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# DIRECT IMAGES OF SEMI-MEROMORPHIC CURRENTS 

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Abstract. - We introduce a calculus for the class $A S M(X)$ of direct images of semi-meromorphic currents on a reduded analytic space $X$, that extends the classical calculus due to Coleff, Herrera and Passare. Our main result is that each element in this class acts as a kind of multiplication on the sheaf $\mathcal{P} \mathcal{M}_{X}$ of pseudomeromorphic currents on $X$. We also prove that $A S M(X)$ as well as $\mathcal{P} \mathcal{M}_{X}$ and certain subsheaves are closed under the action of holomorphic differential operators and interior multiplication by holomorphic vector fields.

Résumé. - Nous introduisons un calcul pour la classe $A S M(X)$ d'images directes de courants semi-méromorphes sur un espace analytique reduit $X$, qui étend le calcul classique de Coleff, Herrera et Passare. Notre résultat principal montre que chaque élément de cette classe agit de manière analogue à une multiplication sur le faisceau $\mathcal{P} \mathcal{M}_{X}$ de courants pseudoméromorphes sur $X$. Nous prouvons également que $A S M(X)$ ainsi que $\mathcal{P} \mathcal{M}_{X}$ et certains sous-faisceaux sont fermés sous l'action des opérateurs différentiels holomorphes et la multiplication intérieure par des champs vectoriels holomorphes.

## 1. Introduction

Let $f$ be a generically nonvanishing holomorphic function on a reduced analytic space $X$ of pure dimension $n$. It was proved by Herrera and Lieberman, [14], that one can define the principal value current

$$
\begin{equation*}
\left[\frac{1}{f}\right] \cdot \xi:=\lim _{\epsilon \rightarrow 0} \int_{|f|^{2}>\epsilon} \frac{\xi}{f}, \tag{1.1}
\end{equation*}
$$

for test forms $\xi$. It follows that $\bar{\partial}[1 / f]$ is a current with support on the zero set $Z(f)$ of $f$; such a current is called a residue current. Coleff and Herrera,

[^0][13], introduced products of principal value and residue currents, like
\[

$$
\begin{equation*}
\left[1 / f_{1}\right] \ldots\left[1 / f_{r}\right] \bar{\partial}\left[1 / f_{r+1}\right] \wedge \ldots \wedge \bar{\partial}\left[1 / f_{m}\right] \tag{1.2}
\end{equation*}
$$

\]

The product of principal value currents is commutative, but when there are residue factors, like $\bar{\partial}\left[1 / f_{j}\right]$, present these products are not (anti-)commutative in general.

In the literature there are various generalizations and related currents, for instance the abstract so-called Coleff-Herrera currents introduced by Björk, see [12], the Bochner-Martinelli type residue currents introduced in [21], and generalizations in, e.g., $[3,5,9]$.

In order to obtain a coherent approach to questions about residue and principal value currents the sheaf $\mathcal{P} \mathcal{M}_{X}$ of pseudomeromorphic currents on $X$ was introduced in [10], and further developed in [7]; this sheaf consists of direct images under holomorphic mappings of products of test forms and currents like (1.2). See Section 2 below for the precise definition. This sheaf is closed under $\bar{\partial}$ and under multiplication by smooth forms. Pseudomeromorphic currents have a geometric nature, similar to positive closed (or normal) currents. For example, the dimension principle states that if the pseudomeromorphic current $\mu$ has bidegree ( $*, p$ ) and support on a variety of codimension larger than $p$, then $\mu$ must vanish. Moreover one can form restrictions $\mathbb{1}_{W} \mu$ of the pseudomeromorphic current $\mu$ to analytic (or constructible) subsets $W \subset X$, such that

$$
\begin{equation*}
\mathbb{1}_{V} \mathbb{1}_{W} \mu=\mathbb{1}_{V \cap W} \mu \tag{1.3}
\end{equation*}
$$

see Section 2.2. The notion of pseudomeromorphic currents plays a decisive role in, for instance, $[7,8,10,11,15,16,18,22,23,24,25]$.

It is well-known that one cannot multiply currents in general. Several attempts to find a working calculus for principal value and residue currents have been made. A famous result by Coleff and Herrera, [13], see also Passare, [20], asserts that (1.2) has all expected (anti-)commutativity properties as long as the common zero set of $f_{1}, \ldots, f_{m}$ has codimension $m$. Various extension are introduced in the references above. In [10] we proved that one can give a reasonable meaning to a product $[1 / f] \mu$ for any holomorphic function $f$ and pseudomeromorphic current $\mu$; more precisely one should consider this as an operator

$$
\begin{equation*}
\mu \mapsto[1 / f] \mu \tag{1.4}
\end{equation*}
$$

on the sheaf $\mathcal{P} \mathcal{M}_{X}$.
We have not found a way to define a reasonable product of general pseudomeromorphic currents. Our first objective in this paper is to study a
generalization of principal value currents leading to an extension of (1.4). Following [7] we say that a current $a$ is almost semi-meromorphic, $a \in$ $\operatorname{ASM}(X)$, if it is the direct image under a modification of a semi-meromorphic current, i.e., a current of the form $\omega[1 / f]$, where $f$ is a holomorphic section of a line bundle and $\omega$ is a smooth form with values in the same bundle. Almost semi-meromorphic currents are pseudomeromorphic and in many ways they generalize principal value currents. For example, it turns out that they form an (anti-)commutative algebra, see Section 4. Moreover $\operatorname{ASM}(X)$ is closed under $\partial$, see Proposition 4.16. Taking $\bar{\partial}$ of $a \in \operatorname{ASM}(X)$, however, yields an almost semi-meromorphic current plus a residue current supported on the Zariski singular support, $\operatorname{ZSS}(a)$, of $a$, which is the smallest analytic set where $a$ is not smooth. Many of the currents in the references above can be considered as (products of) the residues of almost semi-meromorphic currents. Theorem 4.8 states that the mapping (1.4) holds for any almost semi-meromorphic current $a$ instead of $[1 / f]$. More precisely, there is a unique extension to $X$ of the current $a \wedge \mu$, defined in the obvious way in $X \backslash Z S S(a)$, such that its restriction to $Z S S(a)$ is zero.

A second objective is to prove that $\mathcal{P} \mathcal{M}_{X}$ and $\operatorname{ASM}(X)$ are closed under interior multiplication by a holomorphic vector field $\xi$ and under the Lie derivative with respect to $\xi$; see Sections 3 and 4.5 .

In Section 2 we recall basic known properties of the sheaf $\mathcal{P} \mathcal{M}_{X}$ and provide some new results, e.g., Theorem 2.15 gives a new quite natural characterization of pseudomeromorphicity. Section 4 is devoted to the study of $\operatorname{ASM}(X)$.

Ackowledgment. We are grateful to the referee for carefully reading and pointing out unclarities and misprints.

## 2. Pseudomomeromorphic currents

In one complex variable $s$ one can define the principal value current $\left[1 / s^{m}\right]$ for instance as the value

$$
\left[\frac{1}{s^{m}}\right]=\left.\frac{|s|^{2 \lambda}}{s^{m}}\right|_{\lambda=0}
$$

of the current-valued analytic continuation of $\lambda \mapsto|s|^{2 \lambda} / s^{m}$, a priori defined for $\operatorname{Re} \lambda \gg 0$, see, e.g., [3, Lemma 2.1]. We have the relations

$$
\begin{equation*}
\frac{\partial}{\partial s}\left[\frac{1}{s^{m}}\right]=-m\left[\frac{1}{s^{m+1}}\right], \quad s\left[\frac{1}{s^{m+1}}\right]=\left[\frac{1}{s^{m}}\right] . \tag{2.1}
\end{equation*}
$$

It is also well-known that

$$
\begin{equation*}
\bar{\partial}\left[\frac{1}{s^{m}}\right] \cdot \xi \mathrm{d} s=\frac{2 \pi i}{(m-1)!} \frac{\partial^{m-1}}{\partial s^{m-1}} \xi(0) \tag{2.2}
\end{equation*}
$$

for test functions $\xi$ and $m \geqslant 1$; in particular, $\bar{\partial}\left[1 / s^{m}\right]$ has support at $\{s=0\}$. Thus

$$
\begin{equation*}
\bar{s} \bar{\partial}\left[\frac{1}{s^{m}}\right]=0, \quad \mathrm{~d} \bar{s} \wedge \bar{\partial}\left[\frac{1}{s^{m}}\right]=0 \tag{2.3}
\end{equation*}
$$

We say that a function $\chi$ on the real line is a smooth approximand of the characteristic function $\chi_{[1, \infty)}$ of the interval $[1, \infty)$, and write

$$
\chi \sim \chi_{[1, \infty)}
$$

if $\chi$ is smooth, equal to 0 in a neighborhood of 0 and 1 in a neighborhood of $\infty$. It is well-known that $\left[1 / s^{m}\right]=\lim _{\epsilon \rightarrow 0} \chi\left(|s|^{2} / \epsilon\right)\left(1 / s^{m}\right)$.

Let $t_{j}$ be coordinates in an open set $\mathcal{U} \subset \mathbb{C}^{N}$ and let $\alpha$ be a smooth form with compact support in $\mathcal{U}$. Then

$$
\begin{equation*}
\tau=\alpha \wedge\left[\frac{1}{t_{1}^{m_{1}}}\right] \ldots\left[\frac{1}{t_{k}^{m_{k}}}\right] \bar{\partial}\left[\frac{1}{t_{k+1}^{m_{k+1}}}\right] \wedge \ldots \wedge \bar{\partial}\left[\frac{1}{t_{r}^{m_{r}}}\right], \tag{2.4}
\end{equation*}
$$

where $m_{1}, \ldots, m_{r} \geqslant 1$, is a well-defined current, since it is the tensor product of one-variable currents (times $\alpha$ ). We say that $\tau$ is an elementary (pseudomeromorphic) current, and we refer to $\left[1 / t_{j}^{m_{j}}\right]$ and $\bar{\partial}\left[1 / t_{\ell}^{m_{\ell}}\right]$ as its principal value factors and residue factors, respectively. It is clear that (2.4) is commuting in the principal value factors and anti-commuting in the residue factors. We say the intersection of $\mathcal{U}$ and the coordinate plane $\left\{t_{k+1}=\cdots=t_{r}=0\right\}$ is the elementary support of $\tau$. Clearly the support of $\tau$ is contained in the intersection of the elementary support of $\tau$ and the support of $\alpha$.

Remark 2.1. - Since $\partial$ does not introduce new residue factors, $\partial \tau$ is an elementary current, cf. (2.1), whose elementary support either equals the elementary support $H$ of $\tau$ or is empty. Moreover $\bar{\partial} \tau$ is a finite sum of elementary currents, whose elementary supports are either equal to $H$ or coordinate planes of codimension 1 in $H$, cf. (2.2).

