

On the global maximum of the solution to a stochastic heat equation with compact-support initial data¹

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Abstract. Consider a stochastic heat equation $\partial_t u = \kappa \partial_{xx}^2 u + \sigma(u)\dot{w}$ for a space–time white noise \dot{w} and a constant $\kappa > 0$. Under some suitable conditions on the initial function u_0 and σ , we show that the quantities

 $\limsup_{t \to \infty} t^{-1} \sup_{x \in \mathbf{R}} \ln \mathbb{E}(|u_t(x)|^2) \quad \text{and} \quad \limsup_{t \to \infty} t^{-1} \ln \mathbb{E}(\sup_{x \in \mathbf{R}} |u_t(x)|^2)$

are equal, as well as bounded away from zero and infinity by explicit multiples of $1/\kappa$. Our proof works by demonstrating quantitatively that the peaks of the stochastic process $x \mapsto u_t(x)$ are highly concentrated for infinitely-many large values of t. In the special case of the parabolic Anderson model – where $\sigma(u) = \lambda u$ for some $\lambda > 0$ – this "peaking" is a way to make precise the notion of physical intermittency.

Résumé. Nous considérons l'équation de la chaleur stochastique $\partial_t u = \kappa \partial_{xx}^2 u + \sigma(u)\dot{w}$ avec un bruit blanc spatio-temporel \dot{w} et une constante $\kappa > 0$. Sous des conditions adéquates sur la condition initiale u_0 et sur σ , nous montrons que les quantités

$$\limsup_{t \to \infty} t^{-1} \sup_{x \in \mathbf{R}} \ln \mathbb{E}(|u_t(x)|^2) \quad \text{et} \quad \limsup_{t \to \infty} t^{-1} \ln \mathbb{E}\left(\sup_{x \in \mathbf{R}} |u_t(x)|^2\right)$$

sont égales. Par ailleurs, nous les bornons inférieurement et supérieurement par des constantes strictement positives et finies dépendant explicitement de $1/\kappa$. Nos démonstrations reposent sur la preuve quantitative de la forte concentration des pics du processus $x \mapsto u_t(x)$ pour de grandes valeurs de *t* infiniment nombreuses. Dans le cas particulier du modèle d'Anderson parabolique-où $\sigma(u) = \lambda u$ pour un $\lambda > 0$ – ce phénomène de pics est une façon de préciser la notion physique d'intermittence.

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1. Introduction

We consider the stochastic heat equation,

$$\frac{\partial u_t(x)}{\partial t} = \kappa \frac{\partial^2 u_t(x)}{\partial x^2} + \sigma \left(u_t(x) \right) \dot{w}(t,x) \quad \text{for } t > 0 \text{ and } x \in \mathbf{R},$$
(1.1)

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where $\kappa > 0$ is fixed, $\sigma : \mathbf{R} \to \mathbf{R}$ is Lipschitz continuous with $\sigma(0) = 0$, \dot{w} denotes space-time white noise, and the initial data $u_0: \mathbf{R} \to \mathbf{R}$ is nonrandom. There are several areas to which (1.1) has deep and natural connections; perhaps chief among them are the stochastic Burgers' equation [10] and the celebrated KPZ equation of statistical mechanics [11,12]; see also [13], Chapter 9.

It is well known that (1.1) has an almost-surely unique, adapted and continuous solution $\{u_t(x)\}_{t \ge 0, x \in \mathbb{R}}$ ([5], Theorem 6.4, p. 26). In addition, the condition that $\sigma(0) = 0$ implies that if $u_0 \in L^2(\mathbf{R})$, then $u_t \in L^2(\mathbf{R})$ a.s. for all $t \ge 0$; see Dalang and Mueller [6]. Note that our conditions on σ ensure that

$$\left|\sigma(u)\right| \le \operatorname{Lip}_{\sigma}|u| \quad \text{for all } u \in \mathbf{R},\tag{1.2}$$

where

$$\operatorname{Lip}_{\sigma} := \sup_{-\infty < x < x' < \infty} \left| \frac{\sigma(x) - \sigma(x')}{x - x'} \right|.$$
(1.3)

Our goal is to establish the following general growth estimate.

Theorem 1.1. Suppose there exists $L_{\sigma} \in (0, \infty)$ such that $|\sigma(u)| \ge L_{\sigma}|u|$ for all $u \in \mathbf{R}$. Suppose also that $u_0 \neq 0$ is Hölder-continuous of order $\geq 1/2$, nonnegative, and supported in [-K, K] for some finite K > 0. Then, (1.1) has an almost-surely unique, continuous and adapted solution $\{u_t(x)\}_{t\geq 0, x\in \mathbf{R}}$ such that $u_t \in L^2(\mathbf{R})$ a.s. for all $t\geq 0$, and

$$\frac{L_{\sigma}^4}{8\kappa} \le \limsup_{t \to \infty} t^{-1} \sup_{x \in \mathbf{R}} \ln \mathbb{E}(|u_t(x)|^2) = \limsup_{t \to \infty} t^{-1} \ln \mathbb{E}\left(\sup_{x \in \mathbf{R}} |u_t(x)|^2\right) \le \frac{\operatorname{Lip}_{\sigma}^4}{8\kappa}$$

Because of Mueller's comparison principle [14] (see also [7,16]), the nonnegativity of u_0 implies that $\sup_{t,x} E(|u_t(x)|) = \sup_{t,x} E(u_t(x))$, and this quantity has to be finite because u_0 is bounded; confer with (1.5). Consequently,

$$\sup_{x \in \mathbf{R}} \|u_t(x)\|_{L^1(\mathbf{P})} \ll \sup_{x \in \mathbf{R}} \|u_t(x)\|_{L^2(\mathbf{P})} \quad \text{as } t \to \infty.$$
(1.4)

When $\text{Lip}_{\sigma} = L_{\sigma}$, (1.1) becomes the well-studied parabolic Anderson model [1,3]. And (1.4) makes precise the physical notion that the solution to (1.1) concentrates near "very high peaks" [1,3,11,12].

In order to explain the idea behind our proof, we introduce the following.

