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# Hideo Kozono <br> Hermann Sohr <br> <br> On a new class of generalized solutions for the Stokes <br> <br> On a new class of generalized solutions for the Stokes equations in exterior domains 

 equations in exterior domains}

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# On a New Class of Generalized Solutions for the Stokes Equations in Exterior Domains 

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## Introduction

Let $n \geq 2$ and let $\Omega$ be an exterior domain in $\mathbb{R}^{\mathbf{}}$, i.e., a domain having a compact complement $\mathbb{R}^{n} / \Omega$, and assume that the boundary $\partial \Omega$ is of class $C^{2+\mu}$ with $0<\mu<1$. Consider the following boundary value problem for the Stokes equations in $\Omega$ :

$$
\begin{equation*}
-\Delta u+\nabla p=f, \quad \operatorname{div} u=g \text { in } \Omega, \quad u=0 \text { on } \partial \Omega \tag{S}
\end{equation*}
$$

where $u=\left(u^{1}(x), \cdots, u^{n}(x)\right)$ and $p=p(x)$ denote the unknown velocity and pressure, respectively; $f=\left(f^{1}(x), \cdots, f^{n}(x)\right)$ and $g=g(x)$ denote the given external force and the scalar divergence, respectively.

The purpose of the present paper is to extend the well-known concept of generalized solutions $u$ of (S) having a finite Dirichlet integral

$$
\begin{equation*}
\int_{\Omega}|\nabla u(x)|^{2} \mathrm{~d} x<\infty \tag{D}
\end{equation*}
$$

(see, e.g., Chang-Finn [11], Finn [14], Fujita [16], Heywood [21]). We consider here a much larger class of the generalized solutions $u$ of $(S)$ satisfying

$$
\begin{equation*}
\int_{\Omega}|\nabla \cdot u(x)-A|^{q} \mathrm{~d} x<\infty \text { with some matrix } A \tag{q}
\end{equation*}
$$

where $1<q<\infty$. In particular, setting $A \equiv 0$, we treat the class

$$
\begin{equation*}
\int_{\Omega}|\nabla u(x)|^{q} \mathrm{~d} x<\infty \tag{q}
\end{equation*}
$$

which generalizes the Dirichlet integral to $L^{q}$-spaces.

In the class $\left(C L_{q}\right)$, we can investigate the motion of the fluid past an obstacle rotating around its axis. Such a fluid motion is governed by ( $S$ ) with the boundary condition at infinity
$(B . C .)_{\infty}$

$$
u(x) \rightarrow A x+a \quad \text { as } x \rightarrow \infty
$$

where $A$ denotes a skew-symmetric matrix and the vector $a$ is a constant. Another physical phenomenon described in the class ( $C L_{q}$ ) is the flow due to an obstacle embedded in a pure straining tensor: far from the obstacle the fluid is in a pure stretching specified by the rate-of-strain tensor $A$, with $\operatorname{Tr} A=0$. Then the velocity $u$ can be written as

$$
u(x)=A x+u_{0}(x)
$$

where $u_{0}$ represents the changes due to the presence of the obstacle, with $u_{0}(x)$ small for large $|x|$. Such a solution describes a suspension, i.e., the motion of a small particle in the fluid, by which one can calculate an effective viscosity being different from that of the original fluid and determine the radius of particles. Einstein [12] calculated their quantities when the obstacle is a sphere (see Batchelor [3] and Landau-Lifschitz [23]).

In 1850, G.G. Stokes showed that, in general, in two-dimensional exterior domains, there is no solution $u$ of $(S)$ tending to a prescribed non-zero constant vector at infinity. We shall first generalize the "Stokes paradox" to higher dimensions and determine the exact class of solutions in which the paradox holds. Indeed, we shall treat the simpler class $\left(D_{q}\right)$ and show that $u \equiv 0$ is the only solution of $(S)$ with $f \equiv g \equiv 0$, if $1<q \leq n /(n-1)$. In the two-dimensional case, Finn [14] and Heywood [20] obtained similar results. Secondly, we shall give a concrete characterization of the null-space for the solutions of ( $S$ ) in the class $\left(C L_{q}\right)$. Here we shall see that a non-trivial null-space appears when $q$ varies and that the case $q=n /(n-1)$ is critical. Finally, based on these results of the null-space, we shall give a theorem on the existence and uniqueness for the solutions of $(S)$ in the class $\left(C L_{q}\right)$. This theorem holds if one can solve $(S)$ with the boundary condition (B.C. $)_{\infty}$ at infinity.

Our basic tool consists of the two fundamental facts, a regularity theory (Theorem 3.1) and an a priori estimate (Theorem 3.3) in $L^{q}$-spaces for the gradient of solutions of $(S)$. The regularity theorem is useful to show the Stokes paradox in higer dimensions and enables us to see why the critical value $q=n /(n-1)$ appears in solvability of $(S)$. The a priori estimate plays a basic role in characterization of the null-space and range of solutions. Such an estimate has been got by several authors for $n \geq 3$ (Kozono-Sohr [22], Borchers-Miyakawa [7]). Recently, Galdi-Simader [18] obtained a similar result to ours by using the hydrodynamic potentials. Our method is however different from Galdi-Simader's [18]: we are based on the cut-off procedure. Making use of a simple embedding argument about a certain functional space, we shall show the same a priori estimate holding for all $n \geq 2$.

Concerning characterization of the null-space in the class $\left(C L_{q}\right)$, Maslennikova-Timoshin [25] solved ( $S$ ) explicitly in an exterior domain of the unit sphere in $\mathbb{R}^{3}$ and announced a similar result to ours. They have used the special functions (the Legendre functions) for representation of the solution. We shall give a more systematic treatement for generalized solutions of ( $S$ ). Our approach is so different from [25] that we can apply it to all dimensions $n \geq 2$. For another investigation such as strong solutions, see, e.g., Sohr-Varnhorn [30].

## 1. - Main Results

1.1. Before stating our results we introduce some notations. For $1<q<\infty$ $\left(q^{\prime}=q /(q-1)\right),\|\cdot\|_{q}$ and $(\cdot, \cdot)$ denote the usual norm of $L^{q}(\Omega)$ and the inner product between $L^{q}(\Omega)$ and $L^{q^{\prime}}(\Omega)$, respectively. In general we shall denote by $(f, \phi)$ the value of the distribution $f$ at $\phi \in C_{0}^{\infty}(\Omega) . \hat{H}_{0}^{1, q}(\Omega)$ is the completion of $C_{0}^{\infty}(\Omega)$ with respect to the norm $\|\nabla u\|_{q}$. Since $\Omega$ is an exterior domain, $\hat{H}_{0}^{1, q}(\Omega)$ is larger than $H_{0}^{1, q}(\Omega)$. Having introduced $\hat{H}_{0}^{1, q}(\Omega)$, it is also useful to define $\hat{H}^{-1, q}(\Omega):=\hat{H}_{0}^{1, q^{\prime}}(\Omega)^{*}\left(X^{*}\right.$; dual space of $\left.X\right)$, and $\|\cdot\|_{-1, q}$ denotes the norm of $\hat{H}^{-1, q}(\Omega)$ defined by $\|f\|_{-1, q}:=\sup \left\{|(f, \phi)| /\|\nabla \phi\|_{q^{\prime}} ; \phi \in C_{0}^{\infty}(\Omega), \phi \neq 0\right\}$. We shall denote by $C_{0}^{\infty}(\Omega)^{n}, L^{q}(\Omega)^{n}, \cdots$, and $C_{0}^{\infty}(\Omega)^{n^{2}}, L^{q}(\Omega)^{n^{2}}, \cdots$ the corresponding spaces for the vector-valued and the matrix-valued functions, respectively. In such spaces, we shall also use the same notations $\|\cdot\|_{q}$ and $(\cdot, \cdot)$.

