3.8}
\end{equation*}
$$

It is well known (see, e.g., [8], Thm. 5.7.1) that

$$
\begin{equation*}
\lambda_{1}^{\prime}(\Omega, V)=-\int_{\partial \Omega}\left(\frac{\partial \varphi}{\partial n}\right)^{2} V \cdot n \mathrm{~d} \mathcal{H}^{d-1} \tag{3.9}
\end{equation*}
$$

where $\varphi$ is the (normalized) first eigenfunction and $n$ denotes the unit outer normal to $\partial \Omega$. Remark that $\frac{\partial \varphi}{\partial n}$ is well defined as soon as $\Omega$ is regular or convex, since $\varphi \in H^{2}(\Omega)$ in that case.

The computation of $M^{\prime}$ is more delicate and requires additional assumptions.

Proposition 3.4. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open convex set. Then, for every $V \in C^{1}\left(\mathbb{R}^{d}\right)$, the shape derivative of $M$ at $\Omega$ in direction $V$ exists and is given by

$$
M^{\prime}(\Omega, V)=u^{\prime}\left(x_{0}\right)
$$

where $x_{0}$ is a maximum point for $u$ and $u^{\prime}$ is the solution of

$$
\left\{\begin{array}{l}
\Delta u^{\prime}=0 \quad \text { in } \quad \Omega \\
u^{\prime}+\nabla u \cdot V \in H_{0}^{1}(\Omega)
\end{array}\right.
$$

Proof. Let $u_{t}$ denote the torsion function of $\Omega_{t}$, for $t>0$, and consider the function $\psi: \mathbb{R} \times \mathbb{R}^{d} \rightarrow \mathbb{R}^{d}$ defined as

$$
(t, x) \mapsto \psi(t, x):=\nabla u_{t}(x)
$$

By optimality of $x_{0}$ for $u$, we have $\psi\left(0, x_{0}\right)=0$. Moreover, the matrix $D_{x} \psi\left(0, x_{0}\right)$ is invertible: indeed, setting $v:=\sqrt{u}$, we have

$$
D_{x} \psi\left(0, x_{0}\right)=\operatorname{Hess}_{u}\left(x_{0}\right)=2 \sqrt{M(\Omega)} \operatorname{Hess}_{v}\left(x_{0}\right)
$$

and the matrix $H_{e s s}^{v}$ is negative definite everywhere in $\Omega$ (see [11]). Thus, by the implicit function theorem, we infer that, in a neighborhood of $x_{0}$, for $t$ small enough, there exists a unique $x_{t}$ such that $\nabla u_{t}\left(x_{t}\right)=0$; furthermore, $t \mapsto x_{t}$ is differentiable.

We claim that, for every $t$,

$$
\begin{equation*}
u_{t}\left(x_{t}\right)=M\left(\Omega_{t}\right) \tag{3.10}
\end{equation*}
$$

and that, as $t \rightarrow 0$,

$$
\begin{align*}
& \frac{u_{t}\left(x_{0}\right)-u\left(x_{0}\right)}{t} \rightarrow u^{\prime}\left(x_{0}\right)  \tag{3.11}\\
& \frac{u_{t}\left(x_{t}\right)-u_{t}\left(x_{0}\right)}{t} \rightarrow 0 \tag{3.12}
\end{align*}
$$

Once proved the claims we are done, indeed we have

$$
\frac{M\left(\Omega_{t}\right)-M(\Omega)}{t}=\frac{u_{t}\left(x_{t}\right)-u\left(x_{0}\right)}{t}=\frac{u_{t}\left(x_{t}\right)-u_{t}\left(x_{0}\right)}{t}+\frac{u_{t}\left(x_{0}\right)-u\left(x_{0}\right)}{t} \rightarrow u^{\prime}\left(x_{0}\right)
$$

which concludes the proof.
Let us show (3.10). As already pointed out in the proof of Theorem 3.2, exploiting [8] (Prop. 3.6.1) we infer that the sequence $u_{t}$ uniformly converges to $u$, extended to zero in a larger common domain independent of $t$. In particular, every sequence of maximizers $y_{t}$ for $u_{t}$ converges to the unique minimizer $x_{0}$ of $u$ (see, e.g., [4], Prop. 5.2, 7.24 applied to the sequence $-u_{t}$ ). Finally, since $x_{t}$ are isolated critical points in a neighborhood of $x_{0}$, they must coincide with $y_{t}$, so that $u_{t}\left(x_{t}\right)=u_{t}\left(y_{t}\right)=M\left(\Omega_{t}\right)$.