Definition 1.2. We say that a continuous random field $f := \{f(t, x)\}_{t>0, x \in \mathbb{R}}$ has effectively-compact support [in the spatial variable x] if there exists a nonrandom measurable function $p: \mathbf{R}_+ \to \mathbf{R}_+$ of at-most polynomial growth such that:

- (a) $\limsup_{t\to\infty} t^{-1} \ln \int_{|x| \le p(t)} \mathbb{E}(|f(t, x)|^2) \, dx > 0$ and (b) $\limsup_{t\to\infty} t^{-1} \ln \int_{|x| > p(t)} \mathbb{E}(|f(t, x)|^2) \, dx < 0.$

We might refer to the function p as the radius of effective support of f.

One of the ideas here is to use Mueller's comparison principle [14] to compare $\sup_{x \in \mathbf{R}} |u_t(x)|$ with the $L^2(\mathbf{R})$ norm of $x \mapsto u_t(x)$, which is easier to analyze. We carry these steps out in Lemma 3.3. We also appeal to the fact that the compact-support property of u_0 implies that $u_1(x)$ has an effectively-compact support [Proposition 3.7]. This can be interpreted as a kind of optimal regularity theorem. However, these matters need to be handled delicately, as "effectively compact" cannot be replaced by "compact"; see Mueller [14].

Our method for establishing an effectively-compact support property is motivated strongly by ideas of Mueller and Perkins [15]. In the cases that $u_t(x)$ denotes the density of some particles at x at time t, our effectively-compact support property implies that most of the particles accumulate on a very small set. This method might appeal to the reader who is interested in mathematical descriptions of physical intermittency.

$$u_t(x) = (p_t * u_0)(x) + \int_0^t \int_{-\infty}^\infty p_{t-s}(y-x)\sigma(u_s(y))w(\mathrm{d} s \,\mathrm{d} y), \tag{1.5}$$

where $p_{\tau}(z) := (4\kappa\tau\pi)^{-1/2} \exp(-z^2/(4\kappa\tau))$ denotes the heat kernel corresponding to the operator $\kappa \partial^2/\partial x^2$, and the stochastic integral is understood in the sense of Walsh [17]. Some times we write $||X||_p$ in place of $\{E(|X|^p)\}^{1/p}$.

2. A preliminary result

As mentioned in the Introduction, the strategy behind our proof of Theorem 1.1 is to relate the global maximum of the solution to a "closed-form quantity" that resembles $\sup_{x} |u_t(x)|$ for large values of t. That closed-form quantity turns out to be the $L^2(\mathbf{R})$ -norm of $x \mapsto u_t(x)$. Our next result analyses the growth of the mentioned closed-form quantity. We related it to $\sup_{x} |u_t(x)|$ in the next section. The methods of this section follow closely the classical ideas of Choquet and Deny [4] that were developed in a deterministic setting.

Theorem 2.1. Suppose $\sigma : \mathbf{R} \to \mathbf{R}$ is Lipschitz continuous, $\sigma(0) = 0$, and there exists $\mathcal{L}_{\sigma} \in (0, \infty)$ such that $\mathcal{L}_{\sigma}|u| \le |\sigma(u)|$ for all $u \in \mathbf{R}$. If $u_0 \in L^2(\mathbf{R})$ and $u_0 \not\equiv 0$, then (1.1) has an almost-surely unique, continuous and adapted solution $\{u_t(x)\}_{t\geq 0, x\in \mathbf{R}}$ such that $u_t \in L^2(\mathbf{R})$ a.s. for all $t \ge 0$, and

$$\frac{\mathcal{L}_{\sigma}^{4}}{8\kappa} \leq \limsup_{t \to \infty} t^{-1} \ln \mathbb{E} \left(\|u_{t}\|_{L^{2}(\mathbf{R})}^{2} \right) \leq \frac{\operatorname{Lip}_{\sigma}^{4}}{8\kappa}.$$
(2.1)

Proof. It suffices to establish (2.1). Note that

$$E(|u_t(x)|^2) = |(p_t * u_0)(x)|^2 + \int_0^t ds \int_{-\infty}^\infty dy E(|\sigma(u_s(y))|^2) \cdot |p_{t-s}(y-x)|^2.$$
(2.2)

Because $|\sigma(u)| \ge L_{\sigma}|u|$,

$$E(\|u_{t}\|_{L^{2}(\mathbf{R})}^{2}) = \|p_{t} * u_{0}\|_{L^{2}(\mathbf{R})}^{2} + \int_{0}^{t} ds \int_{-\infty}^{\infty} dy E(|\sigma(u_{s}(y))|^{2}) \cdot \|p_{t-s}\|_{L^{2}(\mathbf{R})}^{2}$$

$$\geq \|p_{t} * u_{0}\|_{L^{2}(\mathbf{R})}^{2} + L_{\sigma}^{2} \cdot \int_{0}^{t} E(\|u_{s}\|_{L^{2}(\mathbf{R})}^{2}) \cdot \|p_{t-s}\|_{L^{2}(\mathbf{R})}^{2} ds.$$
(2.3)

We can multiply the preceding by $exp(-\lambda t)$ throughout and integrate [dt] to find that if

$$U(\lambda) := \int_0^\infty e^{-\lambda t} \mathbb{E}\left(\left\|u_t\right\|_{L^2(\mathbf{R})}^2\right) \mathrm{d}t,\tag{2.4}$$

then

$$U(\lambda) \ge \int_0^\infty e^{-\lambda t} \|p_t * u_0\|_{L^2(\mathbf{R})}^2 dt + L_{\sigma}^2 \cdot U(\lambda) \cdot \int_0^\infty e^{-\lambda t} \|p_t\|_{L^2(\mathbf{R})}^2 dt.$$
(2.5)

According to Plancherel's theorem, the following holds for all finite Borel measures μ on **R**:

$$\|p_t * \mu\|_{L^2(\mathbf{R})}^2 = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left|\hat{\mu}(\xi)\right|^2 e^{-2\kappa t\xi^2} d\xi.$$
(2.6)