Let $f \in \hat{H}^{-1, q}(\Omega)^{n}$ and $g \in L_{\mathrm{loc}}^{q}(\bar{\Omega})$, where $g \in L_{\mathrm{loc}}^{q}(\bar{\Omega})$ means that $\int_{\Omega \cap B}|g(x)|^{q} \mathrm{~d} x<\infty$ for all open balls $B$ in $\mathbb{R}^{n}$ with $\Omega \cap B \neq \phi$. A pair $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ with $\left.u\right|_{\partial \Omega}=0$ (in the trace sense) is called a generalized solution of ( $S$ ) if

$$
(\nabla u, \nabla \Phi)-(p, \operatorname{div} \Phi)=(f, \Phi),-(u, \nabla \phi)=(g, \phi)
$$

for all $\Phi \in C_{0}^{\infty}(\Omega)^{n}$ and all $\phi \in C_{0}^{\infty}(\Omega)$, respectively.
1.2. Our result on the generalized Stokes paradox now reads:

Theorem A. (Stokes paradox). Let $n \geq 2$ and $1<q \leq n^{\prime}\left(n^{\prime} \equiv n /(n-1)\right.$ ). Suppose that $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ is a generalized solution of $(S)$ with $f \equiv 0, g \equiv 0$ satisfying $\nabla u \in L^{q}(\Omega)^{n^{2}}$. Then it follows that $u \equiv 0, p \equiv 0$ in $\Omega$.

By Bogovskii's result [6], the pressure $p$ is determined by $u$, and hence we can restate the above theorem without $p$.

THEOREM A'. Let $n \geq 2$ and $1<q \leq n^{\prime}$. Suppose that $u \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n}$ satisfies $\operatorname{div} u=0$ in $\Omega,\left.u\right|_{\partial \Omega}=0$, and $(\nabla u, \nabla \Phi)=0$ for all $\Phi \in C_{0}^{\infty}(\Omega)^{n}$ with $\operatorname{div} \Phi=0$. If, in addition, $\nabla u \in L^{q}(\Omega)^{n^{2}}$, then we have $u \equiv 0$ in $\Omega$.

REMARKS. 1. In the above theorem, we do not assume any integrability condition on $u$ itself. It follows that there is no solution $u$ of $(S)$ with $f \equiv 0, g \equiv 0$ in the class $\left(D_{q}\right)$ for $1<q \leq n^{\prime}$ such that $u(x) \rightarrow a$ as $x \rightarrow \infty$, where $a$ is a non-zero constant vector in $\mathbb{R}^{n}$.
2. Heywood [20] showed the same result in the special case $n=q=2$. Chang-Finn [11] gave a similar result for $n=2$ in the class $u(x)=\mathrm{o}(\log |x|)$ as $x \rightarrow \infty$.
1.3. We next proceed to the characterization of the null-space for $(S)$ in the class $\left(C L_{q}\right)$.

Let us denote by $\mathbb{N}_{q}$ the set of all generalized solutions $\{u, p\} \in$ $H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ of $(S)$ with $f \equiv 0, g \equiv 0$ satisfying $\nabla u-A \in L^{q}(\Omega)^{n^{2}}$ for some matrix $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0 . \mathbb{N}_{q}^{0}$ is the subspace of $\mathbb{N}_{q}$ defined by $\mathbb{N}_{q}^{0} \equiv\left\{\{u, p\} \in \mathbb{N}_{q} ; \nabla u \in L^{q}(\Omega)^{n^{2}}\right\}$.

Our second result now reads:
THEOREM B. (Characterization of the null-space). (i) Let $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Then $\operatorname{dim} \mathbb{N}_{q}=n^{2}-1$ and $\operatorname{dim} \mathbb{N}_{q}^{0}=0$. For every $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$ and $a \in \mathbb{R}^{n}$ satisfying the condition

$$
\begin{equation*}
\int_{\partial \Omega}\left\{(A x+a) \cdot \frac{\partial v}{\partial \nu}-\chi(A x+a) \cdot \nu\right\} \mathrm{d} S=0 \tag{1.1}
\end{equation*}
$$

for all $\{v, \chi\} \in \mathbb{N}_{q}^{0}$, there exists a unique $\{u, p\} \in \mathbb{N}_{q}$ such that

$$
\begin{align*}
& \nabla u-A \in L^{q}(\Omega)^{n^{2}}  \tag{1.2}\\
& u \in C^{0}(\bar{\Omega})^{n}, \lim _{x \rightarrow \infty}|u(x)-(A x+a)|=0 \tag{1.3}
\end{align*}
$$

where $\nu$ denotes the unit outer normal to $\partial \Omega$ and $\mathrm{d} S$ is the surface element of $\partial \Omega$. Conversely, for every $\{u, p\} \in \mathbb{N}_{q}$, there are unique $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$ and $a \in \mathbb{R}^{n}$ such that (1.2) and (1.3) hold.
(ii) Let $n^{\prime}<q<\infty, n \geq 2$. Then $\operatorname{dim} \mathbb{N}_{q}=n^{2}+n-1$ and $\operatorname{dim} \mathbb{N}_{q}^{0}=n$. For every $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$ and $a \in \mathbb{R}^{n}$, there exists a unique $\{u, p\} \in \mathbb{N}_{q}$ such that (1.2) and (1.3) hold if $n \geq 3$, and such that (1.2) and

$$
\int_{\Omega}|\nabla[u(x)-A x-E(x) a]|^{2} \mathrm{~d} x<\infty
$$

hold if $n=2$, where $E=\left(E_{i j}(x)\right)_{i, j=1,2}$ denotes the fundamental tensor of the Stokes equations in $\mathbb{R}^{2}: E_{i j}(x)=(4 \pi)^{-1}\left[\log \left(\frac{\delta_{i j}}{|x|}\right)+\frac{x_{i} x_{j}}{|x|^{2}}\right]$. Conversely, for every $\{u, p\} \in \mathbb{N}_{q}$, there are unique $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$ and $a \in \mathbb{R}^{n}$ such that (1.2)-(1.3) hold if $n \geq 3$, and such that (1.2)-(1.3') hold if $n=2$.
(iii) Let $n=q=2$. Then $\operatorname{dim} \mathbb{N}_{2}=n^{2}-1=3$ and $\operatorname{dim} \mathbb{N}_{2}^{0}=0$. For every $A \in \mathbb{R}^{2^{2}}$ with $\operatorname{Tr} A=0$, there is a unique $\{u, p\} \in \mathbb{N}_{2}$ such that (1.2) holds with $q=2$. Conversely, for every $\{u, p\} \in \mathbb{N}_{2}$, there is a unique $A \in \mathbb{R}^{2^{2}}$ with $\operatorname{Tr} A=0$ such that (1.2) holds with $q=2$.

