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Annales mathématiques Blaise Pascal, tome 1, nº 1 (1994), p. 21-26 ">http://www.numdam.org/item?id=AMBP_1994_1_1_21_0>

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Ann. Math. Blaise Pascal, Vol. 1, N° 1, 1994, pp. 21-26

GENERATING FUNCTION AND ORTHOGONALITY PROPERTY OF A CLASS OF POLYNOMIALS OCCURRING IN QUANTUM MECHANICS

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ABSTRACT: In this paper, we present a generating function and an orthogonality property of a class of polynomials occuring in quantum mechanics.

Key words : Generating function, Orthogonality property, Hermite Polynomials, Quantum mechanics.

AMS (MOS) : Subject classification : 33C25, 81

INTRODUCTION: The object of this paper is to present a generating function and an orthogonality property of the polynomials ${}_{1}F_{1}(-n; b + 3/2; x^{2})$, which occurs in the radical wave function of isotropic harmonic oscillator [4, p. 36, (6.60)].

The generating function for the polynomials ${}_{1}F_{1}(-n; b+3/2; x^{2})$ has been obtained as a particular case of the generating function of *B*-polynomials, which has recently been defined by the author [2]. We obtain the orthogonality property of the polynomials ${}_{1}F_{1}(-n; b+3/2; x^{2})$ as a bonus in our attempts to establish an orthogonality property of *B*polynomials. We shall use the symbol $H_{n}^{b}(x)$ to denote the polynomials ${}_{1}F_{1}(-n; b+3/2; x^{2})$. It is interesting to note that the polynomials $H_n^b(x)$ appear to lead to the generalization of the Hermite polynomials $H_n(x)$ [5, p. 380, (25)].

We visualize at least three orthogonality properties of the B-polynomials for different weight functions on different intervals. However, we have not been successful to establish any of them. The proofs are difficult in view of the general nature of B-polynomials.

In what follows for sake of brevity, the symbol a_p is used to denote $a_1, ..., a_p$, the symbol $1 - a_p - m$ is used to denote $1 - a_1 - m, ..., 1 - a_p - m$ and the notation $\prod_{j=1}^{p} (a_j)_m$

stands for the product $(a_1)_m \dots (a_p)_m$. Further, the expression

$${}_{p}F_{q}\begin{bmatrix}a_{p};z\\b_{q}\end{bmatrix} = \sum_{n=0}^{\infty} \frac{(a_{1})_{n}...(a_{p})_{n}z^{n}}{(b_{1})_{n}...(b_{q})_{n}n!}$$
(1.0)

is known as the generalized hypergeometric series or generalized hypergeometric function. Here p and q are positive integers or zero, and we assume that the variable z, the numerator parameters $a_1, ..., a_p$ and the denominator parameters $b_1, ..., b_q$ take on complex values, provided that no $b_j(j = 1, ..., q)$ is zero or a negative integer.

Recently [2], we have defined the *B*-polynomials :

$$B_{m}(x) = \frac{\prod_{j=1}^{p} (a_{j})_{m}}{\prod_{j=1}^{q} (b_{j})_{m}} r + q + 1^{F}s + p \begin{bmatrix} c_{r}, 1 - b_{q} - m, -m; \frac{\beta}{\alpha}x(-1)^{p-q-1} \\ d_{s}, 1 - a_{p} - m \end{bmatrix} (\alpha)^{m}, \quad (1.1)$$

by means of the generating function :

$${}_{p}F_{q}\left[\begin{array}{c}a_{p};\alpha t\\b_{q}\end{array}\right]r^{F}s\left[\begin{array}{c}c_{r};\beta xt\\d_{s}\end{array}\right] = \sum_{m=0}^{\infty}\frac{(\alpha t)^{m}}{m!}\frac{\prod_{j=1}^{r}(a_{j})_{m}}{\prod_{j=1}^{q}(b_{j})_{m}}$$
$$\cdot_{r}+q+1^{F}s+p\left[\begin{array}{c}c_{r},1-b_{q}-m,-m;\frac{\beta}{\alpha}x(-1)^{p-q-1}\\d_{s},1-a_{p}-m\end{array}\right]$$
(1.2)

The generating function of the polynomials $H_n^b(x)$:

In (1.2), putting $\alpha = \beta = 1, p = q = r = 0, s = 1, d_1 = b + 3/2$, and setting t^2 for t and $-x^2$ for x, we obtain the generating function for ${}_1F_1(-m; b + 3/2; x^2)$:

Generating function and orthogonality

$${}_{0}F_{0}(-;-;t^{2}){}_{0}F_{1}(-;b+3/2;-t^{2}x^{2}) = \sum_{m=0}^{\infty} \frac{t^{2m}}{m!} {}_{1}F_{1}(-m;b+3/2;x^{2})$$
(1.3)

In (1.3), setting
$${}_{0}F_{0}(-;-;t^{2}) = e^{t^{2}}, {}_{0}F_{1}(-;b+3/2;-t^{2}x^{2}) =$$

 $(tx)^{b/2+1/4}\Gamma(b+3/2)J_{b+1/2}(2\sqrt{t}x) \text{ and } {}_{1}F_{1}(-m,b+3/2;x^{2}) = H_{m}^{b}(x), \text{ we have}$

$$e^{t^{2}}(tx)^{b/2+1/4}\Gamma(b+3/2)J_{b+1/2}(2\sqrt{t}x) = \sum_{m=0}^{\infty} \frac{t^{2m}}{m!}H_{m}^{b}(x)$$
(1.4)

The following formulae are required in the proofs :

The integral :

$$\int_{-\infty}^{\infty} x^{2u} e^{-x^2} p^F q \begin{bmatrix} a_p; zx^2 \\ b_q \end{bmatrix} dx$$

= $\Gamma(u+1/2)_{p+1} F_q \begin{bmatrix} a_p; zx^2 \\ b_q \end{bmatrix},$ (1.5)

where p < q + 1 (or p = q + 1 and |z| < 1), u = 0, 1, 2,

The integral (1.5) can easily be established by expressing the hypergeometric function in the integrand as [1, p. 322, (10.1)] and interchanging the order of integration and summation, which is justified due to the absolute convergence of the integral and summation involved in the process, and evaluating the inner-integral with the help of the following integral :

$$\int_{-\infty}^{\infty} x^{2n} e^{-x^2} dx = \Gamma(n+1/2), \ n = 0, 1, 2, \dots$$
(1.6)
The integral:

$$\int_{-\infty}^{\infty} x^{2u} e^{-x^2} p^F q \begin{bmatrix} a_p; zx^2 \\ b_q \end{bmatrix} r F_s \begin{bmatrix} c_r; yx^2 \\ d_s \end{bmatrix} dx$$

$$= \sum_{m=0}^{\infty} \frac{\prod_{j=1}^{r} (c_j)_m}{\prod_{j=1}^{s} (d_j)_m} \frac{y^m}{m!} \Gamma(m+u+1/2)_{p+1} F_q \begin{bmatrix} a_p, m+u+1/2; x \\ b_q \end{bmatrix},$$
(1.7)

where in addition to the conditions of (1.5), r < s + 1 (or r = s + 1 and |y| < 1).

To derive (1.7), we use the series representation of $_{r}F_{s}$ interchange the order of integration and summation and evaluate the resulting integral with the help of (1.5).

The Vandermonde's theorem [3, p. 110, (4.1.2)]:

23

$${}_{2}F_{1}\left[\begin{array}{c} -n, b; 1\\ c \end{array} = \frac{(c-b)_{n}}{(c)_{n}}, n = 0, 1, 2, ...;$$
(1.8)

The modified form of the relation [1, p. 308, (9.37)]:

$$H_{2n}(x) = (-1)^n (2)^{2n} (1/2)_n F_1 \begin{bmatrix} -n; x^2 \\ 1/2 \end{bmatrix}$$
(1.9)

The modified from of the relation [1, p. 312, (6)]:

$$H_{2n+1}(x) = (-1)^n 2^{2n+1} (3/2)_{n-1} F_1 \begin{bmatrix} -n; x^2 \\ 3/2 \end{bmatrix}$$
(1.10)

The Legendre duplication formula [1, p. 58, (2.24)]

$$2^{2x-1}\Gamma(x)\Gamma(x+1/2) = \sqrt{\pi}\Gamma(2x)$$
(1.11)