Therefore, Tonelli's theorem ensures that

$$\int_0^\infty e^{-\lambda t} \|p_t * \mu\|_{L^2(\mathbf{R})}^2 dt = \frac{1}{2\pi} \int_0^\infty \frac{|\hat{\mu}(\xi)|^2}{\lambda + 2\kappa\xi^2} d\xi.$$
(2.7)

We apply this identity twice in (2.3): Once with $\mu := \delta_0$; and once with $d\mu/dx := u_0$. This leads us to the following.

$$U(\lambda) \geq \frac{1}{2\pi} \int_0^\infty \frac{|\hat{u}_0(\xi)|^2}{\lambda + 2\kappa\xi^2} \,\mathrm{d}\xi + \mathcal{L}_{\sigma}^2 \cdot U(\lambda) \cdot \frac{1}{2\pi} \int_0^\infty \frac{\mathrm{d}\xi}{\lambda + 2\kappa\xi^2} \\ = \frac{1}{2\pi} \int_0^\infty \frac{|\hat{u}_0(\xi)|^2}{\lambda + 2\kappa\xi^2} \,\mathrm{d}\xi + \mathcal{L}_{\sigma}^2 \cdot U(\lambda) \cdot \frac{1}{2\sqrt{2\kappa\lambda}}.$$

$$(2.8)$$

Since $u_0 \neq 0$, the first [Fourier] integral is strictly positive. Consequently, the above recursive relation shows that $U(\lambda) = \infty$ if $\lambda \leq L_{\sigma}^4/(8\kappa)$. This and a real-variable argument together imply the first inequality in (2.1). Indeed, we follow the argument in [8] in this way: Suppose, to the contrary, that the first inequality in (2.1) failed. This means that for all $\varepsilon > 0$ there exists $t_0 > 0$ such that for all $t > t_0$,

$$\mathbb{E}\left(\left\|u_{t}\right\|_{L^{2}(\mathbf{R})}^{2}\right) \leq \exp\left(t\left\{\frac{\mathbf{L}_{\sigma}^{4}}{8\kappa} - \varepsilon\right\}\right).$$
(2.9)

We multiply this by $\exp(-\lambda t)$ and integrate [dt] to deduce that $U(\lambda) < \infty$ for all $\lambda > L_{\sigma}^4/(8\kappa) - \varepsilon$. And this contradicts the earlier finding that $U(\lambda) = \infty$ for all $\lambda \le L_{\sigma}^4/(8\kappa)$.

For the other bound we use a Picard-iteration argument in order to obtain an a priori estimate. Let $u_t^{(0)}(x) := u_0(x)$ and iteratively define

$$u_t^{(n+1)}(x) := (p_t * u_0)(x) + \int_0^t \int_{-\infty}^\infty p_{t-s}(y-x)\sigma(u_s^{(n)}(y))w(\mathrm{d} s \,\mathrm{d} y).$$
(2.10)

Since $||p_t * u_0||_{L^2(\mathbf{R})} \le ||u_0||_{L^2(\mathbf{R})}$ and $|\sigma(u)| \le \operatorname{Lip}_{\sigma}|u|$, Hölder's inequality yields

$$\mathbb{E}\left(\left\|u_{t}^{(n+1)}\right\|_{L^{2}(\mathbf{R})}^{2}\right) \leq \left\|u_{0}\right\|_{L^{2}(\mathbf{R})}^{2} + \operatorname{Lip}_{\sigma}^{2} \cdot \int_{0}^{t} \mathbb{E}\left(\left\|u_{s}^{(n)}\right\|_{L^{2}(\mathbf{R})}^{2}\right) \cdot \left\|p_{t-s}\right\|_{L^{2}(\mathbf{R})}^{2} \,\mathrm{d}s.$$

$$(2.11)$$

Therefore, if we set

$$M^{(k)}(\lambda) := \sup_{t \ge 0} \left[e^{-\lambda t} \mathbf{E} \left(\left\| u_t^{(k)} \right\|_{L^2(\mathbf{R})}^2 \right) \right],$$
(2.12)

then it follows that

$$M^{(n+1)}(\lambda) \leq \|u_0\|_{L^2(\mathbf{R})}^2 + \operatorname{Lip}_{\sigma}^2 \cdot M^{(n)}(\lambda) \cdot \int_0^\infty e^{-\lambda(t-s)} \|p_{t-s}\|_{L^2(\mathbf{R})}^2 \, \mathrm{d}s$$

= $\|u_0\|_{L^2(\mathbf{R})}^2 + \frac{\operatorname{Lip}_{\sigma}^2}{2\sqrt{2\kappa\lambda}} M^{(n)}(\lambda).$ (2.13)

Thus, in particular, $\sup_{n\geq 0} M^{(n)}(\lambda) < \infty$ if $\lambda > \operatorname{Lip}_{\sigma}^{4}/(8\kappa)$. We can argue similarly to show also that if $\lambda > \operatorname{Lip}_{\sigma}^{4}/(8\kappa)$, then

$$\sum_{n} \sup_{t \ge 0} \left[e^{-\lambda t} \mathbf{E} \left(\left\| u_t^{(n+1)} - u_t^{(n)} \right\|_{L^2(\mathbf{R})}^2 \right) \right]^{1/2} < \infty.$$
(2.14)

In particular, uniqueness shows that if $\lambda > \text{Lip}_{\sigma}^4/(8\kappa)$, then

$$\lim_{n \to \infty} \sup_{t \ge 0} \left[e^{-\lambda t} \mathbf{E} \left(\left\| u_t^{(n)} - u_t \right\|_{L^2(\mathbf{R})}^2 \right) \right] = 0.$$
(2.15)

Consequently, if $\lambda > \text{Lip}_{\sigma}^4/(8\kappa)$, then

$$\sup_{t\geq 0} \left[e^{-\lambda t} \mathbf{E} \left(\|u_t\|_{L^2(\mathbf{R})}^2 \right) \right] = \lim_{n \to \infty} M^{(n)}(\lambda) \le \sup_{k\geq 0} M^{(k)}(\lambda) < \infty.$$
(2.16)

The second inequality of (2.1) follows readily from this bound.

$$\square$$

3. Proof of Theorem 1.1

Our proof of Theorem 1.1 hinges on a number of steps, which we develop separately. First we recall the following.