REMARK. For $\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}$, we have $\nabla v-\chi I \in L^{q^{\prime}}(\Omega)^{n^{2}}$ and $\operatorname{div}(\nabla v-\chi I)=$ 0 in the sense of distributions in $\Omega$, where $I$ is the identity matrix in $\mathbb{R}^{n^{2}}$. Using the trace theorem as in Miyakawa [26, Proposition 1.2] and Simader-Sohr [29], we see that $\frac{\partial v}{\partial \nu}-\chi I \cdot \nu \in\left(H^{1 / q^{\prime}, q}(\partial \Omega)^{n}\right)^{*}$ and hence (1.1) should be understood in such a generalized sense as the duality between $\frac{\partial v}{\partial \nu}-\chi I \cdot \nu \in$ $\left(H^{1 / q^{\prime}, q}(\partial \Omega)^{n}\right)^{*}$ and $A x+a \in H^{1 / q^{\prime}, q}(\partial \Omega)^{n}$. However, from the regularity theorem in bounded domains (as Cattabriga [10] showes), we get $v \in H_{\mathrm{loc}}^{2, q^{\prime}}(\bar{\Omega})^{n}, \chi \in$ $H_{\text {loc }}^{2, q^{\prime}}(\bar{\Omega})$; therefore (1.1) may be also regarded in the usual sense.
1.4. We are next concerned with the necessary and sufficient condition for the solvability of $(S)$ in the class $\left(C L_{q}\right)$.

THEOREM C. (Inhomogeneous case). (i) Let $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Then for every $f \in \hat{H}^{-1, q}(\Omega)^{n}, g \in L_{\text {loc }}^{q}(\bar{\Omega}), A \in \mathbb{R}^{n^{2}}$ with $g-\operatorname{Tr} A \in L^{q}(\Omega)$ and $a \in \mathbb{R}^{n}$, there exists a generalized solution $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\overline{\mathbf{\Omega}})^{n} \times L^{q}(\Omega)$ of (S) satisfying (1.2) and

$$
\begin{equation*}
\int_{\Omega}|u(x)-(A x+a)|^{n q /(n-q)} \mathrm{d} x<\infty \tag{1.4}
\end{equation*}
$$

if and only if the compatibility condition

$$
\begin{equation*}
(f, v)-(g-\operatorname{Tr} A, \chi)+\int_{\partial \Omega}\left\{(A x+a) \cdot \frac{\partial v}{\partial \nu}-\chi(A x+a) \cdot \nu\right\} \mathrm{d} S=0 \tag{1.5}
\end{equation*}
$$

holds for all $\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}$. Such $\{u, p\}$ is unique and subject to the inequality

$$
\begin{equation*}
\|\nabla u-A\|_{q}+\|p\|_{q} \leq C\left(\|f\|_{-1, q}+\|g-\operatorname{Tr} A\|_{q}+|A|+|a|\right) \tag{1.6}
\end{equation*}
$$

with $C=C(\Omega, n, q)>0$ independent of $u$ and $p$, where $|A|$ and $|a|$ denote the standard Euclidian norms in $\mathbb{R}^{n^{2}}$ and $\mathbb{R}^{n}$, respectively.
(ii) Let $n^{\prime}<q<n, n \geq 3$. Then for every $f \in \hat{H}^{-1, q}(\Omega)^{n}, g \in L_{\mathrm{loc}}^{q}(\bar{\Omega})$, $A \in \mathbb{R}^{n^{2}}$ with $g-\operatorname{Tr} A \in L^{q}(\Omega)$ and $a \in \mathbb{R}^{n}$, there exists a unique generalized solution $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ of $(S)$ such that (1.2) and (1.4) hold. Such $\{u, p\}$ is subject to the inequality (1.6). If in addition $f \in \hat{H}^{-1, r}(\Omega)^{n}, g-\operatorname{Tr} A \in$ $L^{r}(\Omega)$ for some $r>n$, we have also $\nabla u-A \in L^{r}(\Omega)^{n^{2}}, p \in L^{r}(\Omega)$ and (1.3).
(iii) Let $n \leq q<\infty$ for $n \geq 3$ and $2<q<\infty$ for $n=2$. Then for every $f \in \hat{H}^{-1, q}(\Omega)^{n}, g \in L_{\mathrm{loc}}^{q}(\bar{\Omega})$ and $\bar{A} \in \mathbb{R}^{n^{2}}$ with $g-\operatorname{Tr} A \in L^{q}(\Omega)$, there exists at least one generalized solution $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ of $(S)$ satisfying (1.2). Such $\{u, p\}$ is unique modulo $\mathbb{N}_{q}^{0}$ and subject to the inequality

$$
\begin{align*}
& \inf \left\{\|\nabla u-A-\nabla v\|_{q}+\|p-\chi\|_{q} ;\{v, \chi\} \in \mathbb{N}_{q}^{0}\right\}  \tag{1.7}\\
& \leq C\left(\|f\|_{-1, q}+\|g-\operatorname{Tr} A\|_{q}+|A|\right)
\end{align*}
$$

where $C=C(\Omega, n, q)>0$.
(iv) Let $n=q=2$. Then for every $f \in \hat{H}^{-1,2}(\Omega)^{2}, g \in L_{\text {loc }}^{2}(\bar{\Omega})$ and $A \in \mathbb{R}^{2^{2}}$ with $g-\operatorname{Tr} A \in L^{2}(\Omega)$, there exists a unique generalized solution $\left.\{u, p\} \in H_{\mathrm{loc}}^{1,2} \overline{\boldsymbol{\Omega}}\right)^{2} \times L^{2}(\Omega)$ of $(S)$ satisfying (1.2) with $q=2$. Such $\{u, p\}$ is subject to the inequality

$$
\begin{equation*}
\|\nabla u-A\|_{2}+\|p\|_{2} \leq C\left(\|f\|_{-1,2}+\|g-\operatorname{Tr} A\|_{2}+|A|\right) \tag{1.8}
\end{equation*}
$$

where $C=C(\Omega)>0$.
REMARKS. 1. In case (i), the compatibility condition (1.5) is necessary and sufficient for the solvability of $(S)$.
2. In case (ii), the additional condition $f \in \hat{H}^{-1, r}(\Omega)^{n}, g-\operatorname{Tr} A \in$ $L^{r}(\Omega)(r>n)$ enables us to get the smoothness of $u$ and its asymptotic behaviour (1.3). In case (iii) we cannot prescribe $a \in \mathbb{R}^{n}$ so that the uniqueness follows. However, if we assume in addition that $f \in \hat{H}^{-1, \gamma}(\Omega)^{n}, g-\operatorname{Tr} A \in L^{\gamma}(\Omega)$ for some $n^{\prime}<\gamma<n$, then we can prescribe $a \in \mathbb{R}^{n}$ so as to get the unique solvability under the condition (1.3).

## 2. - Preliminaries

2.1. Homogeneous Sobolev space $\hat{H}_{0}^{1, q}(\Omega)$.

In this subsection we shall give a concrete characterization of $\hat{H}_{0}^{1, q}(\Omega)$ and some elementary lemmas for the proof of the main results.