Assertion (3.11) follows by applying the mean value property to the harmonic functions $u_{t}-u$ and $u^{\prime}$ : choose $R>0$ such that $B_{R}\left(x_{0}\right) \subset \Omega_{t}$ for every $t \ll 1$, then we have

$$
\left|\frac{u_{t}\left(x_{0}\right)-u\left(x_{0}\right)}{t}-u^{\prime}\left(x_{0}\right)\right|=\left|f_{B_{R}\left(x_{0}\right)} \frac{u_{t}(x)-u(x)}{t}-u^{\prime}\left(x_{0}\right) \mathrm{d} x\right| \leq C\left\|\left(u_{t}-u\right) / t-u^{\prime}\right\|_{L^{2}\left(B_{R}\left(x_{0}\right)\right)}
$$

The right-hand side vanishes as $t \rightarrow 0$, since the map $t \mapsto u_{t} \in L^{2}\left(\mathbb{R}^{d}\right)$ is differentiable at 0 with derivative $\frac{\mathrm{d}}{\mathrm{d} t} u \bigsqcup_{t=0}=u^{\prime}$ (see, for instance, [8], Thm. 5.3.1).

Similarly, property (3.12) follows by combining the mean value property of $\nabla u_{t}$, the differentiability of $t \mapsto x_{t}$, and the strong convergence of $u_{t}$ to $u$ in $H^{1}\left(\mathbb{R}^{d}\right)$ :

$$
\begin{aligned}
\frac{u_{t}\left(x_{t}\right)-u_{t}\left(x_{0}\right)}{t} & =\nabla u_{t}\left(\xi_{t}\right) \cdot \frac{x_{t}-x_{0}}{t}=\left(f_{B_{R}\left(\xi_{t}\right)} \nabla u_{t}(x) \mathrm{d} x\right) \cdot \frac{x_{t}-x_{0}}{t} \\
& \longrightarrow \nabla u\left(x_{0}\right) \cdot v_{0}=0, \quad \text { as } t \rightarrow 0
\end{aligned}
$$

with $\xi_{t}$ a suitable intermediate point between $x_{0}$ and $x_{t}, R$ a positive radius such that $B_{R}\left(\xi_{t}\right) \subset \Omega_{t}$ for every $t \ll 1$, and $v_{0}$ the derivative $\frac{\mathrm{d}}{\mathrm{d} t} x_{t}\left\lfloor_{t=0}\right.$.

Remark 3.5. We point out that Proposition 3.4 is valid in a more general setting, when the torsion function $u_{\Omega}$ belongs to $W^{1, \infty}(\Omega)$ and has a unique maximum point which is non degenerate (i.e. invertible Hessian at the point). These properties are ensured by the convexity of $\Omega$.

Theorem 3.6. Let $\Omega \subset \mathbb{R}^{d}$ be a bounded open convex set. Then, for every $V \in C^{1}\left(\mathbb{R}^{d}\right)$, the shape derivative of $G$ at $\Omega$ in direction $V$ exists and is given by

$$
\begin{equation*}
G^{\prime}(\Omega, V)=\int_{\partial \Omega}\left[\lambda_{1}(\Omega) \frac{\partial u}{\partial n} \frac{\partial \phi_{x_{0}}}{\partial n}-M(\Omega)\left(\frac{\partial \varphi}{\partial n}\right)^{2}\right] V \cdot n \mathrm{~d} \mathcal{H}^{d-1} \tag{3.13}
\end{equation*}
$$

where $x_{0} \in \Omega$ is a maximum point of $u$ and $\phi_{x_{0}}$ is the (Green function) solution of

$$
\begin{cases}-\Delta \phi_{x_{0}}=\delta_{x_{0}} & \text { in } \quad \mathcal{D}^{\prime}(\Omega) \\ \phi_{x_{0}}=0 & \text { on } \partial \Omega\end{cases}
$$