Proposition 3.1 (Theorem 2.1 and Example 2.9 of [8]). *If* u_0 *is bounded and measurable, then* $u_t(x) \in L^p(\mathbb{P})$ *for all* $p \in [1, \infty)$ *. Moreover,* $\overline{\gamma}(p) < \infty$ *for all* $p \in [1, \infty)$ *and* $\overline{\gamma}(2) \leq \operatorname{Lip}_{\sigma}^4/(8\kappa)$ *, where*

$$\overline{\gamma}(p) := \limsup_{t \to \infty} t^{-1} \sup_{x \in \mathbf{R}} \ln \mathbb{E}(|u_t(x)|^p) < \infty.$$
(3.1)

[Note that, in the preceding, the supremum is outside the expectation.] Next, we record a simple though crucial property of the function $\overline{\gamma}$.

Remark 3.2. Suppose X is a nonnegative random variable with finite moments of all orders. By Hölder's inequality, $p \mapsto \ln E(X^p)$ is convex on $[1, \infty)$. It follows that $\overline{\gamma}$ is convex – in particular continuous – on $[1, \infty)$.

Now we begin our analysis, in earnest, by deriving an upper bound on the $L^{k}(P)$ -norm of the solution $u_{t}(x)$ that includes simultaneously a sharp decay rate in x and a sharp explosion rate in t.

Lemma 3.3. Suppose that $u_0 \neq 0$, and u_0 is supported in [-K, K] for some finite constant K > 0. Then, for all real numbers $k \in [1, \infty)$ and $p \in (1, \infty)$,

$$\limsup_{t \to \infty} t^{-1} \sup_{x \in \mathbf{R}} \left(\frac{x^2}{4t^2} + \frac{k + 1 - (1/p)}{k} \ln \mathbb{E}\left(\left| u_t(x) \right|^k \right) \right) \le \frac{\overline{\gamma}(kp)}{p}.$$
(3.2)

Proof. According to Mueller's comparison principle ([14]; more specifically, see [5], Theorem 5.1, p. 130; see also [7,16]), the solution to (1.1) has the following nonnegativity property: Because $u_0 \ge 0$ then outside a single null set, $u_t \ge 0$ for all $t \ge 0$. Since $u_t(x) \in L^2(P)$ [e.g., by Proposition 3.1], the stochastic integral in (1.5) is a martingale-measure stochastic integral in $L^2(P)$ [say], and consequently has mean zero. And therefore,

$$\left\|u_{t}(x)\right\|_{1} = (p_{t} * u_{0})(x) = \frac{1}{\sqrt{4\kappa\pi t}} \int_{-K}^{K} e^{-(x-y)^{2}/(4\kappa t)} u_{0}(y) \,\mathrm{d}y.$$
(3.3)

Because $(x - y)^2 \ge (x^2/2) - K^2$,

$$\left\| u_t(x) \right\|_1 \le \operatorname{const} \cdot e^{-x^2/(4t)} \quad \text{for all } x \in \mathbf{R} \text{ and } t \ge 1.$$
(3.4)

The constant appearing in the above display depends on K. Next we note that for every $\theta \in (0, \infty)$,

$$\mathbb{E}\left(\left|u_{t}(x)\right|^{k}\right) \leq \theta^{k} + \mathbb{E}\left(\left|u_{t}(x)\right|^{k}; u_{t}(x) \geq \theta\right) \\ \leq \theta^{k} + \left(\mathbb{E}\left(\left|u_{t}(x)\right|^{kp}\right)\right)^{1/p} \cdot \left(\mathbb{P}\left\{u_{t}(x) > \theta\right\}\right)^{1-(1/p)}.$$
(3.5)

Proposition 3.1 implies that

$$\sup_{x \in \mathbf{R}} \left(\mathbb{E}\left(\left| u_t(x) \right|^{kp} \right) \right)^{1/p} \le \exp\left(t \; \frac{\overline{\gamma}(kp) + \mathrm{o}(1)}{p} \right), \tag{3.6}$$

where $o(1) \rightarrow 0$ as $t \rightarrow \infty$. Also, we can apply (3.4) together with the Chebyshev inequality to find that

$$\left(\mathsf{P}\left\{u_t(x) > \theta\right\}\right)^{1-(1/p)} \le \operatorname{const} \cdot \theta^{-1+(1/p)} \exp\left(-\frac{x^2}{4t} \cdot \left[1 - \frac{1}{p}\right]\right).$$
(3.7)

In light of (3.6) and (3.7), we can deduce that the following from (3.5):

$$\mathbf{E}\left(\left|u_{t}(x)\right|^{k}\right) \leq \inf_{\theta > 0}\left(\theta^{k} + \alpha \theta^{-1 + (1/p)}\right),\tag{3.8}$$

where

$$\alpha := \exp\left(-\frac{x^2}{4t} \cdot \left[1 - \frac{1}{p}\right] + t \, \frac{\overline{\gamma}(kp) + \mathrm{o}(1)}{p}\right). \tag{3.9}$$

Some calculus shows that the function $g(\theta) := (\theta^k + \alpha \theta^{-1+1/p}) \mathbf{1}_{(0,\infty)}(\theta)$ attains its minimum at $\theta := ((p-1)/(kp)^{p/(kp+p-1)})$. Consequently,

$$\mathbb{E}(|u_t(x)|^k) \le \alpha^{kp/(kp+p-1)} \left(\frac{p-1}{kp}\right)^{kp/(kp+p-1)} \cdot \left(\frac{1-p-kp}{1-p}\right).$$
(3.10)

We now divide both sides of the above display by $\alpha^{kp/(kp+p-1)}$ and take the appropriate limit to obtain the result. \Box

Our next lemma is a basic estimate of continuity in the variable x. It is not entirely standard as it holds uniformly for all times $t \ge 0$. We emphasize that the constant p is assumed to be an integer. We will deal with this shortcoming subsequently.