### 2.1. Definition and basic properties

Let $X$ be a reduced complex space of pure dimension $n$. Fix a point $x \in X$. We say that a germ $\mu$ of a current at $x$ is pseudomeromorphic at
$x, \mu \in \mathcal{P} \mathcal{M}_{x}$, if it is a finite sum of currents of the form

$$
\begin{equation*}
\pi_{*} \tau=\pi_{*}^{1} \ldots \pi_{*}^{m} \tau \tag{2.5}
\end{equation*}
$$

where $\mathcal{U} \subset X$ is a neighborhood of $x$,

$$
\begin{equation*}
\mathcal{U}_{m} \xrightarrow{\pi^{m}} \ldots \xrightarrow{\pi^{2}} \mathcal{U}_{1} \xrightarrow{\pi^{1}} \mathcal{U}_{0}=\mathcal{U} \tag{2.6}
\end{equation*}
$$

each $\pi^{j}: \mathcal{U}_{j} \rightarrow \mathcal{U}_{j-1}$ is either a modification, a simple projection $\mathcal{U}_{j-1} \times Z \rightarrow$ $\mathcal{U}_{j-1}$, or an open inclusion (i.e., $\mathcal{U}_{j}$ is an open subset of $\mathcal{U}_{j-1}$ ), and $\tau$ is elementary on $\mathcal{U}_{m} \subset \mathbb{C}^{N}$.

By definition the union $\mathcal{P} \mathcal{M}=\mathcal{P} \mathcal{M}_{X}=\cup_{x} \mathcal{P} \mathcal{M}_{x}$ is an open subset (of the étalé space) of the sheaf $\mathcal{C}=\mathcal{C}_{X}$ of currents, and hence it is a subsheaf, which we call the sheaf of pseudomeromorphic currents ${ }^{(1)}$. A section $\mu$ of $\mathcal{P M}$ over an open set $\mathcal{V} \subset X, \mu \in \mathcal{P} \mathcal{M}(\mathcal{V})$, is then a locally finite sum

$$
\begin{equation*}
\mu=\sum\left(\pi_{\ell}\right)_{*} \tau_{\ell} \tag{2.7}
\end{equation*}
$$

where each $\pi_{\ell}$ is a composition of mappings as in (2.6) (with $\mathcal{U} \subset \mathcal{V}$ ) and $\tau_{\ell}$ is elementary. For simplicity we will always suppress the subscript $\ell$ in $\pi_{\ell}$. If $\xi$ is a smooth form, then

$$
\begin{equation*}
\xi \wedge \pi_{*} \tau=\pi_{*}\left(\pi^{*} \xi \wedge \tau\right) \tag{2.8}
\end{equation*}
$$

Thus $\mathcal{P} \mathcal{M}$ is closed under exterior multiplication by smooth forms. Since $\bar{\partial}$ and $\partial$ commute with push-forwards it follows that $\mathcal{P} \mathcal{M}$ is closed under $\bar{\partial}$ and $\partial$, cf. Remark 2.1.

Remark 2.2. - Let $\tau$ be an elementary current with elementary support $H$. Since $H$ is the intersection of an open set $\mathcal{U}$ and a linear subspace, each of its components is irreducible, and it follows that, in fact, $\tau$ is a finite sum of currents $\tau_{\ell}$ such that the support of $\tau_{\ell}$ is contained in an irreducible component of $H$. We may therefore assume that each $\tau_{\ell}$ in (2.7) has irreducible elementary support.

Remark 2.3. - One may assume that each $\tau_{\ell}$ in (2.7) has at most one residue factor. Indeed, in [21], see also [4, Corollary 3.5], it is shown that the Coleff-Herrera product

$$
\bar{\partial}\left[1 / t_{k+1}^{m_{k+1}}\right] \wedge \ldots \wedge \bar{\partial}\left[1 / t_{r}^{m_{r}}\right]
$$

equals the Bochner-Martinelli residue current of $t_{k+1}^{m_{k+1}}, \ldots, t_{r}^{m_{r}}$, which, see, e.g., [3], is the direct image under a modification of a current of the form $\alpha \wedge \bar{\partial}[1 / f]$, cf. Example 4.18 below. It follows, cf. [6, Lemma 3.2], that (2.4)

[^1]is the direct image under another modification of a finite sum of elementary currents with at most one residue factor.

Proposition 2.4. - Assume that $\mu \in \mathcal{P} \mathcal{M}$ has support on the subvariety $V \subset X$.
(1) If the holomorphic function $h$ vanishes on $V$, then $\bar{h} \mu=0$ and $\mathrm{d} \bar{h} \wedge \mu=0$.
(2) If $\mu$ has bidegree $(*, p)$ and codim $V>p$, then $\mu=0$.

This proposition is from [10]; for the adaption to nonsmooth $X$, see [7, Proposition 2.3]. Part (1) means that the action of the current $\mu$ only involves holomorphic derivatives of test forms. We refer to part (2) as the dimension principle. We will also need, [6, Proposition 1.2]:

Proposition 2.5. - If $\pi: X^{\prime} \rightarrow X$ is a modification, then $\pi_{*}: \mathcal{P} \mathcal{M}\left(X^{\prime}\right) \rightarrow \mathcal{P} \mathcal{M}(X)$ is surjective.

### 2.2. Basic operations on pseudomeromorphic currents

Assume that $\mu$ is pseudomeromorphic on $X$ and that $V \subset X$ is a subvariety. It was proved in [10], see also [7], that the restriction of $\mu$ to the open set $X \backslash V$ has a natural pseudomeromorphic extension $\mathbb{1}_{X \backslash V} \mu$ to $X$. In [10] it was obtained as the value

$$
\begin{equation*}
\mathbb{1}_{X \backslash V} \mu:=\left.|f|^{2 \lambda} \mu\right|_{\lambda=0} \tag{2.9}
\end{equation*}
$$

at $\lambda=0$ of the analytic continuation of the current valued function $\lambda \mapsto$ $|f|^{2 \lambda} \mu$, where $f$ is any tuple of holomorphic functions such that $Z(f)=V$. It follows that

$$
\mathbb{1}_{V} \mu:=\mu-\mathbb{1}_{X \backslash V} \mu
$$

has support on $V$. It is proved in [10] that this operation extends to all constructible sets and that (1.3) holds. If $\alpha$ is a smooth form, then

$$
\begin{equation*}
\mathbb{1}_{V}(\alpha \wedge \mu)=\alpha \wedge \mathbb{1}_{V} \mu \tag{2.10}
\end{equation*}
$$

Moreover, if $\pi: X^{\prime} \rightarrow X$ is a modification, a simple projection or an open inclusion and $\mu=\pi_{*} \mu^{\prime}$, then

$$
\begin{equation*}
\mathbb{1}_{V} \mu=\pi_{*}\left(\mathbb{1}_{\pi^{-1} V} \mu^{\prime}\right) \tag{2.11}
\end{equation*}
$$

In this paper it is convenient to express $\mathbb{1}_{X \backslash V} \mu$ as a limit of currents that are pseudomeromorphic themselves.

Lemma 2.6. - Let $V$ be a germ of a subvariety at $x \in X$, let $f$ be a tuple of holomorphic functions whose common zero set is precisely $V$, let $v$ be a positive and smooth function, and let $\chi \sim \chi_{[1, \infty)}$. For each germ of a pseudomeromorphic current $\mu$ at $x$ we have

$$
\begin{equation*}
\mathbb{1}_{X \backslash V} \mu=\lim _{\epsilon \rightarrow 0} \chi\left(|f|^{2} v / \epsilon\right) \mu \tag{2.12}
\end{equation*}
$$

Because of the factor $v$, the lemma holds just as well for a holomorphic section $f$ of a Hermitian vector bundle.

In case $V$ is a hypersurface and $f$ is one single holomorphic function, or section of a line bundle, the lemma follows directly from Lemma 6 in [17] by just taking $T=f \mu$. We will reduce the general case to this lemma. The proof of this lemma relies on the proof of Theorem 1.1 in [17], which is quite involved. For a more direct proof of Lemma 2.6, see the proof of Proposition 3.4 in [1, Chapter 2].

Proof. - Let $\pi: X^{\prime} \rightarrow X$ be a smooth modification such that $\pi^{*} f=$ $f^{0} f^{\prime}$, where $f^{0}$ is a holomorphic section of a Hermitian line bundle $L \rightarrow X^{\prime}$ and $f^{\prime}$ is a nonvanishing tuple of holomorphic sections of $L^{-1}$. In view of Proposition 2.5 we can assume that $\mu=\pi_{*} \mu^{\prime}$, where $\mu^{\prime}$ is pseudomeromorphic on $X^{\prime}$. Then

$$
\left|\pi^{*} f\right|^{2} \pi^{*} v=\left|f^{0}\right|^{2}\left|f^{\prime}\right|^{2} \pi^{*} v
$$

and from $[17$, Lemma 6] we thus have that

$$
\lim _{\epsilon \rightarrow 0} \chi\left(\left|\pi^{*} f\right|^{2} \pi^{*} v / \epsilon\right) \mu^{\prime}=\mathbb{1}_{X^{\prime} \backslash \pi^{-1} V} \mu^{\prime}
$$

In view of (2.11) we get (2.12).
Remark 2.7. - Lemma 2.6 holds even if $\chi=\chi_{[1, \infty)}$. However, in general it is not obvious what $\chi\left(|f|^{2} v / \epsilon\right) \mu$ means. Let $\chi^{\delta}$ be smooth approximands such that $\chi^{\delta} \rightarrow \chi_{[1, \infty)}$. It follows from the proof of Lemma 6 in [17] that for small enough $\epsilon$, depending on $\mu, f$, and $v$, the limit $\lim _{\delta \rightarrow 0} \chi^{\delta}\left(|f|^{2} v / \epsilon\right) \mu$ exists and is independent of the choice of $\chi^{\delta}$; thus we can take it as the definition of $\chi\left(|f|^{2} v / \epsilon\right) \mu$. In fact, it turns out that after a suitable change of real coordinates one can realize $\chi\left(|f|^{2} v / \epsilon\right) \mu$ as a tensor product of two currents. In particular we get

$$
\chi\left(|f|^{2} / \epsilon\right) \frac{1}{f} \cdot \xi=\int_{|f|^{2}>\epsilon} \frac{\xi}{f},
$$

cf. (1.1).
We will need the following observation.

Lemma 2.8. - If $\mu$ has the form (2.7), then

$$
\mathbb{1}_{V} \mu=\sum_{\operatorname{supp} \tau_{\ell} \subset \pi^{-1} V} \pi_{*} \tau_{\ell}
$$

It follows from the proof below that we just as well can take the sum over all $\ell$ such that the elementary supports of $\tau_{\ell}$ are contained in $\pi^{-1} V$.