Let $D$ be a domain in $\mathbb{R}^{n}(n \geq 2)$. We denote by $\|\cdot\|_{q, D}$ and $(\cdot, \cdot)_{D}$ the norm of $L^{q}(D)$ and the inner product between $L^{q}(D)$ and $L^{q^{\prime}}(D)$, respectively. $\hat{H}_{0}^{1, q}(D)$ is the completion of $C_{0}^{\infty}(D)$ with respect to the norm $\|\nabla u\|_{q, D}$. If $D$ is bounded, the Poincare inequality states that $\hat{H}_{0}^{1, q}(D)=H_{0}^{1, q}(D)$, but, in general, $\hat{H}_{0}^{1, q}(D)$ is larger than $H_{0}^{1, q}(D) . \hat{H}^{-1, q}(D)$ is the dual space of $\hat{H}_{0}^{1, q^{\prime}}(D)\left(1 / q+1 / q^{\prime}=1\right)$ whose norm is denoted by $\|\cdot\|_{-1, q, D}$. In case $D=\Omega$, we shall call these norms $\|\cdot\|_{q},(\cdot, \cdot)$, and $\|\cdot\|_{-1, q}$. In what follows $C$ denotes a constant which may change from line to line. In particular, $C=C(*, \cdots, *)$ denotes a constant depending only on the quantities appearing in the parentheses.

The following inequality is simple but very useful for the forthcoming arguments (see also Simader-Sohr [29]).

Variational inequality in $L^{q}$. Let $n \geq 2$ and $1<q<\infty$. Suppose that $u \in L_{\text {loc }}^{q}\left(\mathbb{R}^{n}\right)$ with $\nabla u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$. Then we have

$$
\|\nabla u\|_{q, \mathbb{R}^{n}}
$$

$$
\begin{equation*}
\leq C \sup \left\{\frac{\left|(\nabla u, \nabla \phi)_{\mathbb{R}^{n}}\right|}{\|\nabla \phi\|_{q^{\prime}, \mathbb{R}^{n}}} ; \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), \phi \neq 0\right\} \tag{2.1}
\end{equation*}
$$

with $C=C(n, q)$ independent of $u$.
Indeed, the Calderon-Zygmund inequality gives

$$
\|\nabla \nabla \psi\|_{q^{\prime}, \mathbb{R}^{n}} \leq C\|\Delta \psi\|_{q^{\prime}, \mathbb{R}^{n}}\left(\psi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)\right)
$$

Then, since the space $H \equiv\left\{\Delta \psi ; \psi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)\right\}$ is dense in $L^{q^{\prime}}\left(\mathbb{R}^{n}\right)$, we have for each $i=1, \ldots, n$

$$
\begin{aligned}
& \sup \left\{\frac{\left|(\nabla u, \nabla \phi)_{\mathbb{R}^{n}}\right|}{\|\nabla \phi\|_{q^{\prime}, \mathbb{R}^{n}}} ; \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), \phi \neq 0\right\} \\
& \geq \sup \left\{\frac{\left|\left(\nabla u, \nabla\left(\partial_{i} \psi\right)\right)_{\mathbb{R}^{n}}\right|}{\left\|\nabla\left(\partial_{i} \psi\right)\right\|_{q^{\prime}, \mathbb{R}^{n}}} ; \psi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), \psi \neq 0\right\} \\
& \geq C \sup \left\{\frac{\left|\left(\partial_{i} u, \Delta \psi\right)_{\mathbb{R}^{n}}\right|}{\|\Delta \psi\|_{q^{\prime}, \mathbb{R}^{n}}} ; \psi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), \psi \neq 0\right\} \\
& =C \sup \left\{\frac{\left|\left(\partial_{i} u, g\right)_{\mathbb{R}^{n}}\right|}{\|g\|_{q^{\prime}, \mathbb{R}^{n}}} ; g \in L^{q^{\prime}}\left(\mathbb{R}^{n}\right), g \neq 0\right\} \\
& =C\left\|\partial_{i} u\right\|_{q, \mathbb{R}^{n}}
\end{aligned}
$$

with $C=C(n, q)$, and (2.1) follows.
Based on the above variational inequality, we get the following approximation lemma.

LEMMA 2.1. Let $n \geq 2$ and $1<q<\infty$. Then for every $u \in L_{\text {loc }}^{q}(\bar{\Omega})$ with $\nabla u \in L^{q}(\Omega)^{n}$, there is a sequence $\left\{u_{j}\right\}_{j=1}^{\infty}$ in $C_{0}^{\infty}(\bar{\Omega})$ such that $\nabla u_{j} \rightarrow \nabla u$ in $L^{q}(\Omega)^{n}$, where $C_{0}^{\infty}(\bar{\Omega})$ is the set of all $C^{\infty}$-functions $\phi$ with compact support in $\bar{\Omega}$ ( $\phi$ may not vanish on $\partial \Omega$ ). The same assertion is true with $\Omega$ replaced by $\mathbb{R}^{n}$.

Proof. By the extension theorem (Adams [1]), for each $u \in L_{\mathrm{loc}}^{q}(\overline{\mathbf{\Omega}})$ with $\nabla u \in L^{q}(\Omega)^{n}$, there is a function $\tilde{u} \in L_{\mathrm{loc}}^{q}\left(\mathbb{R}^{n}\right)$ with $\nabla \tilde{u} \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$ such that $\tilde{u}=u$ in $\Omega$, so we may only prove the assertion on $\mathbb{R}^{n}$. Let $L^{1, q}=\left\{u \in L_{\text {loc }}^{q}\left(\mathbb{R}^{n}\right) ; \nabla u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}\right\}$. We denote by $[u]$ the set of all $v \in L^{1, q}$ such that $u-v$ is a constant function on $\mathbb{R}^{n}$, and set $L^{1, q} / \mathbb{R} \equiv\left\{[u] ; u \in L^{1, q}\right\}$ and $G_{q} \equiv\left\{\nabla u \in L^{q}\left(\mathbb{R}^{n}\right)^{n} ;[u] \in L^{1, q} / \mathbb{R}\right\}$. We may regard $G_{q}$ as a closed
subspace of $L^{q}\left(\mathbb{R}^{n}\right)^{n}$; equipped with the norm $\|[u]\|_{L^{1, q / \mathbb{R}}}:=\|\nabla u\|_{q, \mathbb{R}^{n}}, L^{1, q} / \mathbb{R}$ is a Banach space isometric to $G_{q}$. Hence it suffices to prove that the space $W \equiv\left\{\nabla \phi ; \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)\right\}$ is dense in $G_{q}$. To this end, let us consider a map $A_{q}: \nabla u \in G_{q} \rightarrow A_{q}(\nabla u) \in G_{q^{\prime}}^{*}$ defined by $\left\langle A_{q}(\nabla u), \nabla v\right\rangle=(\nabla u, \nabla v)_{\mathbb{R}^{n}}$ for $\nabla v \in G_{q^{\prime}}$, where $\langle\cdot, \cdot\rangle$ denotes the duality between $G_{q^{\prime}}^{*}$ and $G_{q^{\prime}}$. Then by (2.1) we see that $A_{q}$ is injective and that its range is closed in $G_{q^{\prime}}^{*}$. Since $A_{q}^{*}$ coincides with $A_{q^{\prime}}\left(T^{*}\right.$; adjoint operator of $\left.T\right)$, it follows from the closed range theorem that $A_{q}$ is also surjective and hence bijective. Now, suppose that $F \in G_{q}^{*}$ satisfies $\langle F, \nabla \phi\rangle=0$ for all $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$. Since $A_{q^{\prime}}$ is also bijective, there is a unique $\nabla u \in G_{q^{\prime}}$ such that $\langle F, \nabla u\rangle=\left\langle A_{q^{\prime}}(\nabla u), \nabla v\right\rangle=(\nabla u, \nabla v)_{\mathbb{R}^{n}}$ holds for all $\nabla v \in G_{q}$. Then by the assumption and (2.1) we get $\nabla u=0$ and hence $F=0$, which implies that $W$ is dense in $G_{q}$.