Lemma 3.4. Suppose that the initial function u_0 is Hölder continuous of order $\ge 1/2$. Then, for all integers $p \ge 1$ and $\beta > \overline{\gamma}(2p)$ there exists a constant $A_{p,\beta} \in (0,\infty)$ such that the following holds: Simultaneously for all $t \ge 0$,

$$\sup_{j \in \mathbf{Z}} \sup_{j \le x < x' \le j+1} \left\| \frac{u_t(x) - u_t(x')}{|x - x'|^{1/2}} \right\|_{2p} \le A_{p,\beta} e^{\beta t/(2p)}.$$
(3.11)

Proof. Burkholder's inequality [2] and Minkowski's inequality together imply that

$$\begin{aligned} \left\| u_{t}(x) - u_{t}(x') \right\|_{2p} &\leq \left| (p_{t} * u_{0})(x) - (p_{t} * u_{0})(x') \right| \\ &+ z_{2p} \left\| \int_{0}^{t} \mathrm{d}s \int_{-\infty}^{\infty} \mathrm{d}y \left| \sigma \left(u_{s}(y) \right) \right|^{2} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2} \right\|_{p}^{1/2} \\ &\leq \left| (p_{t} * u_{0})(x) - (p_{t} * u_{0})(x') \right| \\ &+ z_{2p}' \left\| \int_{0}^{t} \mathrm{d}s \int_{-\infty}^{\infty} \mathrm{d}y \left| u_{s}(y) \right|^{2} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2} \right\|_{p}^{1/2}, \end{aligned}$$
(3.12)

where z_p is a positive and finite constant that depend only on p, and $z'_p := z_p \operatorname{Lip}_{\sigma}$. On one hand,

$$\sup_{t \ge 0} \sup_{|x-x'| \le \delta} \left| (p_t * u_0)(x) - (p_t * u_0)(x') \right| \le \sup_{|a-b| \le \delta} \left| u_0(a) - u_0(b) \right| \le \operatorname{const} \cdot \delta^{1/2}.$$
(3.13)

On the other hand, the generalized Hölder inequality suggests that if $p \ge 1$ is an integer, then for all $s_1, \ldots, s_p \ge 0$ and $y_1, \ldots, y_p \in \mathbf{R}$,

$$\mathbb{E}\left(\prod_{j=1}^{p} |u_{s_j}(y_j)|^2\right) \le \prod_{j=1}^{p} ||u_{s_j}(y_j)||_{2p}^2.$$
(3.14)

[It might help to recall that the generalized Hölder inequality states that $E(\zeta_1 \cdots \zeta_p) \leq \prod_{j=1}^p \|\zeta_j\|_p$ for all nonnegative random variables ζ_1, \ldots, ζ_p .] Therefore,

$$\left\| \int_{0}^{t} ds \int_{-\infty}^{\infty} dy \left| u_{s}(y) \right|^{2} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2} \right\|_{p}$$

$$\leq \int_{0}^{t} ds \int_{-\infty}^{\infty} dy \left\| u_{s}(y) \right\|_{2p}^{2} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2}.$$
(3.15)

[Write the *p*th power of the left-hand side as the expectation of a product and apply (3.14).]

A proof by contradiction shows that Proposition 3.1 gives the following [see [8] for more details]:

$$c_{\beta} := \sup_{s \ge 0} \sup_{y \in \mathbf{R}} \left[e^{-\beta s} \mathbf{E} \left(\left| u_s(y) \right|^{2p} \right) \right] < \infty \quad \text{for all } \beta > \overline{\gamma}(2p).$$
(3.16)

We omit the details, but state instead that the argument is quite similar to the real-variable method that was employed earlier, in the paragraph that preceeds (2.9).

Consequently,

$$\left\| \int_{0}^{t} ds \int_{-\infty}^{\infty} dy \left| u_{s}(y) \right|^{2} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2} \right\|_{p}$$

$$\leq c_{\beta}^{1/p} \cdot \int_{0}^{t} ds \int_{-\infty}^{\infty} dy e^{\beta s/p} \cdot \left| p_{t-s}(y-x) - p_{t-s}(y-x') \right|^{2}$$

$$\leq c_{\beta}^{1/p} e^{\beta t/p} \cdot \int_{0}^{\infty} ds e^{-\beta s/p} \int_{-\infty}^{\infty} dy \left| p_{s}(y-x) - p_{s}(y-x') \right|^{2}.$$
(3.17)

Since $\hat{p}_s(\xi) = \exp(-\kappa s\xi^2)$, Plancherel's theorem tells us that the right-hand side of the preceding inequality is equal to

$$\frac{c_{\beta}^{1/p} e^{\beta t/p}}{\pi} \cdot \int_{0}^{\infty} ds \, e^{-\beta s/p} \int_{-\infty}^{\infty} d\xi \, e^{-2\kappa s \xi^{2}} [1 - \cos(\xi (x - x'))]$$
$$= \frac{2c_{\beta}^{1/p} e^{\beta t/p}}{\pi} \cdot \int_{0}^{\infty} \frac{[1 - \cos(\xi (x - x'))]}{(\beta/p) + 2\kappa \xi^{2}} \, d\xi.$$
(3.18)

Because $1 - \cos \theta \le \min(1, \theta^2)$, a direct estimation of the integral leads to the following bound:

$$\left\|\int_{0}^{t} ds \int_{-\infty}^{\infty} dy \left|u_{s}(y)\right|^{2} \cdot \left|p_{t-s}(y-x) - p_{t-s}(y-x')\right|^{2}\right\|_{p}$$

$$\leq \operatorname{const} \cdot e^{\beta t/p} \cdot |x-x'|, \qquad (3.19)$$

where the implied constant depends only on p, κ , and β . This, (3.13), and (3.12) together imply the lemma.

The preceding lemma holds for all *integers* $p \ge 1$. In the following, we improve it [at a slight cost] to the case that $p \in (1, 2)$ is a real number.