Proof. - In view of (2.11) we have that

$$
\mathbb{1}_{V} \mu=\sum_{\ell} \pi_{*}\left(\mathbb{1}_{\pi^{-1} V} \tau_{\ell}\right)
$$

If $\operatorname{supp} \tau_{\ell} \subset \pi^{-1} V$, then clearly $\mathbb{1}_{\pi^{-1} V} \tau_{\ell}=\tau_{\ell}$. We now claim that if $\operatorname{supp} \tau_{\ell}$ is not contained in $\pi^{-1} V$, then $\mathbb{1}_{\pi^{-1} V} \tau_{\ell}=0$. If $\operatorname{supp} \tau_{\ell} \not \subset \pi^{-1} V$, the elementary support $H$ of $\tau_{\ell}$ is not contained in $\pi^{-1} V$. Assume that $H$ has codimension $q$. Then $\tau_{\ell}$ is of the form $\tau_{\ell}=\alpha \wedge \tau^{\prime}$, where $\alpha$ is smooth and $\tau^{\prime}$ is elementary of bidegree $(0, q)$. It follows from (2.10) that

$$
\mathbb{1}_{\pi^{-1} V} \tau_{\ell}=\alpha \wedge \mathbb{1}_{\pi^{-1} V} \tau^{\prime}
$$

By Remark 2.2 we may assume that $H$ is irreducible, and therefore $\pi^{-1} V \cap$ $H$ has codimension at least $q+1$ in $\mathcal{U}$. Since $\mathbb{1}_{\pi^{-1} V} \tau^{\prime}$ has support on $\pi^{-1} V \cap H$ it must vanish in view of the dimension principle. Thus the lemma follows.

We now consider another fundamental operation on $\mathcal{P M}$ introduced in [10].

Proposition 2.9 ([10]). - Given a holomorphic function $h$ and a pseudomeromorphic current $\mu$ there is a pseudomeromorphic current $T$ such that $T=(1 / h) \mu$ in the open set where $h \neq 0$ and $\mathbb{1}_{\{h=0\}} T=0$.

Here $h$ may just as well be a holomorphic section of a line bundle. Clearly this current $T$ must be unique and we denote it by $[1 / h] \mu$. In [10] the current $[1 / h] \mu$ was defined as $\left.\left(|h|^{2 \lambda} \mu / h\right)\right|_{\lambda=0}$.

Remark 2.10. - Notice that ${ }^{(2)} h[1 / h] \mu=\mathbb{1}_{\{h \neq 0\}} \mu$; in particular, $h[1 / h] \mu \neq \mu$ in general. For example, $z[1 / z] \bar{\partial}[1 / z]=0$.

Since $[1 / h] \mu=(1 / h) \mu$ in $\{h \neq 0\}$ and $[1 / h] \mu=\mathbb{1}_{\{h \neq 0\}}[1 / h] \mu$, it follows from (2.12) that

$$
\begin{equation*}
\left[\frac{1}{h}\right] \mu=\lim _{\epsilon \rightarrow 0} \chi\left(|h|^{2} v / \epsilon\right) \frac{1}{h} \mu \tag{2.13}
\end{equation*}
$$

[^2]One can also define

$$
\begin{equation*}
\bar{\partial}\left[\frac{1}{h}\right] \wedge \mu:=\bar{\partial}\left(\left[\frac{1}{h}\right] \mu\right)-\left[\frac{1}{h}\right] \bar{\partial} \mu \tag{2.14}
\end{equation*}
$$

i.e., so that "Leibniz's rule" holds. Notice that if $\pi: X^{\prime} \rightarrow X$ is a modification and $\mu=\pi_{*} \mu^{\prime}$, then

$$
\begin{equation*}
\left[\frac{1}{h}\right] \mu=\pi_{*}\left(\left[\frac{1}{\pi^{*} h}\right] \mu^{\prime}\right), \quad \bar{\partial}\left[\frac{1}{h}\right] \wedge \mu=\pi_{*}\left(\bar{\partial}\left[\frac{1}{\pi^{*} h}\right] \wedge \mu^{\prime}\right) . \tag{2.15}
\end{equation*}
$$

This follows, e.g., from (2.8) and (2.13). It is also readily checked that

$$
\begin{equation*}
\bar{\partial}\left(\bar{\partial}\left[\frac{1}{h}\right] \wedge \mu\right)=-\bar{\partial}\left[\frac{1}{h}\right] \wedge \bar{\partial} \mu \tag{2.16}
\end{equation*}
$$

Remark 2.11. - Since $[1 / f][1 / g]=[1 /(f g)]=[1 / g][1 / f]$ it follows from (2.14) that

$$
\bar{\partial}\left[\frac{1}{f}\right] \cdot\left[\frac{1}{g}\right]+\left[\frac{1}{f}\right] \bar{\partial}\left[\frac{1}{g}\right]=\bar{\partial}\left[\frac{1}{g}\right] \cdot\left[\frac{1}{f}\right]+\left[\frac{1}{g}\right] \bar{\partial}\left[\frac{1}{f}\right] .
$$

However, it is not true in general that $[1 / g] \bar{\partial}[1 / f]=\bar{\partial}[1 / f] \cdot[1 / g]$. For instance, $[1 / z] \bar{\partial}[1 / z]=0$, whereas $\bar{\partial}[1 / z] \cdot[1 / z]=\bar{\partial}\left[1 / z^{2}\right]$.

We now consider tensor products and direct images under simple projections.

Lemma 2.12. - If $\mu \in \mathcal{P} \mathcal{M}_{X}$ and $\mu^{\prime} \in \mathcal{P} \mathcal{M}_{X^{\prime}}$, then $\mu \otimes \mu^{\prime} \in \mathcal{P} \mathcal{M}_{X \times X^{\prime}}$.
This is precisely [6, Lemma 3.3]. It is easy to verify that

$$
\begin{equation*}
\mathbb{1}_{V \times V^{\prime}} \mu \otimes \mu^{\prime}=\mathbb{1}_{V} \mu \otimes \mathbb{1}_{V^{\prime}} \mu^{\prime} \tag{2.17}
\end{equation*}
$$

Lemma 2.13. - Assume that $p: Z \times W \rightarrow Z$ is a simple projection. If $\mu$ is in $\mathcal{P}_{Z \times W}$ and $p^{-1} K \cap \operatorname{supp} \mu$ is compact for each compact set $K \subset Z$, then $p_{*} \mu$ is in $\mathcal{P} \mathcal{M}_{Z}$.

Proof. - Since pseudomeromorphicity is a local property, after multiplying $\mu$ if necessary by a suitable cutoff function we can assume that $\mu$ has compact support. By compactness and a partition of unity we then have a finite representation $\mu=\sum_{\ell} \pi_{*} \tau_{\ell}$. Now the lemma follows from the very definition of $\mathcal{P} \mathcal{M}$.

Example 2.14. - Assume that $\tau$ is an elementary current on $X, p$ is a simple projection $X \times X^{\prime} \rightarrow X$, and $\chi$ is any test form in $X^{\prime}$ with total integral 1 . Then the tensor product $\tau \otimes \chi$ is an elementary current in $X \times X^{\prime}$ such that $p_{*}(\tau \otimes \chi)=\tau$.

The following result provides a new, quite natural definition of pseudomeromorphicity.

## Theorem 2.15.

(1) Assume that $X$ is smooth. Then a germ of a current $\mu$ at $x \in X$ is pseudomeromorphic if and only if it is a finite sum

$$
\begin{equation*}
\mu=\sum_{\ell}\left(f_{\ell}\right)_{*} \tau_{\ell} \tag{2.18}
\end{equation*}
$$

where $f_{\ell}: \mathcal{U}_{\ell} \rightarrow X$ are holomorphic mappings and $\tau_{\ell}$ are elementary.
(2) If $X$ is a reduced space of pure dimension and $\pi: X^{\prime} \rightarrow X$ is a smooth modification, then a current $\mu$ on $X$ is pseudomeromorphic if and only if there is a pseudomeromorphic current $\mu^{\prime}$ on $X^{\prime}$ such that $\mu=\pi_{*} \mu^{\prime}$.

Proof. - By definition a germ of a pseudomeromorphic current is of the form (2.18). Now assume that $f: \mathcal{U} \rightarrow X$ is any holomorphic mapping and $\tau$ is elementary in $\mathcal{U} \subset \mathbb{C}^{N}$. Let $F: \mathcal{U} \rightarrow \mathcal{U} \times X$ be the mapping $F(s)=(s, f(s))$. Let $\widetilde{F}$ be $F$ considered as a biholomorphism onto the graph $\Gamma \subset \mathcal{U} \times X$ and let $i: \Gamma \rightarrow \mathcal{U} \times X$ be the natural injection. Then clearly $\widetilde{F}_{*} \tau$ is pseudomeromorphic on $\Gamma$ and in view of [6, Theorem 1.1(i)], $F_{*} \tau=i_{*} \widetilde{F}_{*} \tau$ is pseudomeromorphic in $\mathcal{U} \times X$. Clearly, it has compact support in $\mathcal{U} \times X$. If $p$ is the projection $\mathcal{U} \times X \rightarrow X$, we can therefore apply Lemma 2.13, and conclude that $f_{*} \tau=p_{*} F_{*} \tau$ is pseudomeromorphic in $X$. Thus part (1) is proved. Part (2) is just Proposition 2.5.

Corollary 2.16. - Assume that $f: W \rightarrow X$ is a holomorphic mapping and $X$ is smooth. If $\mu$ is pseudomeromorphic on $W$ with compact support, then $f_{*} \mu$ is pseudomeromorphic on $X$.

Proof. - We may assume that $\mu=\pi_{*} \tau$, where $\pi: \mathcal{U} \rightarrow W$ is a mapping as in the definition of pseudomeromorphicity and $\tau$ is elementary in $\mathcal{U}$. Then we can apply Theorem $2.15(1)$ to the mapping $f \circ \pi: \mathcal{U} \rightarrow X$. It follows that $f_{*} \mu=f_{*} \pi_{*} \tau=(f \circ \pi)_{*} \tau$ is pseudomeromorphic in $X$.

Remark 2.17. - Notice that in the proof of Theorem 2.15 we only used [6, Theorem $1.1(\mathrm{i})]$, which asserts that $i_{*}$ maps $\mathcal{P} \mathcal{M}_{W}$ into $\mathcal{P} \mathcal{M}_{X}$ if $i: W \rightarrow X$ is an embedding of a reduced pure-dimensional space $W$ into a manifold $X$, in the relatively simple case when $W$ is a smooth submanifold. The general case now follows from Corollary 2.16. Part (ii) of $[6$, Theorem 1.1] is a partial converse: If $\mu=i_{*} \nu$ is pseudomeromorphic in $X$ and $\mathbb{1}_{W_{\text {sing }}} \mu=0$, then $\nu$ is pseudomeromorphic on $W$. The proof of this fact relies on the possibility to make a so-called strong resolution. This means that there is a resolution $X^{\prime} \rightarrow X$ that is a biholomorphism outside $W$, and such that the strict transform of $W$ is a smooth resolution of $W$.

## 3. Action of holomorphic differential operators and vector fields

Let $X$ be a reduced analytic space of pure dimension. We already know that $\partial$ maps $\mathcal{P} \mathcal{M}_{X}$ into itself. We shall now consider a more general statement, and to this end we need the following result that is interesting in itself.

Proposition 3.1. - Assume that $\mu \in \mathcal{P M}_{x}$ where $x \in X$. If $h \in \mathcal{O}_{x}$ is not identically zero on any irreducible component of $X$ at $x$, then there is $\mu^{\prime} \in \mathcal{P} \mathcal{M}_{x}$ such that $h \mu^{\prime}=\mu$.