The following well known relations [1, pp. 275, 323]:

$$_{0}F_{0}(-;-;x) = e^{x}$$
 (1.12)

$${}_{0}F_{1}\begin{bmatrix} -; -\frac{x^{2}}{4}\\ 1/2 \end{bmatrix} = \cos x \tag{1.13}$$

$$x_0 F_1 \begin{bmatrix} -; -\frac{x^2}{4} \\ 3/2 \end{bmatrix} = sinx \tag{1.14}$$

$$(-k)_n = \begin{cases} 0, n > k \\ (k = 1, 2, 3, ...) \\ (-1)^n n!, k = n \end{cases}$$
(1.15)

2. ORTOGONALITY PROPERTY OF THE POLYNOMIALS $H_{x}^{b}(x)$.

The polynomials $H_n^b(x)$ are orthogonal with weight $x^{2(b+1)}e^{-x^2}$ on the interval $(-\infty,\infty)$, i.e.

$$\int_{-\infty}^{\infty} x^{2(b+1)} e^{-x^2} H_n^b(x) H_k^b(x) dx = \begin{cases} 0, k \neq n \\ \frac{\Gamma(b+3/2)n!}{(b+3/2)n}, k = n \end{cases}$$
(2.1)
where $b = -1, 0, 1, 2, \dots$

<u>PROOF</u>. In (1.7), setting $y = z = 1, u = b + 1, p = q = r = s = 1, a_1 = -n, b_1 = b + 3/2; c_1 = -k, d_1 = b + 3/2$, we have

$$\int_{-\infty}^{\infty} x^{2(b+1)} e^{-x^2} {}_1F_1(-n;b+3/2;x^2) {}_1F_1(-k;b+3/2;x^2) dx$$

= $\sum_{m=0}^{\infty} \frac{(-k)_m}{(b+3/2)_m} \frac{1}{m!} \Gamma(m+b+3/2) {}_2F_1 \left[\frac{-n,m+b+3/2;1}{b+3/2} \right]$ (2.2)

Now, using the notation $H_n^b(x)$ for ${}_1F_1(-n; b+3/2; x^2)$ and Vandermonde's theorem (1.8), (2.2) reduces to the form :

$$\int_{-\infty}^{\infty} x^{2(b+1)} e^{-x^2} H_n^b(x) H_k^b(x) dx$$

= $\sum_{m=0}^{\infty} \frac{\Gamma(b+3/2)(-k)_m(-m)_n}{m!(b+3/2)_n}$ (2.3)

From (1.15), it is evident that all terms of the series (2.3) are zero for $m > k \neq n$ and $m < n \neq k$.

If
$$k = n = m$$
, we have

$$\int_{-\infty}^{\infty} x^{2(b+1)} e^{-x^2} \left\{ H_n^b(x) \right\}^2 dx = \frac{\Gamma(b+3/2)n!}{(b+3/2)_n}$$
(2.4)
This proves (2.1)

3. THE POLYNOMIALS $H_n^b(x)$ AND THE HERMITE POLYNOMIALS $H_n(x)$.

(a) Generating functions

(i) In (1.3), putting b = -1, and applying (1.9), (1.11), (1.12) and (1.13), it reduces to the generating function [1, p. 174, 2(a)] for the Hermite polynomials.

(ii) In (1.3), setting b = 0, and using (1.10), (1.11), (1.12) and (1.14), it yields the generating function [1, p. 174, 2(b)] for the Hermite polynomials.

(b) Orthogonality properties

(i) In (2.1), putting b = -1, and applying (1.9), (1.11), (1.12) and (1.13), we obtain the following orthogonality property of the Hermite polynomials :

$$\int_{-\infty}^{\infty} e^{-x^2} H_{2n}(x) H_{2k}(x) dx = \begin{cases} 0, k \neq n \\ 2^{2n}(2n)! \sqrt{\pi}, k = n \end{cases}$$
(3.1)

(ii) In (2.1), setting b = 0, and using (1.10), (1.11), (1.12) and (1.14), it yields the following orthogonality property of the Hermite polynomials:

$$\int_{-\infty}^{\infty} e^{-x^2} H_{2n+1}(x) H_{2k+1}(x) dx = \begin{cases} 0, k \neq n \\ 2^{2n+1}(2n+1)! \sqrt{\pi}, k = n \end{cases}$$
(3.2)

From (3.1) and (3.2), the orthogonality property of the Hermite polynomials [1, pp. 170-171, (5.17) - (5.22)] follows.

ACKNOWLEDGEMENT

I wish to express my sincere thanks to the referee for his useful suggestions for the revision of this paper.

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manuscrit reçu le 26 Octobre 1992