REMARK. Simader [28] gave another proof of this lemma by using the Poincaré inequality on annular domains and a scaling argument.

The following concrete characterization of the space $\hat{H}_{0}^{1, q}(\Omega)$ is essentially due to Galdi-Simader [18, Theorem 1.1]. Based on Lemma 2.1, we give here another proof.

LEMMA 2.2 (Galdi-Simader). (i) For $1<q<n$, we have

$$
\hat{H}_{0}^{1, q}(\Omega)=\left\{u \in L^{n q /(n-q)}(\Omega) ; \quad \nabla u \in L^{q}(\Omega)^{n},\left.u\right|_{\partial \Omega}=0\right\} .
$$

(ii) For $n \leq q<\infty$, we have

$$
\hat{H}_{0}^{1, q}(\Omega)=\left\{u \in L_{\mathrm{loc}}^{q}(\bar{\Omega}) ; \quad \nabla u \in L^{q}(\Omega)^{n},\left.u\right|_{\partial \Omega}=0\right\}
$$

If $n<q$, the function $u \in \hat{H}_{0}^{1, q}(\Omega)$ is continuous on $\bar{\Omega}$ (after redefinition on a set of measure zero of $\Omega$ ) and satisfies

$$
u(x)=\mathrm{O}\left(|x|^{1-n / q}\right) \text { as } x \rightarrow \infty
$$

Proof. Let $H_{q}$ be the space defined by the right-hand side of (i) and (ii). By the Sobolev inequality, it is easy to see that $\hat{H}_{0}^{1, q}(\Omega) \subset H_{q}$ and so we may only prove the converse inclusion. To this end, we introduce an extension operator $\Gamma$. Take $R>0$ so that $\partial \Omega \subset B_{R} \equiv\left\{x \in \mathbb{R}^{n} ;|x|<R\right\}$ and consider a continuous extension operator $\Gamma: H^{1-1 / q, q}(\partial \Omega) \rightarrow H^{1, q}(\Omega)$ satisfying $\operatorname{supp} \Gamma \phi \subset B_{R}$ for all $\phi \in H^{1-1 / q, q}(\partial \Omega)$.
(i) Case $1<q<n$. Let $u \in H_{q}$. Then by Lemma 2.1 , there is a sequence $\left\{u_{j}\right\}_{j=1}^{\infty}$ in $C_{0}^{\infty}(\bar{\Omega})$ such that $\nabla u_{j} \rightarrow \nabla u$ in $L^{q}(\Omega)^{n}$. Since $u \in L^{n q /(n-q)}(\Omega)$, it follows from the Sobolev inequality that $u_{j} \rightarrow u$ in $L^{n q /(n-q)}(\Omega)$. Then by the trace theorem, we get $\left.u_{j}\right|_{\partial \Omega} \rightarrow 0$ in $H^{1-1 / q, q}(\partial \Omega)$. Setting $w_{j}=u_{j}-\Gamma\left(\left.u_{j}\right|_{\partial \Omega}\right)$, we get $w_{j} \in H_{0}^{1, q}(\Omega)$ and it follows from the continuity of $\Gamma$ that

$$
\left\|\nabla w_{j}-\nabla u\right\|_{q} \leq\left\|\nabla u_{j}-\nabla u\right\|_{q}+C\left\|\left.u_{j}\right|_{\partial \Omega}\right\|_{H^{1-1 / q, q}(\partial \Omega)}
$$

with $C$ independent of $j$. Hence $\nabla w_{j} \rightarrow \nabla u$ in $L^{q}(\Omega)^{n}$, and since $C_{0}^{\infty}(\Omega)$ is dense in $H_{0}^{1, q}(\Omega)$, we obtain $u \in \hat{H}_{0}^{1, q}(\Omega)$.
(ii) Case $n \leq q<\infty$. Let $u \in H_{q}$. Then it follows from Lemma 2.1 and a standard argument that there are sequences $\left\{u_{j}\right\}_{j=1}^{\infty}$ in $C_{0}^{\infty}(\bar{\Omega})$ and $\left\{c_{j}\right\}_{j=1}^{\infty}$ in $\mathbb{R}$ such that $\nabla u_{j} \rightarrow \nabla u$ in $L^{q}(\Omega)^{n}$ and $u_{j}+c_{j} \rightarrow u$ in $L_{\text {loc }}^{q}(\bar{\Omega})$. We shall next approximate the sequence $\left\{c_{j}\right\}_{j=1}^{\infty}$ in terms of a sequence of functions in $C_{0}^{\infty}(\Omega)$ with respect to the norm $\|\nabla \cdot\|_{q}$. Take $\zeta \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ satisfying $0 \leq \varsigma \leq 1, \zeta(x)=1$ for $|x| \leq 1$ and $\zeta(x)=0$ for $|x| \geq 2$ and set $\zeta_{k}(x)=\zeta(x / k)(k=1,2, \cdots)$. The sequence $\left\{s_{k}\right\}_{k=1}^{\infty}$ will be called a sequence of $n$-dimensional cut-off functions. Then we have $\zeta_{k}(x)=1$ for $|x| \leq k$ and $\left\|\nabla_{\zeta_{k}}\right\|_{q, \mathbb{R}^{n}} \leq C k^{-1+n / q}(k=1,2, \cdots)$ with $C$ independent of $k$. Since $n \leq q$, by Mazur's theorem ([33, p. 120 Theorem 2), we can choose a sequence $\left\{\bar{\zeta}_{k}\right\}_{k=1}^{\infty}$ of convex combinations of $\zeta_{k}^{\prime} s$ so that

$$
\nabla \bar{\zeta}_{k} \rightarrow 0 \text { in } L^{q}\left(\mathbb{R}^{n}\right)^{n}, \bar{\zeta}_{k} \rightarrow 1 \text { locally uniformly in } \mathbb{R}^{n} .
$$

Hence there is a subsequence $\left\{\bar{\zeta}_{k(j)}\right\}_{j=1}^{\infty}$ of $\left\{\bar{\zeta}_{k}\right\}_{k=1}^{\infty}$ such that $c_{j}\left\|\nabla \bar{\zeta}_{k(j)}\right\|_{q, \mathbb{R}^{n}} \rightarrow 0$ as $j \rightarrow \infty$. Defining $\bar{u}_{j}=u_{j}+c_{j} \bar{S}_{k(j)}(j=1,2, \cdots)$, we have $\bar{u}_{j} \in C_{0}^{\infty}(\bar{\Omega})$ and $\nabla \bar{u}_{j} \rightarrow \nabla u$ in $L^{q}(\Omega)^{n}, \bar{u}_{j} \rightarrow u$ in $L_{\text {loc }}^{q}(\bar{\Omega})$. Now, making use of a sequence $w_{j}=\bar{u}_{j}-\Gamma\left(\left.\bar{u}_{j}\right|_{\partial \Omega}\right)(j=1,2, \cdots)$ as in the case of (i), we can prove similarly as above that $u \in \hat{H}_{0}^{1, q}(\Omega)$.

Finally, the asymptotic behaviour $u(x)=\mathrm{O}\left(|x|^{1-n / q}\right), x \rightarrow \infty$ for $u \in \hat{H}_{0}^{1, q}(\Omega)$ with $q>n$ follows from Friedman [15, p. 23 Theorem 9.2].