Remark 3.2. - By a partition of unity we can get a global such $\mu^{\prime}$ if $\mu$ and $h$ are global. If $\mu$ has compact support in $\mathcal{U} \subset X$ we can choose $\mu^{\prime}$ with compact support in $\mathcal{U}$.

Remark 3.3. - If $\mu$ has support on $V$ we may assume as well that $\mu^{\prime}$ has. Indeed, $\mu=\mathbb{1}_{V} \mu=\mathbb{1}_{V} h \mu^{\prime}=h \mathbb{1}_{V} \mu^{\prime}$, so we can replace a given solution $\mu^{\prime}$ by $\mathbb{1}_{V} \mu^{\prime}$.

Example 3.4. - Proposition 3.1 is not true if $h$ is anti-holomorphic. In fact, if $\bar{z} \mu^{\prime}=1$, then $[1 / z] \mu^{\prime}$ is equal to $1 /|z|^{2}$ outside 0 . Thus $\lim _{\epsilon \rightarrow 0} \chi\left(|z|^{2} / \epsilon\right) \mu^{\prime} / z$ does not exist, and hence $\mu^{\prime}$ cannot be pseudomeromorphic, cf. Proposition 2.9 and (2.13).

Proof of Proposition 3.1. - First assume that $\tau$ is an elementary pseudomeromorphic current in $\mathbb{C}_{t}^{N}$ and $h$ is a monomial. By induction it is enough to assume that $h=t_{1}$. If $t_{1}$ is a residue factor in $\tau$, then we just raise the power of $t_{1}$ in that factor one unit. Otherwise we take $\tau^{\prime}=\left(1 / t_{1}\right) \tau$. Then $h \tau^{\prime}=\tau$.

We may assume that $\mu=\pi_{*} \tau$, where $\pi: \mathcal{U} \rightarrow X$ and $\tau$ is elementary of the form (2.4). By Hironaka's theorem we can find a modification $\nu: \mathcal{U}^{\prime} \rightarrow$ $\mathcal{U}$ such that, locally in $\mathcal{U}^{\prime}, \nu^{*} \pi^{*} h$ is a monomial and $\nu^{*} t_{j}$ are monomials (times nonvanishing functions). By a partition of unity in $\mathcal{U}^{\prime}$ and repeated use of (2.15) it follows that $\tau$ is a finite sum of currents $\nu_{*} \tau^{\prime}$, where

$$
\tau^{\prime}:=\nu^{*} \alpha \wedge\left[\frac{1}{\nu^{*} t_{1}^{m_{1}}}\right] \ldots\left[\frac{1}{\nu^{*} t_{k}^{m_{k}}}\right] \bar{\partial}\left[\frac{1}{\nu^{*} t_{k+1}^{m_{k+1}}}\right] \wedge \ldots \wedge \bar{\partial}\left[\frac{1}{\nu^{*} t_{r}^{m_{r}}}\right] .
$$

Each such term is a sum of elementary currents $\tau_{\ell}$ in view of (2.14). By the first part of the proof there are elementary currents $\tau_{\ell}^{\prime}$ in $\mathcal{U}^{\prime}$ such that $\nu^{*} \pi^{*} h \tau_{\ell}^{\prime}=\tau_{\ell}$. Now the proposition follows in view of (2.8).

Theorem 3.5. - Assume that $X$ is smooth at $x \in X$.
(1) If $z$ is a local holomorphic coordinate system at $x$ and

$$
\begin{equation*}
\mu=\sum_{|I|=p}^{\prime} \mu_{I} \wedge \mathrm{~d} z_{I} \tag{3.1}
\end{equation*}
$$

is a germ in $\mathcal{P} \mathcal{M}_{x}$, then each $\mu_{I}$ is in $\mathcal{P} \mathcal{M}_{x}$.
(2) If $\xi$ is a germ of a holomorphic vector field, then the contraction $\xi \neg \mu$ and the Lie derivative $L_{\xi} \mu$ are in $\mathcal{P} \mathcal{M}_{x}$.

Notice that (2) is not true for anti-holomorphic vector fields. For example, $\mu=(\partial / \partial \bar{z}) \neg \bar{\partial}(1 / z)$ is a nonzero current of degree 0 with support at 0 . In view of the dimension principle, it cannot be pseudomeromorphic.

Proof. - We will first assume that $\mu$ has bidegree $(n, *)$ so that $\mu=$ $\hat{\mu} \wedge \mathrm{d} z$, where $\hat{\mu}$ has bidegree $(0, *)$, and show that $\hat{\mu}$ is pseudomeromorphic. We may assume that $\mu=\pi_{*}(\tau \wedge \mathrm{~d} s)$, where $\pi: \mathcal{U} \rightarrow X$ is a mapping as in the definition of pseudomeromorphicity, $s$ are local coordinates in $\mathcal{U} \subset \mathbb{C}^{m}$, and $\tau$ is elementary. Since $\pi$ has generically surjective differential, we can write $s=\left(s^{\prime}, s^{\prime \prime}\right)=\left(s_{1}^{\prime}, \ldots, s_{n}^{\prime}, s_{n+1}^{\prime \prime}, \ldots, s_{m}^{\prime \prime}\right)$ so that $h:=\operatorname{det}\left(\partial \pi / \partial s^{\prime}\right)=\operatorname{det}\left(\partial z / \partial s^{\prime}\right)$ is generically nonvanishing in $\mathcal{U}$. By Proposition 3.1 and Remark 3.2 there is a pseudomeromorphic $\tau^{\prime}$ with compact support in $\mathcal{U}$ such that $h \tau^{\prime}=\tau$ in $\mathcal{U}$. Now

$$
\begin{aligned}
\hat{\mu} \wedge \mathrm{d} z & =\pi_{*}(\tau \wedge \mathrm{~d} s)=\pi_{*}\left(\tau^{\prime} \wedge h \mathrm{~d} s^{\prime} \wedge \mathrm{d} s^{\prime \prime}\right)=\pi_{*}\left(\tau^{\prime} \wedge \pi^{*} \mathrm{~d} z \wedge \mathrm{~d} s^{\prime \prime}\right) \\
& = \pm \pi_{*}\left(\tau^{\prime} \wedge \mathrm{d} s^{\prime \prime}\right) \wedge \mathrm{d} z
\end{aligned}
$$

Thus $\hat{\mu}= \pm \pi_{*}\left(\tau^{\prime} \wedge \mathrm{d} s^{\prime \prime}\right)$ is pseudomeromorphic. In general, $\mu_{I} \wedge \mathrm{~d} z=$ $\pm \mu \wedge \mathrm{d} z_{I^{c}}$, where $I^{c}$ is the complementary multiindex of $I$. It follows from above that $\mu_{I}$ is pseudomeromorphic. Thus (1) follows.

The first statement of (2) follows immediately from (1), and the second one follows since $L_{\xi} \mu=\partial(\xi \neg \mu)+\xi \neg(\partial \mu)$.

### 3.1. The sheaves $\mathcal{P} \mathcal{M}_{X}^{Z}$ and $\mathcal{W}_{X}^{Z}$

Let $X$ be a reduced analytic space, let $Z \subset X$ be a (reduced) subspace of pure dimension, and denote by $\mathcal{P} \mathcal{M}_{X}^{Z}$ the subsheaf of $\mathcal{P} \mathcal{M}_{X}$ of currents that have support on $Z$. We say that $\mu \in \mathcal{P} \mathcal{M}_{X}^{Z}$ has the standard extension property, $S E P$, on $Z$ if $\mathbb{1}_{W} \mu=0$ in $\mathcal{U}$ for each subvariety $W \subset \mathcal{U} \cap Z$ of positive codimension, where $\mathcal{U}$ is any open set in $X$. Let $\mathcal{W}_{X}^{Z}$ be the subsheaf of $\mathcal{P} \mathcal{M}_{X}^{Z}$ of currents with the SEP on $Z$. In case $Z=X$ we usually write $\mathcal{W}_{X}$ rather than $\mathcal{W}_{X}^{X}$.

Example 3.6. - Note that an elementary current in $\mathcal{U}$ with elementary support $H$ is in $\mathcal{W}_{\mathcal{U}}^{H}$.

It is easy to see that Theorem 3.5 holds for $\mathcal{P} \mathcal{M}_{X}^{Z}$ as well, since neither $\partial$ nor contraction can increase support. Somewhat less obvious is that also the SEP is preserved.

Theorem 3.7. - The sheaf $\mathcal{W}_{X}^{Z}$ is invariant under $\partial$, and the statements in Theorem 3.5 hold for $\mathcal{W}_{X}^{Z}$ instead of $\mathcal{P M}$.

This theorem is a consequence of the following general equalities.
Proposition 3.8. - Assume that $\mu$ is a pseudomeromorphic current on $X$. If $V \subset X$ is any analytic subset, then

$$
\begin{equation*}
\mathbb{1}_{V} \partial \mu=\partial \mathbb{1}_{V} \mu \tag{3.2}
\end{equation*}
$$

If $\xi$ is a holomorphic vector field, then

$$
\begin{equation*}
\mathbb{1}_{V} \xi \neg \mu=\xi \neg \mathbb{1}_{V} \mu \tag{3.3}
\end{equation*}
$$

Proof. - Note that (3.3) follows in view of (2.12). Let us therefore focus on (3.2). By (1.3) it is enough to consider $V=Z(h)$, where $h$ is a nontrivial holomorphic function. Take $\chi \sim \chi_{[1, \infty)}$ and let $\chi_{\epsilon}=\chi\left(|h|^{2} / \epsilon\right)$. Now

$$
\begin{equation*}
\chi_{\epsilon} \partial \mu=\partial\left(\chi_{\epsilon} \mu\right)-\partial \chi_{\epsilon} \wedge \mu . \tag{3.4}
\end{equation*}
$$

If the last term tends to 0 when $\epsilon \rightarrow 0$, after taking limits we get that $\mathbb{1}_{h \neq 0} \partial \mu=\partial\left(\mathbb{1}_{h \neq 0} \mu\right)$, which is equivalent to (3.2). Let $\hat{\chi}(t)=t \chi^{\prime}(t)+\chi(t)$, and notice that also $\hat{\chi} \sim \chi_{[1, \infty)}$. According to Proposition 3.1 there is a pseudomeromorphic $\mu^{\prime}$ such that $\mu=h \mu^{\prime}$. The last term in (3.4) is therefore $\chi^{\prime}\left(|h|^{2} / \epsilon\right) \bar{h} \partial h \wedge \mu / \epsilon=\chi^{\prime}\left(|h|^{2} / \epsilon\right)|h|^{2} \partial h \wedge \mu^{\prime} / \epsilon=\hat{\chi}\left(|h|^{2} / \epsilon\right) \partial h \wedge \mu^{\prime}-\chi_{\epsilon} \partial h \wedge \mu^{\prime}$, which tends to $\mathbb{1}_{h \neq 0} \partial h \wedge \mu^{\prime}-\mathbb{1}_{h \neq 0} \partial h \wedge \mu^{\prime}=0$.