Proof. First, we rewrite in terms of $\phi_{x_{0}}$ the derivative $M^{\prime}(\Omega, V)$, whose existence is ensured by Proposition 3.4:

$$
\begin{equation*}
M^{\prime}(\Omega, V)=\left\langle-\Delta \phi_{x_{0}}, u^{\prime}\right\rangle=-\int_{\partial \Omega} u^{\prime} \frac{\partial \phi_{x_{0}}}{\partial n} \mathrm{~d} \mathcal{H}^{d-1}=\int_{\partial \Omega} \frac{\partial u}{\partial n} \frac{\partial \phi_{x_{0}}}{\partial n} V \cdot n \mathrm{~d} \mathcal{H}^{d-1} \tag{3.14}
\end{equation*}
$$

Formula (3.13) follows by combining (3.8), (3.9), and (3.14).

As a consequence of Theorem 3.6 we obtain the following optimality condition: if $\Omega \subset \mathbb{R}^{d}$ bounded open convex set is a critical shape for $G$, then

$$
\begin{equation*}
\lambda_{1}(\Omega) \frac{\partial u}{\partial n} \frac{\partial \phi_{x_{0}}}{\partial n}-M(\Omega)\left(\frac{\partial \varphi}{\partial n}\right)^{2}=0 \quad \text { a.e. on } \partial \Omega . \tag{3.15}
\end{equation*}
$$

Now we partially answer a question raised in [22] where the author asked whether the disk could be the maximizer for $G$.

Corollary 3.7. The equilateral triangle gives a better value than the disk; however it is not a critical shape for $G$.

Proof. Let $T \subset \mathbb{R}^{2}$ be the equilateral triangle with side of length 1 and vertices in $(-1 / 2,-1 /(2 \sqrt{3}))$, $(1 / 2,-1 /(2 \sqrt{3}))$, and $(0,1 / \sqrt{3})$, so that the center is at the origin. In this case, $u, \varphi, M$, and $\lambda_{1}$ can be
explicitly computed and read

$$
\begin{aligned}
u(x, y) & =\frac{1}{2 \sqrt{3}}\left(y+\frac{1}{2 \sqrt{3}}\right)\left(y-\sqrt{3} x-\frac{1}{\sqrt{3}}\right)\left(y+\sqrt{3} x-\frac{1}{\sqrt{3}}\right) \\
& =\frac{1}{2 \sqrt{3}}\left(y^{3}-3 x^{2} y-\frac{\sqrt{3}}{2} y^{2}-\frac{\sqrt{3}}{2} x^{2}+\frac{1}{6 \sqrt{3}}\right) \\
\varphi(x, y) & =\left(\frac{2}{\sqrt{3}}\right)^{3 / 2}\left[\sin \left(\frac{4 \pi}{3}(1-\sqrt{3} y)\right)-2 \cos (2 \pi x) \sin \left(\frac{2 \pi}{3}(1-\sqrt{3} y)\right)\right] \\
M(T) & =u(0,0)=\frac{1}{36}, \quad \lambda_{1}(T)=\frac{16}{3} \pi^{2}
\end{aligned}
$$

Therefore $G(T)=4 \pi^{2} / 27 \simeq 1.4622$ while, for the unit disk $\mathbb{D}$, we have $G(\mathbb{D})=j_{0,1}^{2} / 4 \simeq 1.4458$ which proves the first part of the claim.

Now assume by contradiction that $T$ is a critical shape for $G$. Since the normal derivative of $u$ on $\partial T$ is never zero except at the vertices (where both $\nabla u$ and $\nabla \varphi$ vanish), we may recast the optimality condition (3.15) as

$$
\begin{equation*}
\frac{\partial \phi_{0}}{\partial n}=h \quad \text { a.e. on } \partial T \tag{3.16}
\end{equation*}
$$

with $h:=M(\Omega)\left(\lambda_{1}(\Omega)\right)^{-1}\left(\frac{\partial u}{\partial n}\right)^{-1}\left(\frac{\partial \varphi}{\partial n}\right)^{2}$. In particular, if we multiply both sides by an arbitrary harmonic function $w$ and integrate over $\partial T$, we obtain

$$
\begin{equation*}
-w(0)=\int_{\partial T} w h \mathrm{~d} \mathcal{H}^{1} \tag{3.17}
\end{equation*}
$$