We shall next consider the complex interpolation space $[X, Y]_{\theta}(0 \leq \theta \leq 1)$. For all $1<q, r<\infty$, the norms $\|\nabla u\|_{q}$ and $\|\nabla u\|_{r}$ are consistent on $C_{0}^{\infty}(\Omega)$, so the pair $\left\{\hat{H}_{0}^{1, q}(\Omega), \hat{H}_{0}^{1, r}(\Omega)\right\}$ is interpolation couple. See Reed-Simon [27, p. 35]. Using the Riesz-Thorin theorem [32, 1.18.7], we obtain from Lemma 2.2 the following result:

If $1<q<n, 1<r<n$ and if $n \leq q<\infty, n \leq r<\infty$, then

$$
\begin{equation*}
\left[\hat{H}_{0}^{1, q}(\Omega), \hat{H}_{0}^{1, r}(\Omega)\right]_{\theta}=\hat{H}_{0}^{1, s}(\Omega) \tag{2.2}
\end{equation*}
$$

where $1 / s=(1-\theta) / q+\theta / r, 0 \leq \theta \leq 1$.
In the whole space $\mathbb{R}^{n}$, we shall prove the corresponding result without restriction on $q$ and $r$.

LEMMA 2.3. Let $n \geq 2$ and $1<q<\infty, 1<r<\infty$. Then we have

$$
\left[\hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right), \quad \hat{H}_{0}^{1, r}\left(\mathbb{R}^{n}\right)\right]_{\theta}=\hat{H}_{0}^{1, s}\left(\mathbb{R}^{n}\right)
$$

where $1 / s=(1-\theta) / q+\theta / r, 0 \leq \theta \leq 1$.
Proof. Let $E_{q} \equiv\left\{\nabla u \in L^{q}\left(\mathbb{R}^{n}\right)^{n} ; u \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)\right\}$. Then we may regard $E_{q}$ as a closed subspace of $L^{q}\left(\mathbb{R}^{n}\right)^{n}$. Hence $E_{q}$ is a Banach space with the norm
$\|\nabla u\|_{E_{q}}:=\|\nabla u\|_{q, \mathbb{R}^{n}}$ for $\nabla u \in E_{q}$, and isometric to $\hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)$. Now it suffices to show that

$$
\begin{equation*}
\left[E_{q}, E_{r}\right]_{\theta}=E_{s} \text { for } q, r, s \text { and } \theta \text { as above. } \tag{2.3}
\end{equation*}
$$

To this end, we need to solve the equation $\Delta \chi=\operatorname{div} u$ in $\mathbb{R}^{n}$ in the following weak sense:

For every $u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$, there is a unique $\chi \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
(\nabla \chi, \nabla \phi)_{\mathbb{R}^{n}}=(u, \nabla \phi)_{\mathbb{R}^{n}} \text { for all } \phi \in \hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right) \tag{2.4}
\end{equation*}
$$

Based on (2.1), we see as in the proof of Lemma 2.1 that the map $B_{q}: \nabla u \in E_{q} \rightarrow B_{q}(\nabla u) \in E_{q^{\prime}}^{*}$ defined by $\left\langle B_{q}(\nabla u), \nabla v\right\rangle:=(\nabla u, \nabla v)_{\mathbb{R}^{n}}$ for $\nabla v \in E_{q^{\prime}}$ is a bijective operator. Here $\langle\cdot, \cdot\rangle$ denotes the duality between $E_{q^{\prime}}^{*}$ and $E_{q^{\prime}}$. Since the map $\nabla \phi \in E_{q^{\prime}} \rightarrow(u, \nabla \phi)_{\mathbb{R}^{n}} \in \mathbb{R}$ is a continuous functional on $E_{q^{\prime}}$, we can solve (2.4) uniquely for every given $u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$.

Now, it is easy to see that the map $Q: u \rightarrow \nabla \chi$ defined by the relation (2.4) is a projection operator from $L^{q}\left(\mathbb{R}^{n}\right)^{n}$ onto $E_{q}$. Then (2.3) follows from Bergh-Löfström [4, Theorem 6.4.2].

We need further the following two lemmas.
LEMMA 2.4. Let $1<q<\infty$ and $h \in L^{q}\left(\mathbb{R}^{n}\right)$. If

$$
\sup \left\{\frac{\left|(h, \Delta \phi)_{\mathbb{R}^{n}}\right|}{\|\Delta \phi\|_{r^{\prime}, \mathbb{R}^{n}}} ; \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right), \phi \neq 0\right\}<\infty
$$

for some $1<r<\infty$, then we have also $h \in L^{r}\left(\mathbb{R}^{n}\right)$.
Proof. Here we follow Simader-Sohr [29]. Since the space $H \equiv\{\Delta \phi ; \phi \in$ $\left.C^{\infty}\left(\mathbb{R}^{n}\right)\right\}$ is a dense subspace in $L^{r^{\prime}}\left(\mathbb{R}^{n}\right)$, we see by the assumption that the map $\Delta \phi \in H \rightarrow(h, \Delta \phi)_{\mathbb{R}^{n}} \in \mathbb{R}$ is uniquely extended as a continuous functional on $L^{r^{\prime}}\left(\mathbb{R}^{n}\right)$. Hence there is a unique $\eta \in L^{r}\left(\mathbb{R}^{n}\right)$ such that $(\eta, \Delta \phi)_{\mathbb{R}^{n}}=(h, \Delta \phi)_{\mathbb{R}^{n}}$ holds for all $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$. Since $w:=h-\eta \in L_{\text {loc }}^{1}\left(\mathbb{R}^{n}\right)$, Weyl's lemma states that the function $w$ is of class $C^{\infty}$ and harmonic in $\mathbb{R}^{n}$ in the classical sense. Applying the mean value property to $w$ on the ball $B_{|x|}(x)$ centered at $x \neq 0$ with radius $|x|$, and then using the Hölder inequality, we obtain the estimate

$$
|w(x)| \leq C\left(\|h\|_{q, \mathbb{R}^{n}}|x|^{-n / q}+\|\eta\|_{r, \mathbb{R}^{n}}|x|^{-n / r}\right)
$$

where $C=C(n, q, r)$. Then it follows from the Liouville theorem that $w \equiv 0$ in $\mathbb{R}^{n}$ and hence $h \in L^{r}\left(\mathbb{R}^{n}\right)$.

LEMMA 2.5 (Embedding argument). Let $\Omega_{0}$ be a subdomain of $\Omega$ with closure $\bar{\Omega}_{0}$ contained in $\Omega$. Then for each $1<q<\infty$, there is a constant $C=C\left(\Omega, \Omega_{0}, n, q\right)$ such that

$$
\|f\|_{-1, q, \mathbb{R}^{n}} \leq C\|f\|_{-1, q, \Omega}
$$

holds for all $f \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)$ with supp $f \subset \bar{\Omega}_{0}$.