## 4. Almost semi-meromorphic currents

We say that a current on $X$ is semi-meromorphic if it is of the form $\omega[1 / f]$, where $f$ is a generically nonvanishing holomorphic section of a line bundle $L \rightarrow X$ and $\omega$ is a smooth form with values in $L$. For simplicity we will often omit the brackets [ ] indicating principal value in the sequel. Since furthermore $\omega[1 / f]=[1 / f] \omega$ when $\omega$ is smooth we can write just $\omega / f$.

### 4.1. The algebra $A S M(X)$

Let $X$ be a pure-dimensional reduced analytic space. We say that a current $a$ is almost semi-meromorphic in $X, a \in \operatorname{ASM}(X)$, if there is a modification $\pi: X^{\prime} \rightarrow X$ such that

$$
\begin{equation*}
a=\pi_{*}(\omega / f) \tag{4.1}
\end{equation*}
$$

where $\omega / f$ is semi-meromorphic in $X^{\prime}$. We say that $a$ is almost smooth in $X$ if one can choose $f$ to be nonvanishing. We can assume that $X^{\prime}$ is smooth because otherwise we take a smooth modification $\pi^{\prime}: X^{\prime \prime} \rightarrow X^{\prime}$ and consider the pullbacks of $f$ and $\omega$ to $X^{\prime \prime}$, cf. (2.15). If nothing else is said we tacitly assume that $X^{\prime}$ is smooth.

Notice that if $\mathcal{U} \subset X$ is an open subset, then the restriction $a_{\mathcal{U}}$ of $a \in$ $A S M(X)$ to $\mathcal{U}$ is in $A S M(\mathcal{U})$. In fact, if (4.1) holds, then $\mathcal{U}^{\prime}:=\pi^{-1} \mathcal{U} \rightarrow \mathcal{U}$ is a modification of $\mathcal{U}$, and $a_{\mathcal{U}}$ is the direct image of the restriction of $\omega / f$ to $\mathcal{U}^{\prime}$.

If $V$ has positive codimension in $\mathcal{U} \subset X$, then $\pi^{-1} V$ has positive codimension in $\mathcal{U}^{\prime}$ and $\mathbb{1}_{V} a=\pi_{*}\left(\mathbb{1}_{\pi^{-1} V}(\omega / f)\right)=\pi_{*}\left(\omega \mathbb{1}_{\pi^{-1} V}(1 / f)\right)=0$ in $\mathcal{U}$, cf. (2.11), (2.10), and the dimension principle. Thus $A S M(X)$ is contained in $\mathcal{W}(X)$.

Remark 4.1. - One can introduce a notion "locally almost semi-meromorphic current" and consider the associated sheaf. However, for the moment we have no need for such a concept.

Example 4.2. - Assume that $X=\{z w=0\} \subset \mathbb{C}^{2}$. Let $a: X \rightarrow \mathbb{C}$ be 1 and 0 on the $z$-axis and the $w$-axis, respectively, except at the origin. Then $a$ is almost smooth. Indeed the normalization $\nu: \widetilde{X} \rightarrow X$ consists of two disjoint components and $a=\nu_{*} \tilde{a}$, where $\tilde{a}$ is 0 and 1 , respectively, on these components.

Given a modification $\pi: X^{\prime} \rightarrow X$, let $\operatorname{sing}(\pi) \subset X^{\prime}$ be the (analytic) set where $\pi$ is not a biholomorphism. By the definition of a modification it has positive codimension. Let $a$ be given by (4.1) and let $Z \subset X^{\prime}$ be the zero set of $f$. By assumption also $Z$ has positive codimension. Notice that $a \in \operatorname{ASM}(X)$ is smooth outside $\pi(Z \cup \operatorname{sing}(\pi))$ which has positive codimension in $X$. We let $Z S S(a)$, the Zariski-singular support of $a$, be the smallest Zariski-closed set $V \subset X$ such that $a$ is smooth outside $V$.

Example 4.3. - Assume that $a \in A S M(X)$ is almost smooth. Then $a=\pi_{*} \omega$, where $\omega$ is smooth, and thus $Z S S(a) \subset \pi(\operatorname{sing}(\pi))$. This inclusion may be strict. For example if $a$ is smooth, then $Z S S(a)$ is empty. In this case
$\omega=\pi^{*} a$ outside $\operatorname{sing}(\pi)$ and since both sides are smooth $\operatorname{across} \operatorname{sing}(\pi)$, by continuity, then $\omega=\pi^{*} a$ everywhere in $X^{\prime}$.

Given two modifications $X_{1} \rightarrow X$ and $X_{2} \rightarrow X$, there is a modification $\pi: X^{\prime} \rightarrow X$ that factorizes over both $X_{1}$ and $X_{2}$, i.e., we have $X^{\prime} \rightarrow X_{j} \rightarrow$ $X$ for $j=1,2$. Therefore, given $a_{1}, a_{2} \in \operatorname{ASM}(X)$ we can assume that $a_{j}=\pi_{*}\left(\omega_{j} / f_{j}\right), j=1,2$. It follows that

$$
a_{1}+a_{2}=\pi_{*}\left(\frac{\omega_{1}}{f_{1}}+\frac{\omega_{2}}{f_{2}}\right)=\pi_{*} \frac{f_{2} \omega_{1}+f_{1} \omega_{2}}{f_{1} f_{2}}
$$

so that $a_{1}+a_{2}$ is in $\operatorname{ASM}(X)$ as well. Moreover, $A:=\pi_{*}\left(\omega_{1} \wedge \omega_{2} / f_{1} f_{2}\right)$ is an almost semi-meromorphic current that coincides with $a_{1} \wedge a_{2}$ outside the set $\pi\left(\operatorname{sing}(\pi) \cup V\left(f_{1}\right) \cup V\left(f_{2}\right)\right)$. If we had other representations $a_{j}=$ $\pi_{*}^{\prime}\left(\omega_{j}^{\prime} / f_{j}^{\prime}\right), j=1,2$, we would get an almost semi-meromorphic $A^{\prime}$ that coincides generically with $a_{1} \wedge a_{2}$ on $X$. Since almost semi-meromorphic have the SEP, thus $A=A^{\prime}$. Hence we can define $a_{1} \wedge a_{2}$ as $A$. Similarly, since

$$
a_{2} \wedge a_{1}=(-1)^{\operatorname{deg} a_{1} \operatorname{deg} a_{2}} a_{1} \wedge a_{2}, \quad a_{1} \wedge\left(a_{2}+a_{3}\right)=a_{1} \wedge a_{2}+a_{1} \wedge a_{3}
$$

and

$$
a_{1} \wedge\left(a_{2} \wedge a_{3}\right)=\left(a_{1} \wedge a_{2}\right) \wedge a_{3}
$$

hold generically on $X$ and because of the SEP they hold on $X$. Thus $\operatorname{ASM}(X)$ is an algebra.

Remark 4.4. - Notice that the almost smooth currents form a subalgebra of $A S M(X)$.

Example 4.5. - Clearly $Z S S\left(a_{1} \wedge a_{2}\right) \subset Z S S\left(a_{1}\right) \cup Z S S\left(a_{2}\right)$ but the inclusion may be strict. Take for instance $z_{1} / z_{2}$ and $z_{2} / z_{3}$.

Example 4.6. - The most basic example of an (almost semi-)meromorphic current is the principal value current associated with a meromorphic form. Let $f$ a be meromorphic $k$-form on $X$, i.e., locally $f=g / h$ where $h$ is a holomorphic function that is generically nonvanishing and $g$ is a holomorphic ( $k, 0$ )-form. By definition $g / h=g^{\prime} / h^{\prime}$ if and only if $g^{\prime} h-g h^{\prime}$ vanishes outside a set of positive codimension. In that case

$$
\begin{equation*}
g\left[\frac{1}{h}\right]=g^{\prime}\left[\frac{1}{h^{\prime}}\right] \tag{4.2}
\end{equation*}
$$

outside a set of positive codimension. By the dimension principle therefore (4.2) holds everywhere. Thus there is a well-defined almost semimeromorphic current $[f]$ associated with $f$. Notice that $Z S S([f])$ is contained in the pole set of the meromorphic form $f$, so unless $X$ is smooth
it may have codimension larger than 1. Actually, $\operatorname{ZSS}([f])$ is equal to the pole set of $f$. In fact, by continuity $\bar{\partial} f=0$ where $f$ is smooth, and by a classical result proved by Malgrange (at least for functions), [19], then $f$ is holomorphic there.

The following lemma will be crucial in what follows.
Lemma 4.7. - If $a$ is almost semi-meromorphic in $X$, then there is a representation (4.1) such that $f$ is nonvanishing in $X^{\prime} \backslash \pi^{-1} Z S S(a)$.

Proof. - Let $V=Z S S(a)$ and assume that we have a representation (4.1) and that $X^{\prime}$ is smooth. Let $Z$ be the union of the irreducible components of the divisor defined by $f$ that are not fully contained in $\pi^{-1} V$. Since $X^{\prime}$ is smooth, $Z$ is a Cartier divisor and thus the divisor of a section $f^{\prime}$ of some line bundle $L^{\prime} \rightarrow X^{\prime}$. It follows that $g:=f / f^{\prime}$ is a holomorphic section of $L \otimes\left(L^{\prime}\right)^{-1}$ in $X^{\prime}$ that is nonvanishing in $X^{\prime} \backslash \pi^{-1} V$. Outside $\operatorname{sing}(\pi) \cup Z \cup \pi^{-1} V$ we have that

$$
\begin{equation*}
\omega=f \pi^{*} a=f^{\prime} g \pi^{*} a \tag{4.3}
\end{equation*}
$$

By continuity, (4.3) must hold in $X^{\prime} \backslash \pi^{-1} V$ since both sides are smooth there.

We claim that $\widetilde{\omega}:=\omega / f^{\prime}$ is smooth in $X^{\prime}$. Taking this for granted, then

$$
\begin{equation*}
\pi_{*} \frac{\widetilde{\omega}}{g} \tag{4.4}
\end{equation*}
$$

is in $\operatorname{ASM}(X)$ and the zero set of $g$ is contained in $\pi^{-1} V$. Since (4.4) coincides with $a$ outside $V \cup \pi(\operatorname{sing}(\pi))$ it follows by the SEP that (4.4) indeed is equal to $a$ in $X$. Thus the lemma follows.

The claim is a local statement in $X^{\prime}$ so given a point in $X^{\prime}$ we can choose local coordinates $t$ in a neighborhood $\mathcal{U}$ of that point and consider each coefficient of the form $\omega$ with respect to these coordinates. Thus we may assume that $\omega$ is a function and that $\omega=f^{\prime} \gamma$ where $\gamma=g \pi^{*} a$ is smooth in $\mathcal{U} \backslash \pi^{-1} V$, cf. (4.3) and the comment thereafter. For all multiindices $\alpha$ thus

$$
\begin{equation*}
\frac{\partial^{\alpha} \omega}{\partial \bar{t}^{\alpha}} \bar{\partial} \frac{1}{f^{\prime}}=0 \tag{4.5}
\end{equation*}
$$

in $\mathcal{U} \backslash \pi^{-1} V$, since $f^{\prime} \bar{\partial}\left(1 / f^{\prime}\right)=0$. By assumption $Z \cap \pi^{-1} V$ has positive codimension in $Z$. By the dimension principle it follows that (4.5) holds in $\mathcal{U}$ for all $\alpha$, since $\bar{\partial}\left(1 / f^{\prime}\right)$ has support on $Z$. From [2, Theorem 1.2] we conclude that $\widetilde{\omega}$ is smooth in $\mathcal{U}$. It follows that $\widetilde{\omega}$ is smooth in $X^{\prime}$.