Lemma 3.5. Suppose the conditions of Lemma 3.4 are met. Then for all $p \in (1, 2)$ and $\delta \in (0, 1)$ there exists a constant $B_{p,\delta} \in (0, \infty)$ such that the following holds: Simultaneously for all $t \ge 0$ and $x, x' \in \mathbf{R}$ with $|x - x'| \le 1$,

$$\mathbf{E}\left(\left|u_{t}(x)-u_{t}(x')\right|^{2p}\right) \leq B_{p,\delta} \cdot \left|x-x'\right|^{p} \cdot \mathrm{e}^{(1+\delta)\lambda_{p}t},\tag{3.20}$$

where

$$\lambda_p := (2-p)\overline{\gamma}(2) + (p-1)\overline{\gamma}(4). \tag{3.21}$$

Proof. We start by writing $E(|u_t(x) - u_t(x')|^{2p})$ as

$$\mathbf{E}(|u_t(x) - u_t(x')|^{2(2-p)}|u_t(x) - u_t(x')|^{4(p-1)}).$$
(3.22)

We can apply Hölder's inequality to conclude that for all $p \in (1, 2), t \ge 0$, and $x, x' \in \mathbf{R}$,

$$\mathbb{E}(|u_t(x) - u_t(x')|^{2p}) \le [\mathbb{E}(|u_t(x) - u_t(x')|^2)]^{2-p} [\mathbb{E}(|u_t(x) - u_t(x')|^4)]^{p-1}.$$
(3.23)

We now use Lemma 3.4 to obtain the following:

$$\left[\mathsf{E}(\left|u_{t}(x) - u_{t}(x')\right|^{2}) \right]^{2-p} \leq |x - x'|^{(2-p)} A_{1,\beta_{1}}^{2(2-p)} \mathrm{e}^{\beta_{1}(2-p)t}$$
(3.24)

and

$$\left[\mathbb{E} \left(\left| u_t(x) - u_t(x') \right|^4 \right) \right]^{p-1} \le \left| x - x' \right|^{2(p-1)} A_{2,\beta_2}^{4(p-1)} e^{\beta_2(p-1)t}, \tag{3.25}$$

where $A_{1,\beta_1}, A_{2,\beta_2} \in (0,\infty)$ and $\beta_1 > \bar{\gamma}(2)$ and $\beta_2 > \bar{\gamma}(4)$ are fixed and finite constants. The proof now follows by combining the above and choosing β_1 and β_2 such that $(1 + \delta)\bar{\gamma}(2) > \beta_1 > \bar{\gamma}(2)$ and $(1 + \delta)\bar{\gamma}(4) > \beta_2 > \bar{\gamma}(4)$. \Box

The preceding lemma allows for a uniform modulus of continuity estimate, which we record next.

Lemma 3.6. Suppose the conditions of Lemma 3.4 are met. Then for all $p \in (1, 2)$ and $\varepsilon, \delta \in (0, 1)$ there exists $C_{p,\varepsilon,\delta} \in (0, \infty)$ such that simultaneously for all $t \ge 0$,

$$\sup_{j \in \mathbf{Z}} \left\| \sup_{j \le x < x' \le j+1} \frac{|u_t(x) - u_t(x')|^2}{|x - x'|^{1 - \varepsilon}} \right\|_p \le C_{p,\varepsilon,\delta} \cdot e^{(1+\delta)\lambda_p t},$$
(3.26)

where λ_p was defined in (3.21).

Proof. The proof consists of an application of the Kolmogorov continuity theorem. Recall that the spatial dimension is 1 and we are choosing a continuous version of the solution $(t, x) \mapsto u_t(x)$. Since p > 1 in Lemma 3.5, we can use a suitable version of Kolmogorov continuity theorem, for example Theorem 4.3 of reference [5], p. 10, to obtain the result. The stated dependence of the constant, $C_{p,\varepsilon,\delta}$ is consequence of the explicit form of inequality (3.20) and the proof of Theorem 4.3 in [5].

Before we begin our proof of Theorem 1.1, we prove that under some condition the $L^2(P)$ -norm of the solution has an effectively-compact support.

Proposition 3.7. If the conditions of Theorem 1.1 are met, then there exists a finite and positive constant m such that $u_t(x)$ has an effectively-compact support with radius of effective support p(t) = mt.

Proof. We begin by noting that for all m, t > 0,

$$\int_{|x|>mt} |u_t(x)|^2 \,\mathrm{d}x \le \int_{|x|>mt} u_t(x) \,\mathrm{d}x + \int_{\substack{|x|>mt\\u_t(x)\ge 1}} |u_t(x)|^2 \,\mathrm{d}x.$$
(3.27)

Therefore,

$$\mathbb{E}\left(\int_{|x|>mt} |u_t(x)|^2 \,\mathrm{d}x\right) \le \int_{|x|>mt} (p_t * u_0)(x) \,\mathrm{d}x + \int_{|x|>mt} \mathbb{E}\left(|u_t(x)|^2; u_t(x) \ge 1\right) \,\mathrm{d}x.$$
(3.28)

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Since u_0 has compact support, (3.4) implies that

$$\int_{|x|>mt} (p_t * u_0)(x) \, \mathrm{d}x = O(\mathrm{e}^{-m^2 t/2}) \quad \text{as } t \to \infty.$$
(3.29)

Next we estimate the final integral in (3.28).