By taking as test functions $w \equiv 1$ and $w=\operatorname{Re}\left(z^{6}\right)$, we get

$$
\left\{\begin{array}{l}
\int_{\partial T} h \mathrm{~d} \mathcal{H}^{1}=-1 \\
\int_{\partial T} \operatorname{Re}\left(z^{6}\right) h \mathrm{~d} \mathcal{H}^{1}=0
\end{array}\right.
$$

Exploiting the symmetry of the domain and of the functions involved, these conditions can be rephrased as follows:

$$
\left\{\begin{array}{l}
\int_{\Sigma} h \mathrm{~d} \mathcal{H}^{1}=-\frac{1}{3} \\
\int_{\Sigma} \operatorname{Re}\left(z^{6}\right) h \mathrm{~d} \mathcal{H}^{1}=0
\end{array}\right.
$$

where $\Sigma$ denotes the basis of the triangle, i.e. the segment $\Sigma=[-1 / 2,1 / 2] \times\{-1 /(2 \sqrt{3})\}$. This system may be rewritten as

$$
\left\{\begin{array}{l}
\sigma:=\int_{0}^{1 / 2} \frac{(1+\cos (2 \pi x))^{2}}{x^{2}-1 / 4} \mathrm{~d} x=-\frac{27}{8}  \tag{3.18}\\
\tau:=\int_{0}^{1 / 2} P(x) \frac{(1+\cos (2 \pi x))^{2}}{x^{2}-1 / 4} \mathrm{~d} x=0
\end{array}\right.
$$

with

$$
P(x):=\operatorname{Re}\left(z^{6}\right)\left\lfloor_{\Sigma}=\left(x^{6}-15 x^{4} y^{2}+15 x^{2} y^{4}-y^{6}\right)\left\lfloor_{y=-1 /(2 \sqrt{3})}=x^{6}-\frac{5}{4} x^{4}+\frac{5}{48} x^{2}-\frac{1}{1728}\right.\right.
$$

Making use of the factorization

$$
P(x)=\left(x^{2}-\frac{1}{4}\right)\left(x^{4}-x^{2}-\frac{7}{48}\right)-\frac{1}{27},
$$

we obtain

$$
\tau=\int_{0}^{1 / 2}\left(x^{4}-x^{2}-\frac{7}{48}\right)(1+\cos (2 \pi x))^{2} \mathrm{~d} x-\frac{\sigma}{27}
$$

It is easy to check that

$$
\begin{aligned}
\int_{0}^{1 / 2}(1+\cos (2 \pi x))^{2} \mathrm{~d} x & =\frac{3}{4} \\
\int_{0}^{1 / 2} x^{2}(1+\cos (2 \pi x))^{2} \mathrm{~d} x & =\frac{1}{16}-\frac{15}{32 \pi^{2}} \\
\int_{0}^{1 / 2} x^{4}(1+\cos (2 \pi x))^{2} \mathrm{~d} x & =\frac{3}{320}-\frac{15}{64 \pi^{2}}+\frac{189}{128 \pi^{4}} .
\end{aligned}
$$

Thus we get

$$
\begin{aligned}
\tau+\frac{\sigma}{27} & =\frac{3}{320}-\frac{15}{64 \pi^{2}}+\frac{189}{128 \pi^{4}}-\frac{1}{16}+\frac{15}{32 \pi^{2}}-\frac{7}{48} \cdot \frac{3}{4} \\
& =\frac{3}{320}-\frac{1}{16}-\frac{7}{64}+\frac{15}{32 \pi^{2}}-\frac{15}{64 \pi^{2}}+\frac{189}{128 \pi^{4}} \\
& =-\frac{13}{80}+\frac{15}{64 \pi^{2}}+\frac{189}{128 \pi^{4}} .
\end{aligned}
$$

Since $\pi$ is not algebraic, the last relation is in contradiction with (3.18), which in turn implies $\tau+\sigma / 27=-1 / 8$. Therefore we conclude that the equilateral triangle is not a critical shape.