### 4.2. Action of $\operatorname{ASM}(X)$ on $\mathcal{P M}_{X}$

We will now extend Proposition 2.9 to general almost semi-meromorphic currents.

Theorem 4.8. - Assume that $a \in \operatorname{ASM}(X)$. For each $\mu \in \mathcal{P} \mathcal{M}(X)$ there is a unique pseudomeromorphic current $T$ in $X$ that coincides with $a \wedge \mu$ in $X \backslash Z S S(a)$ and such that $\mathbb{1}_{Z S S(a)} T=0$.

Let $V=Z S S(a)$. If such an extension $T$ exists then $T=\mathbb{1}_{X \backslash V} T=$ $\mathbb{1}_{X \backslash V} a \wedge \mu$ and so $T$ is unique. Moreover, if $h$ is a holomorphic tuple such that $Z(h)=V$, then

$$
\begin{equation*}
T=\lim _{\epsilon \rightarrow 0} \chi\left(|h|^{2} v / \epsilon\right) a \wedge \mu \tag{4.6}
\end{equation*}
$$

in view of Lemma 2.6. We will denote the extension $T$ by $a \wedge \mu$ as well.
Proof. - As observed above, if the extension $T$ exists, then (4.6) holds. Conversely, if the limit in (4.6) exists as a pseudomeromorphic current $T$ on $X$, then it must coincide with $a \wedge \mu$ in $X \backslash V$. In particular, $\chi\left(|h|^{2} v / \epsilon\right) T=$ $\chi\left(|h|^{2} v / \epsilon\right) a \wedge \mu$ for each $\epsilon>0$ and hence, taking limits and using Lemma 2.6, we get $\mathbb{1}_{X \backslash V} T=T$, i.e., $\mathbb{1}_{Z S S(a)} T=0$. To prove the theorem it is thus enough to verify that the limit in (4.6) exists as a pseudomeromorphic current.

In view of Lemma 4.7 we may assume that $a$ has the form (4.1), where $Z=Z(f)$ is contained in $\pi^{-1} V$ and $\omega / f=\pi^{*} a$ in $X^{\prime} \backslash \pi^{-1} V$. Let $\chi_{\epsilon}=$ $\chi\left(|h|^{2} v / \epsilon\right)$, so that $\pi^{*} \chi_{\epsilon}=\chi\left(\left|\pi^{*} h\right| \pi^{*} v / \epsilon\right)$. By Proposition 2.5 there is $\mu^{\prime} \in$ $\mathcal{P} \mathcal{M}\left(X^{\prime}\right)$ such that $\pi_{*} \mu^{\prime}=\mu$. Thus

$$
\chi_{\epsilon} a \wedge \mu=\chi_{\epsilon} a \wedge \pi_{*} \mu^{\prime}=\pi_{*}\left(\pi^{*} \chi_{\epsilon} \pi^{*} a \wedge \mu^{\prime}\right)=\pi_{*}\left(\pi^{*} \chi_{\epsilon} \frac{\omega}{f} \wedge \mu^{\prime}\right)
$$

In view of Proposition 2.9 and Lemma 2.6,

$$
\pi^{*} \chi_{\epsilon} \frac{\omega}{f} \wedge \mu^{\prime} \rightarrow \mathbb{1}_{X^{\prime} \backslash \pi^{-1} V} \frac{\omega}{f} \wedge \mu^{\prime}
$$

when $\epsilon \rightarrow 0$. In particular, the limit is a pseudomeromorphic current. Thus the limit in (4.6) exists and is pseudomeromorphic.

Notice that the definition of $a \wedge \mu$ is local, so that it commutes with restrictions to open subsets of $X$. Thus for each $a \in A S M(X)$ we get a linear sheaf mapping

$$
\begin{equation*}
\mathcal{P} \mathcal{M}_{X} \rightarrow \mathcal{P} \mathcal{M}_{X}, \quad \mu \mapsto a \wedge \mu \tag{4.7}
\end{equation*}
$$

Proposition 4.9. - Assume that $a \in A S M(X)$. If $W$ is an analytic subset of $\mathcal{U} \subset X$ and $\mu \in \mathcal{P} \mathcal{M}(\mathcal{U})$, then

$$
\begin{equation*}
\mathbb{1}_{W}(a \wedge \mu)=a \wedge \mathbb{1}_{W} \mu \tag{4.8}
\end{equation*}
$$

Proof. - On the one hand (4.8) holds in the open set $\mathcal{U} \backslash Z S S(a)$ by (2.10) since $a$ is smooth there. On the other hand both sides vanish on $Z S S(a)$, so (4.8) holds in all of $\mathcal{U}$; indeed $\mathbb{1}_{Z S S(a)}\left(a \wedge \mathbb{1}_{W} \mu\right)=0$ by definition, cf. Theorem 4.8, and $\mathbb{1}_{Z S S(a)} \mathbb{1}_{W}(a \wedge \mu)=\mathbb{1}_{W} \mathbb{1}_{Z S S(a)}(a \wedge \mu)=0$ in view of (1.3).

Proposition 4.10. - Each $a \in \operatorname{ASM}(X)$ induces a linear mapping

$$
\begin{equation*}
\mathcal{W}_{X}^{Z} \rightarrow \mathcal{W}_{X}^{Z}, \quad \mu \mapsto a \wedge \mu \tag{4.9}
\end{equation*}
$$

Proof. - To begin with, certainly $a \wedge \mu$ has support on $Z$ if $\mu$ has. Let $\mathcal{U}$ be an open subset of $X$ and assume that $W \subset \mathcal{U} \cap Z$ has positive codimension in $\mathcal{U} \cap Z$. Then $\mathbb{1}_{W}(a \wedge \mu)=a \wedge \mathbb{1}_{W} \mu=0$ if $\mathbb{1}_{W} \mu=0$, cf. (4.8).

Example 4.11. - Assume that $\mu$ is in $\mathcal{W}_{X}$. Then $\mu^{\prime}:=[1 / h] \mu$ is in $\mathcal{W}$ as well and if $h$ is generically nonvanishing, then $h \mu^{\prime}=h[1 / h] \mu=\mathbb{1}_{\{h \neq 0\}} \mu=$ $\mu$, cf. Remark 2.10.

Proposition 4.12. - Assume that $a_{1}, a_{2} \in A S M(X)$ and $\mu \in \mathcal{P} \mathcal{M}_{X}$. Then

$$
\begin{equation*}
a_{1} \wedge a_{2} \wedge \mu=(-1)^{\operatorname{deg} a_{1} \operatorname{deg} a_{2}} a_{2} \wedge a_{1} \wedge \mu \tag{4.10}
\end{equation*}
$$

Proof. - Notice that both sides of (4.10) coincide outside $Z S S\left(a_{1}\right) \cup$ $Z S S\left(a_{2}\right)$ and the restictions to $Z S S\left(a_{1}\right) \cup Z S S\left(a_{2}\right)$ vanish.

In particular, one of the $a_{j}$ may be a smooth form. We conclude that both (4.7) and (4.9) are $\mathcal{E}$-linear.

Proposition 4.13. - If $a_{1}, a_{2} \in A S M(X)$ and $\mu \in \mathcal{W}_{X}$, then

$$
\begin{equation*}
a_{1} \wedge a_{2} \wedge \mu=\left(a_{1} \wedge a_{2}\right) \wedge \mu, \quad\left(a_{1}+a_{2}\right) \wedge \mu=a_{1} \wedge \mu+a_{2} \wedge \mu \tag{4.11}
\end{equation*}
$$

In fact, (4.11) holds outside $V:=Z S S\left(a_{1}\right) \cup Z S S\left(a_{2}\right)$ and since $\mathbb{1}_{V} \mu=0$ the equalities follow from (4.8).

Example 4.14. - Both equalities in (4.11) may fail for a general $\mu \in$ $\mathcal{P} \mathcal{M}_{X}$. Let $a_{1}=1 / z_{1}, a_{2}=z_{1} / z_{2}, a_{3}=1 / z_{2}$, and $\mu=\bar{\partial}\left(1 / z_{1}\right)$. Then $\left(a_{1} a_{2}\right) \mu=\left(1 / z_{2}\right) \bar{\partial}\left(1 / z_{1}\right)$, but $a_{2} \mu=0$, and so $a_{1} a_{2} \mu=0$. Moreover

$$
\left(a_{1}+a_{3}\right) \mu=\frac{z_{2}+z_{1}}{z_{1} z_{2}} \bar{\partial} \frac{1}{z_{1}}=0
$$

but

$$
a_{1} \mu+a_{3} \mu=\frac{1}{z_{1}} \bar{\partial} \frac{1}{z_{1}}+\frac{1}{z_{2}} \bar{\partial} \frac{1}{z_{1}}=\frac{1}{z_{2}} \bar{\partial} \frac{1}{z_{1}} .
$$

### 4.3. Vector-valued almost semi-meromorphic currents

We will need to consider almost semi-meromorphic currents that take values in a holomorphic vector bundle $E \rightarrow X$. We say that $a \in \operatorname{ASM}(X, E)$ if there is a representation (4.1), where as before $f$ is a holomorphic section of $L \rightarrow X^{\prime}$ and now $\omega$ takes values in $L \otimes \pi^{*} E$. Clearly then $a$ is a current with values in $E$. If $\eta$ is a test form with values in the dual bundle $E^{*}$, then $a \cdot \eta=\pi_{*}\left((\omega / f) \cdot \pi^{*} \eta\right)$. Let $e_{j}$ be a local frame for $E$ in $\mathcal{U}$ and let $\xi$ be a test function with support in $\mathcal{U}$. If $\xi^{\prime}=\pi^{*} \xi, e_{j}^{\prime}=e_{j} \circ \pi$ and $\omega=\omega_{1} e_{1}^{\prime}+\omega_{2} e_{2}^{\prime}+\ldots$, then

$$
\begin{equation*}
\xi a=\sum_{j} \pi_{*}\left(\xi^{\prime} \omega_{j} / f\right) e_{j} \tag{4.12}
\end{equation*}
$$

Proposition 4.15. - Assume that $X$ is smooth. There are natural isomorphisms

$$
\begin{equation*}
A S M^{p, *}(X, E) \simeq A S M^{0, *}\left(X, \Lambda^{p} T_{1,0}^{*}(X) \otimes E\right) \tag{4.13}
\end{equation*}
$$

Proof. - First notice that if $F, G$ are vector bundles of the same rank over $X^{\prime}$ and $h$ is a holomorphic section of $\operatorname{Hom}(F, G)$ that is generically invertible, then there is a holomorphic section $g$ of $\operatorname{Hom}(G, F) \otimes \operatorname{det} G \otimes$ $(\operatorname{det} F)^{-1}$ such that $h g=s \cdot I_{G}$, where $s$ is a generically nonvanishing section of $\operatorname{det} G \otimes(\operatorname{det} F)^{-1}$.