Thanks to (3.4) and Chebyshev's inequality,

$$P\{u_t(x) \ge 1\} \le \text{const} \cdot e^{-x^2/(4t)},$$
(3.30)

uniformly for all $x \in \mathbf{R}$ and $t \ge 1$. Also, from Proposition 3.1, there exists a constant $b \in (0, \infty)$ such that

$$\sup_{x \in \mathbf{R}} \mathbb{E}(|u_t(x)|^4) \le b e^{bt/4} \quad \text{for all } t \ge 1.$$
(3.31)

Using the preceding two inequalities, the right-hand side of inequality (3.28) reduces to

$$E\left(\int_{|x|>mt} |u_t(x)|^2 dx\right)$$

$$\leq O(e^{-m^2t/2}) + \text{const} \cdot \int_{|x|>mt} \sqrt{E(|u_t(x)|^4)} e^{-x^2/(8t)} dx$$

$$\leq O(e^{-m^2t/2}) + \text{const} \cdot b^{1/2} e^{bt/8} \cdot \int_{|x|>mt} e^{-x^2/(8t)} dx.$$
(3.32)

We now choose and fix $m > \sqrt{b}$ to obtain from the preceding that

$$\limsup_{t \to \infty} t^{-1} \ln \mathbb{E}\left(\int_{|x| > mt} |u_t(x)|^2 \, \mathrm{d}x\right) < 0.$$
(3.33)

This implies part (b) of Definition 1.2 with p(t) = mt. We now prove the remaining part of Definition 1.2. From Theorem 2.1 and the preceding, we obtain for infinitely-many values of $t \to \infty$:

$$\exp\left(\left[\frac{\mathrm{L}_{\sigma}^{4}}{8\kappa} + \mathrm{o}(1)\right]t\right) \leq \mathrm{E}\left(\int_{-\infty}^{\infty} |u_{t}(x)|^{2} \,\mathrm{d}x\right)$$
$$= \mathrm{E}\left(\int_{-mt}^{mt} |u_{t}(x)|^{2} \,\mathrm{d}x\right) + \mathrm{o}(1).$$
(3.34)

This finishes the proof.

We will need the following elementary real-variable lemma from the theory of slowly-varying functions. It is without doubt well known; we include a derivation for the sake of completeness only.

Lemma 3.8. For every $q, \eta \in (0, \infty)$,

$$\int_{e}^{\infty} \exp\left(-\frac{q(\ln x)^{\eta+1}}{t}\right) dx = O\left(t^{1/\eta} \exp\left\{(t/q)^{1/\eta}\right\}\right)$$

$$as \ t \to \infty.$$
(3.35)

Proof. The proof uses some standard tricks. First we write the integral as

$$\int_{e}^{\infty} e^{-q(\ln x)^{\eta+1}/t} \, \mathrm{d}x = \int_{1}^{\infty} e^{-qz^{\eta+1}/t} e^{z} \, \mathrm{d}z.$$
(3.36)

 \square

Next we change variables $[w := z/\theta]$, for an arbitrary $\theta > 0$, and find that

$$\int_{e}^{\infty} e^{-q(\ln x)^{\eta+1}/t} dx = \theta \int_{1/\theta}^{\infty} \exp\left(-\frac{q\theta^{\eta+1}}{t}w^{\eta+1} + \theta w\right) dw.$$
(3.37)

Upon choosing $\theta := (t/q)^{1/\eta}$, we obtain

$$-\frac{q\theta^{\eta+1}}{t}w^{\eta+1} + \theta w = \left(\frac{t}{q}\right)^{1/\eta} (w - w^{\eta+1}),$$

and this yields

$$\int_{e}^{\infty} e^{-q(\ln x)^{\eta+1}/t} \, \mathrm{d}x = (t/q)^{1/\eta} \int_{(q/t)^{1/\eta}}^{\infty} e^{(t/q)^{1/\eta} \cdot (w-w^{\eta+1})} \, \mathrm{d}w.$$
(3.38)

Therefore, for *t* sufficiently large, we split the integral on the right-hand side of the previous display as follows:

$$\int_{e}^{\infty} e^{-q(\ln x)^{\eta+1}/t} \, \mathrm{d}x = (t/q)^{1/\eta} (I_1 + I_2), \tag{3.39}$$

where

$$I_{1} := \int_{(q/t)^{1/\eta}}^{1} \exp((t/q)^{1/\eta} \cdot (w - w^{\eta + 1})) dw,$$

$$I_{2} := \int_{1}^{\infty} \exp(-(t/q)^{1/\eta} \cdot w(w^{\eta} - 1)) dw.$$
(3.40)

Clearly,

$$I_2 \le 1 + \int_2^\infty \exp\left(-\left(2^\eta - 1\right)(t/q)^{1/\eta} \cdot w\right) \mathrm{d}w = \mathcal{O}(1).$$
(3.41)

The lemma follows because the integrand of I_1 is at most $\exp((t/q)^{1/\eta})$.

We are now ready to establish Theorem 1.1.

Proof of Theorem 1.1. The proof of the first inequality of the theorem is a continuation of the proof Proposition 3.7. Indeed, from (3.34), we obtain

$$\exp\left(\left[\frac{\mathbf{L}_{\sigma}^{4}}{8\kappa} + \mathbf{o}(1)\right]t\right) \leq E\left(\int_{-\infty}^{\infty} |u_{t}(x)|^{2} dx\right)$$
$$\leq E\left(\int_{-mt}^{mt} |u_{t}(x)|^{2} dx\right) + \mathbf{o}(1)$$
$$\leq 2mt \cdot \sup_{x \in \mathbf{R}} E\left(|u_{t}(x)|^{2}\right) + \mathbf{o}(1), \tag{3.42}$$

valid for $t \to \infty$. We obtain first inequality of the theorem after taking the appropriate limit.

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Next we prove the second inequality of the theorem by first observing that for every $j \ge 1$, every increasing sequence of real numbers $\{a_j\}_{j=1}^{\infty}$ with $\sup_{j\ge 1}(a_{j+1}-a_j) \le 1$, $p \in (1, 2)$, $\varepsilon \in (0, 1)$, and $t \ge 0$,

$$\sup_{a_j \le x \le a_{j+1}} |u_t(x)|^{2p} = \sup_{a_j \le x \le a_{j+1}} |u_t(a_j) + u_t(x) - u_t(a_j)|^{2p}$$

$$\le 2^{2p-1} \Big(|u_t(a_j)|^{2p} + \sup_{a_j \le x \le a_{j+1}} |u_t(x) - u_t(a_j)|^{2p} \Big)$$

$$\le 2^{2p-1} \Big(|u_t(a_j)|^{2p} + (a_{j+1} - a_j)^{p(1-\varepsilon)} \Omega_j^p \Big),$$
(3.43)

where

$$\Omega_j := \sup_{a_j \le x < x' \le a_{j+1}} \frac{|u_t(x) - u_t(x')|^2}{|x - x'|^{1 - \varepsilon}}.$$
(3.44)