Remark 3.8. We point out that the choice of any test function of the form $w=\operatorname{Re} z^{n}$, for $n=1, \ldots, 5$, in (3.17) does not provide any contradiction. Moreover the numerical values of $\sigma$ and $\tau$ defined in the above proof are not so far of what appears in (3.18). In some sense, the equilateral triangle is not far from being a critical point.

We conclude the Section with the proof of Proposition 3.1.
Proof of Proposition 3.1. Assume by contradiction that $\Omega$ is a maximizer for $G$. Given $x \in \Omega$ and $\epsilon>0$ a small parameter, we set $\Omega_{\epsilon}(x)$ the perforated domain $\Omega \backslash \bar{B}_{\epsilon}(x)$ and we denote by $u_{\epsilon, x}$ the associated torsion function.

In the limit as $\epsilon \rightarrow 0$, we have the following asymptotic expansions for $\lambda_{1}\left(\Omega_{\epsilon}(x)\right)$ and $u_{\epsilon, x}$ in terms of $\lambda_{1}(\Omega)$ and $u(c f .[6,14]$, (Chap. 8)):

$$
\begin{aligned}
& \lambda_{1}\left(\Omega_{\epsilon}(x)\right)=\left\{\begin{array}{llc}
\lambda_{1}(\Omega)+\frac{2 \pi}{|\log \epsilon|} \varphi^{2}(x)+o(1 /|\log \epsilon|) & \text { if } & d=2 \\
\lambda_{1}(\Omega)+\epsilon^{d-2}(d-2) \omega_{d} \varphi^{2}(x)+o\left(\epsilon^{d-2}\right) & \text { if } & d>2
\end{array}\right. \\
& u_{\epsilon, x}(y)=\left\{\begin{array}{lcc}
u(y)-\frac{2 \pi}{|\log \epsilon|} u(x) \phi_{x}(y)+o(1 /|\log \epsilon|) & \text { if } & d=2 \\
u(y)-\epsilon^{d-2}(d-2) \omega_{d} u(x) \phi_{x}(y)+o\left(\epsilon^{d-2}\right) & \text { if } & d>2,
\end{array}\right.
\end{aligned}
$$

where $\omega_{d}$ is the measure of the $(d-1)$-sphere and $\phi_{x}(y)$ is the Green function of the Laplace operator vanishing on the boundary of $\Omega$. In particular, choosing $x$ different from the maximum point $x_{0}$ of $u$ and evaluating $u_{\epsilon, x}$ at $x_{0}$, we obtain

$$
M\left(\Omega_{\epsilon}(x)\right) \geq \begin{cases}M(\Omega)-\frac{2 \pi}{|\log \epsilon|} u(x) \phi_{x}\left(x_{0}\right)+o(r(\epsilon)) & \text { if } \quad d=2 \\ M(\Omega)-\epsilon^{d-2}(d-2) \omega_{d} u(x) \phi_{x}\left(x_{0}\right)+o(r(\epsilon)) & \text { if } \quad d>2\end{cases}
$$

where $r(\epsilon)=1 /|\log \epsilon|$ if $d=2$ and $\epsilon^{d-2}$ otherwise. In particular, we obtain the lower bound

$$
\begin{equation*}
G\left(\Omega_{\epsilon}(x)\right) \geq G(\Omega)+R(x) r(\epsilon)+o(r(\epsilon)) \tag{3.19}
\end{equation*}
$$

with

$$
\begin{equation*}
R(x):=M(\Omega) \varphi^{2}(x)-\lambda_{1}(\Omega) u(x) \phi_{x_{0}}(x) \tag{3.20}
\end{equation*}
$$

(here we have used the symmetry of the Green function: $\left.\phi_{x}\left(x_{0}\right)=\phi_{x_{0}}(x)\right)$. To get a contradiction, it suffices to find a point $x$ in which $R(x)>0$. Taking $x$ close to the boundary, say $x=x_{1}-\delta n\left(x_{1}\right)$ for some $x_{1} \in \partial \Omega$ and $0<\delta \ll 1$, and recalling that $\varphi$ vanishes on $\partial \Omega$, we may write

$$
\varphi(x)=-\delta \frac{\partial \varphi}{\partial n}\left(x_{1}\right)+\frac{\delta^{2}}{2} \frac{\partial^{2} \varphi}{\partial n^{2}}\left(x_{1}\right)+o\left(\delta^{2}\right)
$$