For simplicity we assume that $E$ is a trivial line bundle; the general case is proved in the same way. Now, let $F=\pi^{*} \Lambda^{p} T_{1,0}^{*}(X)$ and $G=$ $\Lambda^{p} T_{1,0}^{*}\left(X^{\prime}\right)$. Then we have a natural mapping $h: F \rightarrow G$ as above, defined by just mapping the frame element $\mathrm{d} z_{I}$ to its pullback $\pi^{*} \mathrm{~d} z_{I}$. Clearly $h$ is an isomorphism where $\pi: X^{\prime} \rightarrow X$ is biholomorphic.

Now, if $a \in A S M^{0, *}\left(X, \Lambda^{p} T_{1,0}^{*}(X)\right)$, then we have the representation $a=\pi_{*}(\omega / f)$, where $\omega$ takes values in $F \otimes L$. Then $h \omega$ is a $(p, *)$-form in $X^{\prime}$ with values in $L$. It follows that $a^{\prime}:=\pi_{*}(h \omega / f)$ is an element in $A S M^{p, *}(X)$. We claim that $a^{\prime}=a$. By the SEP it is enough to verify the identity where $\pi$ is a biholomorphism. Let $z$ be coordinates in an open subset $\mathcal{U} \subset X \backslash \pi(\operatorname{sing} \pi)$, and let $\xi$ be a test function with support in $\mathcal{U}$. Then, cf. (4.12),

$$
\begin{aligned}
\xi a & =\sum_{|I|=p}^{\prime} \pi_{*}\left(\xi^{\prime} \omega_{I} / f\right) \wedge \mathrm{d} z_{I}=\pi_{*}\left(\xi^{\prime} \sum_{|I|=p}^{\prime} \omega_{I} / f \wedge \pi^{*} \mathrm{~d} z_{I}\right)=\pi_{*}\left(\xi^{\prime} h \omega / f\right) \\
& =\xi \pi_{*}(h \omega / f)=\xi a^{\prime}
\end{aligned}
$$

Conversely, since $h^{-1}=g / s$, if $a^{\prime} \in \operatorname{ASM}^{p, *}(X)$, then $a^{\prime}=\pi_{*}(\tilde{\omega} / f)$, where $\tilde{\omega}$ is a $(p, *)$-form with values in $L$, then $g \tilde{\omega}$ takes values in
$F \otimes \operatorname{det} G \otimes(\operatorname{det} F)^{-1} \otimes L$ and $s f$ takes values in $\operatorname{det} G \otimes(\operatorname{det} F)^{-1} \otimes L$, so that $a=\pi_{*}(g \tilde{\omega} / s f)$ is an element in $A S M^{0, *}\left(X, \Lambda^{p} T_{0,1}^{*}(X)\right)$. Again one verifies that they coincide in $X \backslash \pi(\operatorname{sing} \pi)$.

Notice that if $p=1$, then $s$ is a section of the relative canonical bundle $K_{X^{\prime} / X}=K_{X^{\prime}} \otimes \pi^{*} K_{X}^{-1}$.

### 4.4. Residues of almost semi-meromorphic currents

We shall now study the effect of $\partial$ and $\bar{\partial}$ on almost semi-meromorphic currents.

Proposition 4.16. - If $a \in \operatorname{ASM}(X)$, then $\partial a \in \operatorname{ASM}(X)$ and $b:=$ $\mathbb{1}_{X \backslash Z S S(a)} \bar{\partial} a \in A S M(X)$.

Thus we have the decomposition

$$
\begin{equation*}
\bar{\partial} a=b+r \tag{4.14}
\end{equation*}
$$

where $r:=\mathbb{1}_{Z S S(a)} \bar{\partial} a$ has support on $Z S S(a)$.
Proof. - Assume that $a=\pi_{*}(\omega / f)$ and let $D=D^{\prime}+\bar{\partial}$ be a Chern connection on $L \rightarrow X^{\prime}$. Then

$$
\partial a=\pi_{*}\left(\partial \frac{\omega}{f}\right)=\pi_{*} \frac{f \cdot D^{\prime} \omega-D^{\prime} f \wedge \omega}{f^{2}}
$$

which is in $\operatorname{ASM}(X)$.
In view of Lemma 4.7 we may assume that $Z(f) \subset \pi^{-1} V$, where $V=$ ZSS (a). Now

$$
\begin{equation*}
\bar{\partial} a=\pi_{*} \frac{\bar{\partial} \omega}{f}+\pi_{*} \bar{\partial} \frac{1}{f} \wedge \omega . \tag{4.15}
\end{equation*}
$$

By (2.11),

$$
\begin{align*}
\mathbb{1}_{X \backslash V} \bar{\partial} a & =\pi_{*}\left(\mathbb{1}_{\pi^{-1}(X \backslash V)} \frac{\bar{\partial} \omega}{f}\right)+\pi_{*}\left(\mathbb{1}_{\pi^{-1}(X \backslash V)} \bar{\partial} \frac{1}{f} \wedge \omega\right)  \tag{4.16}\\
& =\pi_{*}\left(\frac{\bar{\partial} \omega}{f}\right) ;
\end{align*}
$$

thus $\mathbb{1}_{X \backslash V} \bar{\partial} a \in A S M(X)$. For the last equality we have used Proposition 2.9 and the fact that $\bar{\partial}(1 / f)$ has support on $\pi^{-1} V$.

In the same way we have: If $a \in A S M(X, E)$ then (4.14) holds, where $b=$ $\mathbb{1}_{X \backslash Z S S(a)} \bar{\partial} a$ is in $A S M(X, E)$ and $r=\mathbb{1}_{Z S S(a)} \bar{\partial} a$ is a pseudomeromorphic current with support on $Z S S(a)$ that takes values in $E$.

Clearly the decomposition (4.14) is unique. We call $r=r(a)$ the residue (current) of $a$. Notice that if $a$ is almost smooth, then $r(a)=0$.

Remark 4.17. - If $a=\pi_{*}(\omega / f)$ is any representation of $a$, then still (4.15) holds, and since the first term is in $\operatorname{ASM}(X)$ we conclude that

$$
r(a)=\pi_{*}\left(\bar{\partial} \frac{1}{f} \wedge \omega\right) .
$$

Notice that the current $\bar{\partial}(1 / f)$ is the residue of the principal value current $1 / f$. Similarly, the residue currents introduced, e.g., in $[3,9,21]$ can be considered as residues of certain almost semi-meromorphic currents, generalizing $1 / f$.

Example 4.18. - Let us describe the construction of the residue currents in [3]. Let $f$ be a holomorphic section of a Hermitian vector bundle $E \rightarrow X$, and let $\sigma$ be the section over $X \backslash Z(f)$ of the dual bundle $E^{*}$ with minimal norm such that $f \sigma=1$. We can find a modification $\pi: X^{\prime} \rightarrow X$ that is a biholomorphism $X^{\prime} \backslash \pi^{-1} Z(f) \simeq X \backslash Z(f)$ such that $\pi^{*} f=f^{0} f^{\prime}$, where $f^{0}$ is a holomorphic section of a line bundle $L \rightarrow X^{\prime}$, $\operatorname{div} f^{0}$ is contained in $\pi^{-1} Z(f)$, and $f^{\prime}$ is a nonvanishing section of $\pi^{*} E \otimes L^{-1}$. Then

$$
\pi^{*} \sigma=\sigma^{\prime} / f^{0}
$$

where $\sigma^{\prime}$ is a smooth section of $\pi^{*} E^{*} \otimes L$. Thus

$$
\pi^{*}\left(\sigma \wedge(\bar{\partial} \sigma)^{k-1}\right)=\frac{\sigma^{\prime} \wedge\left(\bar{\partial} \sigma^{\prime}\right)^{k-1}}{\left(f^{0}\right)^{k}}
$$

is a section of $\Lambda^{k}\left(\pi^{*} E \oplus T_{0,1}^{*}\left(X^{\prime}\right)\right)$ in $X^{\prime} \backslash \pi^{-1} Z(f)$; for the reader's convenience note that $\bar{\partial} \sigma$ has even degree in $\Lambda^{k}\left(\pi^{*} E \oplus T_{0,1}^{*}\left(X^{\prime}\right)\right)$. It follows that

$$
U_{k}:=\sigma \wedge(\bar{\partial} \sigma)^{k-1}
$$

has an extension to an almost semi-meromorphic section of $\Lambda^{k}\left(E \oplus T_{0,1}^{*}(X)\right)$, as the push-forward of $\sigma^{\prime} \wedge\left(\bar{\partial} \sigma^{\prime}\right)^{k-1} /\left(f^{0}\right)^{k}$. Clearly $Z S S\left(U_{k}\right) \subset Z(f)$. Now the residue current $R$ in [3] is the residue of the almost semi-meromorphic current $U=\sum_{k} U_{k}$. More precisely, if $\delta_{f}$ denotes interior multiplication by $f$, then $\left(\delta_{f}-\bar{\partial}\right) U=1-R$, i.e., $\bar{\partial} U=R+\delta_{f} U-1$, where $R$ is the residue and $\delta_{f} U-1$ is almost semi-meromorphic. If $E$ is trivial with trivial metric, the coefficients of $R$ are the Bochner-Martinelli residue currents introduced in [21].

Clearly Theorem 4.8 extends to vector-valued currents. As a consequence of this theorem we can define products of residues of almost semi-meromorphic currents and pseudomeromorphic currents:

Definition 4.19. - For $a \in A S M(X, E)$ and $\mu \in \mathcal{P} \mathcal{M}_{X}$ we define

$$
\begin{equation*}
\bar{\partial} a \wedge \mu:=\bar{\partial}(a \wedge \mu)-(-1)^{\operatorname{deg} a} a \wedge \bar{\partial} \mu, \tag{4.17}
\end{equation*}
$$

where $a \wedge \mu$ and $a \wedge \bar{\partial} \mu$ are defined as in Theorem 4.8. Moreover we define

$$
r(a) \wedge \mu:=\mathbb{1}_{Z S S(a)} \bar{\partial} a \wedge \mu
$$

Thus $\bar{\partial} a \wedge \mu$ is defined so that the Leibniz rule holds. It is easily checked that

$$
\begin{equation*}
r(a) \wedge \mu=\lim _{\epsilon \rightarrow 0} \bar{\partial} \chi\left(|h|^{2} v / \epsilon\right) a \wedge \mu \tag{4.18}
\end{equation*}
$$

if $Z(h)=Z S S(a)$. In particular this gives a way of defining products of $\bar{\partial}$ and residues of almost semi-meromorphic currents. For example, the Coleff-Herrera product $\bar{\partial}\left(1 / f_{1}\right) \wedge \ldots \wedge \bar{\partial}\left(1 / f_{p}\right)$ can be defined by inductively applying (4.17). In [5] the first author defined products of more general residue currents in this way.