Proof. (i) Case $1<q \leq n^{\prime}$. Then we have $n \leq q^{\prime}$. Let us take a subdomain $\Omega_{1}$ of $\Omega$ so that $\bar{\Omega}_{0} \subset \Omega_{1}$ and so that $D \equiv \Omega / \bar{\Omega}_{1}$ is a bounded domain in $\mathbb{R}^{n}$. We show first that the space

$$
S_{D} \equiv\left\{\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right) ; \int_{D} \phi(x) \mathrm{d} x=0\right\}
$$

is dense in $\hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$. Indeed, taking the sequence $\left\{\varsigma_{k}\right\}_{k=1}^{\infty}$ of $n$-dimensional cut-off functions as in the proof of Lemma 2.2, we see $\zeta_{k}(x)=1$ for $|x| \leq k$ and $\left\|\nabla \zeta_{k}\right\|_{q^{\prime}, \mathbb{R}^{n}} \leq C k^{-1+n / q^{\prime}}$ with $C$ independent of $k$. Letting $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$, we set $\phi_{k}(x)=\phi(x)-(\operatorname{vol} D)^{-1}\left(\int_{D} \phi(y) \mathrm{d} y\right) \cdot \zeta_{k}(x), \quad(k=1,2, \cdots)$. For large $k$, we have $\phi_{k} \in S_{D}$, so we may assume that $\phi_{k} \in S_{D}$ for all $k \geq 1$. Since $\left\|\nabla \phi_{k}-\nabla \phi\right\|_{q^{\prime}, \mathbb{R}^{n}} \leq C\|\phi\|_{L^{\infty}\left(\mathbb{R}^{n}\right)} \cdot k^{-1+n / q^{\prime}}$ and since $C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ is dense in $\hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$, we see that $S_{D}$ is dense in $\hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$ if $q^{\prime}>n$ i.e., if $1<q<n^{\prime}$. In case $q^{\prime}=n$ i.e., in case $q=n^{\prime}$, again by Mazur's theorem, we can choose a sequence $\left\{\bar{\phi}_{k}\right\}_{k=1}^{\infty}$ of convex combinations of $\phi_{k}^{\prime} s$ so that $\nabla \bar{\phi}_{k} \rightarrow \nabla \phi$ in $L^{n}\left(\mathbb{R}^{n}\right)^{n}$ as $k \rightarrow \infty$, and we see that $S_{D}$ is also dense in $\hat{H}_{0}^{1, n}\left(\mathbb{R}^{n}\right)$.

Let $f \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)$ with supp $f \subset \bar{\Omega}_{0}$. Taking a function $\eta \in C^{\infty}\left(\mathbb{R}^{n}\right)$ satisfying $0 \leq \eta \leq 1, \eta(x)=1$ for $x \in \bar{\Omega}_{1}$ and $\eta(x)=0$ for $x \in \mathbb{R}^{n} / \Omega$, we have

$$
\begin{equation*}
\left|(f, \phi)_{\mathbb{R}^{n}}\right|=\left|(f, \eta \phi)_{\Omega}\right| \leq C\|f\|_{-1, q}\left(\|(\nabla \eta) \phi\|_{q^{\prime}}+\|\nabla \phi\|_{q^{\prime}, \mathbb{R}^{n}}\right) \tag{2.5}
\end{equation*}
$$

for all $\phi \in S_{D}$ with $C$ independent of $\phi$. Since $\operatorname{supp} \nabla \eta \subset D$ and since $\int_{D} \phi(x) \mathrm{d} x=0$, we have by the Poincaré inequality on $D$ that $\|(\nabla \eta) \phi\|_{q^{\prime}} \leq$ ${ }^{D}\|\nabla \phi\|_{q^{\prime}, D}$. Hence from (2.5) it holds

$$
\left|(f, \phi)_{\mathbb{R}^{n}}\right| \leq C\|f\|_{-1, q}\|\nabla \phi\|_{q^{\prime}, \mathbb{R}^{n}} \text { for all } \phi \in S_{D}
$$

Since $S_{D}$ is dense in $\hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$, the above inequality holds for all $\phi \in \hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$, from which we get the desired result in case $1<q \leq n^{\prime}$.
(ii) Case $n^{\prime}<q<\infty$. Since $1<q^{\prime}<n$, we can take $r \in\left(q^{\prime}, \infty\right)$ so that $1 / r=1 / q^{\prime}-1 / n$. Then it follows from the Sobolev inequality in $\mathbb{R}^{n}$ that

$$
\|(\nabla \eta) \phi\|_{q^{\prime}} \leq C\|\phi\|_{q^{\prime}, D} \leq C\|\phi\|_{r, \mathbb{R}^{n}} \leq C\|\nabla \phi\|_{q^{\prime}, \mathbb{R}^{n}}
$$

for all $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$. Now we get the desired result by making use of (2.5) with $\phi \in S_{D}$ replaced by $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$.

### 2.2. Stokes equations in bounded domains.

In this subsection, we recall the $L^{q}$-theory for the Stokes equations in bounded domains due to Cattabriga [10].

THEOREM 2.6 (Cattabriga). Let $n \geq 2$ and $G \subset \mathbb{R}^{n}$ be a bounded domain with boundary $\partial G$ of class $C^{2+\mu}(0<\mu<1)$. Let $1<q<\infty$. Then for every $f \in \hat{H}^{-1, q}(G)^{n}$ and $g \in L^{q}(G)$ with $\int_{G} g(x) \mathrm{d} x=0$, there is a unique pair $\{u, p\} \in H_{0}^{1, q}(G)^{n} \times L^{q}(G)$ with $\int_{G} p(x) \mathrm{d} x=0$ such that

$$
\begin{equation*}
-\Delta u+\nabla p=f, \operatorname{div} u=g \text { in } G \tag{2.6}
\end{equation*}
$$

in the sense of distributions. Such $\{u, p\}$ is subject to the inequality

$$
\begin{equation*}
\|\nabla u\|_{q, G}+\|p\|_{q, G} \leq C\left(\|f\|_{-1, q, G}+\|g\|_{q, G}\right) \tag{2.7}
\end{equation*}
$$

where $C=C(G, n, q)$.
REMARK. Since $G$ is bounded, we have $\hat{H}^{-1, q}(G)=H_{0}^{1, q^{\prime}}(G)^{*}$. Cattabriga [10] gave the above result for $n=3$ under the weaker assumption that $\partial G$ is of class $C^{2}$. Galdi-Simader [18] extended Catabriga's result for $n \geq 2$. Another proof was given by Kozono-Sohr [22] (see also Borchers-Miyakawa [7]).

The following corollary is an immediate consequence of Theorem 2.6.
COROLLARY 2.7 (Regularity in bounded domains). Under the same assumption on $G, q, f$ and $g$ as in Theorem 2.6, suppose that $\{u, p\} \in$ $H_{0}^{1, q}(G)^{n} \times L^{q}(G)$ satisfies (2.6) in the sense of distributions. If, in addition, $f \in \hat{H}^{-1, r}(G)^{n}$ and $g \in L^{r}(G)$ for some $1<r<\infty$, then we have also $u \in H_{0}^{1, r}(G)^{n}$ and $p \in L^{r}(G)$.

### 2.3. Stokes equations in $\mathbb{R}^{n}$.

In this subsection, we shall give a result on $\mathbb{R}^{n}$ corresponding to that of the preceding subsection.