Furthermore, by combining (1.5) with the relation $\Delta \varphi=\Delta_{\partial \Omega} \varphi+H_{\partial \Omega} \frac{\partial \varphi}{\partial n}+\frac{\partial^{2} \varphi}{\partial n^{2}}$ on $\partial \Omega$, we get

$$
\frac{\partial^{2} \varphi}{\partial n^{2}}\left(x_{1}\right)=-H_{\partial \Omega}\left(x_{1}\right) \frac{\partial \varphi}{\partial n}\left(x_{1}\right)
$$

Arguing in the same way for $u$ and $\phi_{x_{0}}$, we obtain the developments

$$
\begin{aligned}
u(x) & =-\delta \frac{\partial u}{\partial n}\left(x_{1}\right)+\frac{\delta^{2}}{2} \frac{\partial^{2} u}{\partial n^{2}}\left(x_{1}\right)+o\left(\delta^{2}\right) \\
\phi_{x_{0}}(x) & =-\delta \frac{\partial \phi_{x_{0}}}{\partial n}\left(x_{1}\right)+\frac{\delta^{2}}{2} \frac{\partial^{2} \phi_{x_{0}}}{\partial n^{2}}\left(x_{1}\right)+o\left(\delta^{2}\right),
\end{aligned}
$$

and the equalities

$$
\begin{aligned}
\frac{\partial^{2} u}{\partial n^{2}}\left(x_{1}\right) & =-1-H_{\partial \Omega}\left(x_{1}\right) \frac{\partial u}{\partial n}\left(x_{1}\right) \\
\frac{\partial^{2} \phi_{x_{0}}}{\partial n^{2}}\left(x_{1}\right) & =-H_{\partial \Omega}\left(x_{1}\right) \frac{\partial \phi_{x_{0}}}{\partial n}\left(x_{1}\right)
\end{aligned}
$$

In view of these computations, we infer that

$$
R(x)=\left(\delta^{2}+\delta^{3} H_{\partial \Omega}\left(x_{1}\right)\right)\left[M(\Omega)\left(\frac{\partial \varphi}{\partial n}\right)^{2}-\lambda_{1}(\Omega) \frac{\partial u}{\partial n} \frac{\partial \phi_{x_{0}}}{\partial n}\right]\left(x_{1}\right)-\frac{\delta^{3}}{2} \lambda_{1}(\Omega) \frac{\partial \phi_{x_{0}}}{\partial n}\left(x_{1}\right)+o\left(\delta^{3}\right)
$$

By optimality of $\Omega$, the equality (3.15) holds true at $x_{1}$, so that

$$
R(x)=-\frac{\delta^{3}}{2} \lambda_{1}(\Omega) \frac{\partial \phi_{x_{0}}}{\partial n}\left(x_{1}\right)+o\left(\delta^{3}\right)
$$

By the Hopf's principle $\frac{\partial \phi_{x_{0}}}{\partial n}$ is strictly negative on the boundary $\partial \Omega$, therefore $R(x)$ is strictly positive. In particular, in view of (3.19) we conclude that, if $x$ is close enough to the boundary,

$$
G\left(\Omega_{\epsilon}(x)\right)>G(\Omega)
$$

which is in contradiction with the maximality of $\Omega$.

Final comment. In the conference "Shape optimization and isoperimetric and functional inequalities", held in CIRM - Luminy (France) on November 21-25, 2016, we had the occasion to discuss with Buttazzo and Velichkov, who kindly gave us the proof of Theorem 2.2.

In the same conference, we announced our results to M. van den Berg, since they disproved some conjectures he gave in [22]. Few weeks later, M. van den Berg posted [23] on ArXiv, where he gave a different proof of our Theorem 3.3. Finally, we point out that the recent contribution [5] generalizes some of our results to more general operators.

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    1 Institut Élie Cartan de Lorraine, UMR 7502, Université de Lorraine CNRS, Vandoeuvre les Nancy Cedex, France.
    2 Département de Mathématiques, Université Laval Québec, Québec, QC G1V 0A6, Canada.
    *Corresponding author: ilaria.lucardesi@univ-lorraine.fr