Notice that in general $a_{1} \wedge \bar{\partial} a_{2}$ is not equal to $\pm \bar{\partial} a_{2} \wedge a_{1}$, cf. Remark 2.11, and neither is

$$
\begin{equation*}
\bar{\partial} a_{1} \wedge \bar{\partial} a_{2}= \pm \bar{\partial} a_{2} \wedge \bar{\partial} a_{1} \tag{4.19}
\end{equation*}
$$

in general; take, e.g., $a_{1}=1 / z$ and $a_{2}=1 / z w$.
Theorem 4.20. - Assume that $a_{1}, \ldots, a_{p}$ are almost semi-meromorphic currents of degree $\left(*, k_{1}-1\right), \ldots,\left(*, k_{p}-1\right)$, respectively, and that

$$
\begin{equation*}
\operatorname{codim}\left(Z S S\left(a_{i_{1}}\right) \cap \cdots \cap Z S S\left(a_{i_{r}}\right)\right) \geqslant k_{i_{1}}+\cdots+k_{i_{r}} \tag{4.20}
\end{equation*}
$$

for all $\left\{i_{1}, \ldots, i_{r}\right\} \subset\{1, \ldots, p\}$. Then

$$
\begin{align*}
\bar{\partial} a_{1} \wedge \ldots \wedge & \bar{\partial} a_{j} \wedge \bar{\partial} a_{j+1} \wedge \ldots \wedge \bar{\partial} a_{p}  \tag{4.21}\\
& =(-1)^{\left(\operatorname{deg} a_{j}+1\right)\left(\operatorname{deg} a_{j+1}+1\right)} \bar{\partial} a_{1} \wedge \ldots \wedge \bar{\partial} a_{j+1} \wedge \bar{\partial} a_{j} \wedge \ldots \wedge \bar{\partial} a_{p}
\end{align*}
$$

Remark 4.21. - In fact, one can modify the proof below so that one can replace any factor $\bar{\partial} a_{i}$ in (4.21) by $a_{i}$. More precisely, let $b_{i}$ be either $a_{i}$ or $\bar{\partial} a_{i}$ for $i=1, \ldots, p$. Then

$$
\begin{align*}
b_{1} \wedge \ldots \wedge b_{j} \wedge b_{j+1} \wedge \ldots & \ldots b_{p}  \tag{4.22}\\
& =(-1)^{\operatorname{deg} b_{j} \cdot \operatorname{deg} b_{j+1}} b_{1} \wedge \ldots \wedge b_{j+1} \wedge b_{j} \wedge \ldots \wedge b_{p} .
\end{align*}
$$

Remark 4.22. - If the almost semimeromorphic parts of $\bar{\partial} a_{i}$ vanish, then it is enough to assume

$$
\begin{equation*}
\operatorname{codim}\left(Z S S\left(a_{1}\right) \cap \cdots \cap Z S S\left(a_{p}\right)\right) \geqslant k_{1}+\cdots+k_{p} \tag{4.23}
\end{equation*}
$$

Indeed, note that in this case the currents in (4.21) have support on $V:=$ $Z S S\left(a_{1}\right) \cap \cdots \cap Z S S\left(a_{p}\right)$. Thus it is enough to prove (4.21) in a neighborhood of $x \in V$, and there (4.23) implies (4.20).

In particular, the Coleff-Herrera product $\bar{\partial}\left(1 / f_{1}\right) \wedge \ldots \wedge \bar{\partial}\left(1 / f_{p}\right)$ is (anti-)commutative in its factors if the codimension of $\left\{f_{1}=\cdots=f_{p}=0\right\}$ is at least $p$.

Proof. - Let $V_{j}=Z S S\left(a_{j}\right)$. Moreover, let $b_{i}$ be either an almost semimeromorphic current or $\bar{\partial}$ of an semi-meromorphic current for $i=1, \ldots, r$, cf. Remark 4.21, and assume that $\alpha$ is smooth. Then note that

$$
\begin{align*}
b_{1} \wedge \ldots \wedge b_{\ell} \wedge \alpha \wedge b_{\ell+1} \wedge \ldots & \ldots b_{r}  \tag{4.24}\\
& =(-1)^{\operatorname{deg} \alpha\left(\operatorname{deg} b_{1}+\cdots+\operatorname{deg} b_{\ell}\right)} \alpha \wedge b_{1} \wedge \ldots \wedge b_{r}
\end{align*}
$$

Assume that

$$
\begin{align*}
& \bar{\partial} a_{1} \wedge \ldots \wedge \bar{\partial} a_{j-1} \wedge a_{j} \wedge \bar{\partial} a_{j+1} \wedge \ldots \wedge \bar{\partial} a_{p}  \tag{4.25}\\
= & (-1)^{\operatorname{deg} a_{j}\left(\operatorname{deg} a_{j+1}+1\right)} \bar{\partial} a_{1} \wedge \ldots \wedge \bar{\partial} a_{j-1} \wedge \bar{\partial} a_{j+1} \wedge a_{j} \wedge \bar{\partial} a_{j+2} \wedge \ldots \wedge \bar{\partial} a_{p}
\end{align*}
$$

Applying $\bar{\partial}$ to (4.25) yields (4.21) in view of (4.17).
To prove (4.25) we will proceed by induction. First assume that $p=2$. Then in view of (4.24),

$$
\begin{equation*}
a_{1} \wedge \bar{\partial} a_{2}=(-1)^{\operatorname{deg} a_{1}\left(\operatorname{deg} a_{2}+1\right)} \bar{\partial} a_{2} \wedge a_{1} \tag{4.26}
\end{equation*}
$$

where $a_{1}$ or $a_{2}$ is smooth, i.e., outside $V_{1} \cap V_{2}$. Because of the assumption (4.20), (4.26) holds in all of $X$ by the dimension principle. Next, assume that (4.25) holds for $p=\ell$. In view of (4.24), (4.25) holds for $p=\ell+1$, where $a_{j}$ or $a_{j+1}$ is smooth. Moreover, by (4.24) and the assumption that (4.25) holds for $p=\ell$, (4.25) holds for $p=\ell+1$, where (at least) one of $a_{1}, \ldots, a_{j-1}, a_{j+2}, \ldots, a_{\ell+1}$ is smooth. Thus (4.25) holds for $p=\ell+1$ outside $V_{1} \cap \cdots \cap V_{\ell+1}$, and thus by (4.20) and the dimension principle it holds in all of $X$. Hence (4.25) and thus (4.21) hold for all $p$.

The following example shows that $r(a)=0$ does not imply that $r(a) \wedge \mu=$ 0 . This points out the importance of keeping in mind that $\mu \mapsto r(a) \wedge \mu$ is an operator on $\mathcal{P} \mathcal{M}_{X}$ rather than a "product".

Example 4.23. - Let us consider the setting in Example 4.18. Assume in addition that $Z(f)$ has codimension at least 2. Note that then $r(\sigma)=0$ by the dimension principle, since it has bidegree $(0,1)$ and support on $Z(f)$, which has codimension $\geqslant 2$. However, if $\tau$ is the almost semi-meromorphic part of $\bar{\partial} U$, then $r(\sigma) \wedge \tau$ is the residue current $R$ from [3] which is nonzero, cf. Example 4.18.

Remark 4.24. - There are other (weighted) approaches to products of residue currents, see, e.g. [20, 26], which coincide with the products above under suitable conditions.

### 4.5. Action of holomorphic differential operators and vector fields

Finally we prove that $A S M(X)$ is preserved under the action of holomorphic vector fields.

THEOREM 4.25. - Let $\xi$ be a holomorphic vector field on a smooth manifold $X$. If $a \in A S M(X)$, then the contraction $\xi \neg a$ and the Lie derivative $L_{\xi} a$, a priori defined on $X \backslash Z S S(a)$, have extensions as elements in ASM(X).

Since the extensions, if they exist, must be unique, we can simply say that $\xi \neg a$ and $L_{\xi} a$ are in $A S M(X)$.

Proof. - Let $\pi: X^{\prime} \rightarrow X$ be a modification so that $a$ has the form (4.1). Then $\xi^{\prime}:=\pi^{*} \xi$ is a global section of $\pi^{*} T(X)$, that is the natural lifting of $\xi$ to $T\left(X^{\prime}\right)$ over $X^{\prime} \backslash \operatorname{sing}(\pi)$. By duality the mapping $\pi^{*} T_{1,0}^{*}(X) \rightarrow$ $T_{1,0}^{*}\left(X^{\prime}\right)$ from the proof of Proposition 4.15 induces a holomorphic mapping $T\left(X^{\prime}\right) \rightarrow \pi^{*} T(X)$ that is the identity outside $\operatorname{sing}(\pi)$. If $h$ denotes this dual map, by the first part of the same proof there is a holomorphic mapping $g: \pi^{*} T(X) \rightarrow T\left(X^{\prime}\right) \otimes K_{X^{\prime} / X}$ such that $h g=s I_{\pi^{*} T(X)}$, where $s$ is a holomorphic section of $K_{X^{\prime} / X}$. Thus $g \xi^{\prime} / s$ is a semi-meromorphic vector field on $X^{\prime}$ that coincides with $\xi^{\prime}$ on $X^{\prime} \backslash \operatorname{sing}(\pi)$. Moreover, $b:=s \xi^{\prime}$ is smooth. Outside $\pi(\operatorname{sing}(\pi)) \cup Z S S(a)$ we now have that

$$
\xi \neg a=\pi_{*}\left(\frac{\xi^{\prime} \neg \omega}{f}\right)=\pi_{*}\left(\frac{b \neg \omega}{s f}\right)
$$

and it is clear that the right hand side defines an almost semi-meromorphic current in $X$. Finally, $L_{\xi} a=\xi \neg(\partial a)+\partial(\xi \neg a)$ is in $A S M(X)$ in view of Proposition 4.16.

By similar arguments one can prove that $\mathcal{L} a$ is in $A S M(X)$ if $a$ is an almost semi-meromorphic $(0, q)$-current and $\mathcal{L}$ is any (global) holomorphic differential operator. More precisely, one can show that $\mathcal{L} a=$ $\pi_{*}\left(s^{-N} \mathcal{L}^{\prime}(\omega / f)\right)$ for some $N$, where $s$ is the section of $K_{X^{\prime} / X}$ in the proof above and $\mathcal{L}^{\prime}$ is a holomorphic differential operator (with values in $K_{X^{\prime} / X}^{N}$ ).

Corollary 4.26. - Let $X$ be an open subset of $\mathbb{C}_{z}^{n}$. If

$$
\begin{equation*}
a=\sum_{|I|=p}^{\prime} a_{I} \wedge \mathrm{~d} z_{I} \tag{4.27}
\end{equation*}
$$

is in $\operatorname{ASM}(X)$, then each $a_{I}$ is in $\operatorname{ASM}(X)$. If $a \in A S M(X)$ has bidegree $(0, *)$, then $\partial a / \partial z_{j}$ is in $A S M(X)$ for each $j$.

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[^1]:    ${ }^{(1)}$ The definition here is from [7]; in the original definition in [10] simple projections were not included.

[^2]:    ${ }^{(2)}$ We have not exluded the possibility that $h$ vanishes identically on some (or all) irreducible components of $X$.