Consequently,

$$\mathbb{E}\Big(\sup_{a_j \le x \le a_{j+1}} |u_t(x)|^{2p}\Big) \le 2^{2p-1} \Big(\mathbb{E}\big(|u_t(a_j)|^{2p}\big) + (a_{j+1} - a_j)^{p(1-\varepsilon)} \mathbb{E}\big(\Omega_j^p\big)\big).$$
(3.45)

We use inequality (3.2) of Lemma 3.3 with k := 2p and $x := a_i$ to find that

$$\mathbb{E}\left(\left|u_{t}(a_{j})\right|^{2p}\right) \leq \operatorname{const} \cdot \exp\left(\beta_{p} \cdot \left[t\frac{\overline{\gamma}(2p^{2}) + \mathrm{o}(1)}{p} - \frac{a_{j}^{2}}{4t^{2}}\right]\right),\tag{3.46}$$

where

$$\beta_p := \frac{p}{p+1 - (1/p)},\tag{3.47}$$

the implied constant does not depend on j or t, and $o(1) \rightarrow 0$ as $t \rightarrow \infty$, uniformly for all j. Also, Lemma 3.6 implies that

$$\sup_{j\geq 1} \mathbb{E}\left(\Omega_{j}^{p}\right) \leq C_{p,\varepsilon,\delta} \cdot e^{p(1+\delta)\lambda_{p}t},$$
(3.48)

where δ is an arbitrarily-small positive constant, which we will choose and fix appropriately later on. We can combine the preceding inequalities to deduce that

$$\mathbb{E}\Big(\sup_{a_j \le x \le a_{j+1}} |u_t(x)|^{2p}\Big) \le \operatorname{const} \cdot e^{-\beta_p a_j^2/(4t^2)} \cdot e^{\beta_p t(\overline{\gamma}(2p^2) + o(1))/p} + \operatorname{const} \cdot (a_{j+1} - a_j)^{p(1-\varepsilon)} e^{p(1+\delta)\lambda_p t}.$$
 (3.49)

Choose and fix an integer $\nu \ge 1$. We apply the preceding with $p(1 - \varepsilon) > 1$; we also choose the a_l 's so that $a_1 := 0$, $0 \le a_{j+1} - a_j \le 1$ for all $j \ge 1$, and $a_j := (\log j)^{\nu}$ for all j sufficiently large. Because $a_{j+1} - a_j = O((\ln j)^{\nu}/j)$ as $j \to \infty$,

$$\sum_{j=1}^{\infty} (a_{j+1} - a_j)^{p(1-\varepsilon)} < \infty.$$
(3.50)

Also, for all J > 1 + e, sufficiently large,

$$\sum_{j=J}^{\infty} e^{-\beta_p a_j^2/(4t)} \le \int_{J-1}^{\infty} e^{-\beta_p (\ln x)^{2\nu}/(4t^2)} dx$$
$$= O(t^{1/(2\nu-1)} e^{(4t^2/\beta_p)^{1/(2\nu-1)}}) \quad (t \to \infty),$$
(3.51)

where we have used Lemma 3.8 for the last equality. We can choose $\nu := \frac{1}{2}(\delta^{-1} + 1)$ so that $1/(2\nu - 1) = \delta$. We can combine these terms to deduce the following:

$$E\left(\sup_{x \ge a_J} |u_t(x)|^{2p}\right) \le \sum_{j=J}^{\infty} E\left(\sup_{a_j \le x \le a_{j+1}} |u_t(x)|^{2p}\right) = O\left(t^{\delta} e^{(4t^2/\beta_p)^{\delta} + \beta_p t(\overline{\gamma}(2p^2) + o(1))/p} + e^{p(1+\delta)\lambda_p t}\right).$$
(3.52)

A similar – though slightly simpler – argument can be used to derive the very same upper bound for the quantity $E(\sup_{0 \le x \le a_I} |u_t(x)|^{2p})$. We now use symmetry and let $\delta \downarrow 0$,

$$\limsup_{t \to \infty} t^{-1} \ln \mathbb{E}\left(\sup_{x \in \mathbf{R}} \left| u_t(x) \right|^{2p}\right) \le \max\left\{\frac{\beta_p \overline{\gamma}(2p^2)}{p}, p\lambda_p\right\}.$$
(3.53)

Let us substitute the evaluation of β_p in terms of p to find that

$$\limsup_{t \to \infty} t^{-1} \ln \mathbb{E} \left(\sup_{x \in \mathbf{R}} \left| u_t(x) \right|^{2p} \right) \le \max \left\{ \frac{\overline{\gamma}(2p^2)}{p+1-(1/p)}, p\lambda_p \right\}.$$
(3.54)

This and Jensen's inequality together prove that

$$\limsup_{t \to \infty} t^{-1} \ln \mathbb{E}\left(\sup_{x \in \mathbf{R}} \left| u_t(x) \right|^2\right) \le \frac{1}{p} \max\left\{\frac{\overline{\gamma}(2p^2)}{p+1-(1/p)}, p\lambda_p\right\},\tag{3.55}$$

and this valid for all $p \in (1, 2)$. As $p \downarrow 1$, $\lambda_p \to \overline{\gamma}(2)$. Moreover, $\overline{\gamma}(2p^2) \to \overline{\gamma}(2)$ because $\overline{\gamma}$ is convex and hence continuous on $[1, \infty)$ [Remark 3.2]. It follows that

$$\limsup_{t \to \infty} t^{-1} \ln \mathbb{E} \left(\sup_{x \in \mathbf{R}} \left| u_t(x) \right|^2 \right) \le \overline{\gamma}(2), \tag{3.56}$$

and this is $\leq \text{Lip}_{\sigma}^4/(8\kappa)$ by Proposition 3.1. The latter proposition implies the theorem because $\text{E}(\sup_x |u_t(x)|^2) \geq \sup_x \text{E}(|u_t(x)|^2)$.

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