LEMMA 2.8 (Regularity theory in $\mathbb{R}^{n}$ ). Let $n \geq 2,1<q<\infty$ and let $f \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)^{n}, g \in L^{q}\left(\mathbb{R}^{n}\right)$. Suppose that $\{u, p\} \in L_{\mathrm{loc}}^{q}\left(\mathbb{R}^{n}\right)^{n} \times L^{q}\left(\mathbb{R}^{n}\right)$ with $\nabla u \in L^{q}\left(\mathbb{R}^{n}\right)^{n^{2}}$ satisfy

$$
\begin{equation*}
-\Delta u+\nabla p=f, \operatorname{div} u=g \text { in } \mathbb{R}^{n} \tag{2.8}
\end{equation*}
$$

in the sense of distributions. If, in addition, $f \in \hat{H}^{-1, r}\left(\mathbb{R}^{n}\right)^{n}$ and $g \in L^{r}\left(\mathbb{R}^{n}\right)$ for some $1<r<\infty$, then we have also $\nabla u \in L^{r}\left(\mathbb{R}^{n}\right)^{n^{2}}$ and $p \in L^{r}\left(\mathbb{R}^{n}\right)$.

PROOF. By Lemma 2.1, there are sequences $\left\{u_{j}\right\}_{j=1}^{\infty}$ in $C_{0}^{\infty}\left(\mathbb{R}^{n}\right)^{n}$ and $\left\{p_{j}\right\}_{j=1}^{\infty}$ in $C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ such that

$$
\begin{equation*}
\nabla u_{j} \rightarrow \nabla u \text { in } L^{q}\left(\mathbb{R}^{n}\right)^{n^{2}}, p_{j} \rightarrow p \text { in } L^{q}\left(\mathbb{R}^{n}\right) \tag{2.9}
\end{equation*}
$$

Set $f_{j}:=-\Delta u_{j}+\nabla p_{j}$ and $g_{j}:=\operatorname{div} u_{j}(j=1,2, \cdots)$. Then we have by (2.8) and (2.9)

$$
\begin{equation*}
\left(f_{j}, \Phi\right)_{\mathbb{R}^{n}} \rightarrow(f, \Phi)_{\mathbb{R}^{n}},\left(g_{j}, \phi\right)_{\mathbb{R}^{n}} \rightarrow(g, \phi) \tag{2.10}
\end{equation*}
$$

for all $\Phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)^{n}$ and all $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$, respectively. On the other hand, using the fundamental solution $F_{n}$ of $-\Delta$ in $\mathbb{R}^{n}$, we can represent $u_{j}$ and $p_{j}$ as

$$
u_{j}=F_{n} * f_{j}-F_{n} * \nabla p_{j}, p_{j}=-\operatorname{div} F_{n} *\left(f_{j}+\nabla g_{j}\right),
$$

where $*$ denotes the convolution. Then it follows that

$$
\left(p_{j}, \Delta \phi\right)_{\mathbb{R}^{n}}=-\left(f_{j}, \nabla \phi\right)_{\mathbb{R}^{n}}+\left(g_{j}, \Delta \phi\right)_{\mathbb{R}^{n}} \text { for all } \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right) .
$$

Letting $j \rightarrow \infty$ and then using the Calderon-Zygmund inequality, we have by (2.9)-(2.10) that

$$
\left|(p, \Delta \phi)_{\mathbb{R}^{n}}\right| \leq C\left(\|f\|_{-1, r, \mathbb{R}^{n}}+\|g\|_{r, \mathbb{R}^{n}}\right)\|\Delta \phi\|_{r^{\prime}, \mathbb{R}^{n}}
$$

for all $\phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$. Hence Lemma 2.4 states that $p \in L^{r}\left(\mathbb{R}^{n}\right)$. Concerning the regularity of $\nabla u$, we have similarly

$$
\left(\partial_{k} u_{j}, \Delta \Phi\right)_{\mathbb{R}^{n}}=\left(f_{j}, \partial_{k} \Phi\right)_{\mathbf{R}^{n}}+\left(p_{j}, \operatorname{div}\left(\partial_{k} \Phi\right)\right)_{\mathbf{R}^{n}},(k=1, \cdots, n),
$$

for all $\Phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)^{n}$. Since $p \in L^{r}\left(\mathbb{R}^{n}\right)$, the same argument as above yields

$$
\left|\left(\partial_{k} u, \Delta \Phi\right)_{\mathbb{R}^{n}}\right| \leq C\left(\|f\|_{-1, r, \mathbb{R}^{n}}+\|p\|_{r, \mathbb{R}^{n}}\right)\|\Delta \Phi\|_{r^{\prime}, \mathbb{R}^{n}}, \quad(k=1, \cdots, n),
$$

for all $\Phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)^{n}$. Again from Lemma 2.4 , we get $\partial_{k} u \in L^{\gamma}\left(\mathbb{R}^{n}\right)^{n}, \quad(k=$ $1, \cdots, n$ ).

Concerning the existence and uniqueness of solutions in the class $\hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)$, we have

Lemma 2.9 (A priori estimate in $\mathbb{R}^{n}$ ). Let $n \geq 2$ and $1<q<\infty$. Then for every $f \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)^{n}$ and $g \in L^{q}\left(\mathbb{R}^{n}\right)$, there is a unique $\{u, p\} \in$ $\hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)^{n} \times L^{q}\left(\mathbb{R}^{n}\right)$ such that (2.8) holds in the sense of distributions. Such $\{u, p\}$ is subject to the inequality

$$
\|\nabla u\|_{q, \mathbb{R}^{n}}+\|p\|_{q, \mathbb{R}^{n}} \leq C\left(\|f\|_{-1, q, \mathbb{R}^{n}}+\|g\|_{q, \mathbb{R}^{n}}\right),
$$

where $C=C(n, q)$.
Proof. By the definition of the space $\hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$, we see that the operator $-\nabla: \hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right) \rightarrow L^{q^{\prime}}\left(\mathbb{R}^{n}\right)^{n}$ is injective and has a closed range. Hence by the closed range theorem, the adjoint operator div $=(-\nabla)^{*}: L^{q}\left(\mathbb{R}^{n}\right)^{n} \rightarrow \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)$ is surjective. Since the null space $\operatorname{Ker}(\mathrm{div})$ of div is a closed subspace
in $L^{q}\left(\mathbb{R}^{n}\right)^{n}$, for each $h \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right)$, there is at least one $u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$ such that $-(u, \nabla \phi)_{\mathbb{R}^{n}}=(h, \phi)_{\mathbb{R}^{n}}$ holds for all $\phi \in \hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right)$ and that $\|u\|_{q, \mathbb{R}^{n}} \leq C\|h\|_{-1, q, \mathbb{R}^{n}}$ with $C$ independent of $h$. Let us now use the properties of the space $E_{q}$ and the bijective operator $B_{q}: E_{q} \rightarrow E_{q^{\prime}}^{*}$ in the proof of Lemma 2.3. Since $u \in L^{q}\left(\mathbb{R}^{n}\right)^{n}$, the $\operatorname{map} \nabla \phi \in E_{q^{\prime}} \rightarrow-(u, \nabla \phi)_{\mathbb{R}^{n}} \in \mathbb{R}$ is an element in $E_{q^{*}}^{*}$, so we can choose $\pi \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)$ so that

$$
\begin{equation*}
(\nabla \pi, \nabla \phi)_{\mathbb{R}^{n}}=-(u, \nabla \phi)_{\mathbb{R}^{n}}=(h, \phi)_{\mathbb{R}^{n}} \text { for all } \phi \in \hat{H}_{0}^{1, q^{\prime}}\left(\mathbb{R}^{n}\right) . \tag{2.11}
\end{equation*}
$$

By (2.1) such a $\pi$ is uniquely determined by $h$ and so (2.11) defines a bounded linear operator $J_{q}: h \in \hat{H}^{-1, q}\left(\mathbb{R}^{n}\right) \rightarrow \pi \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)$. A direct calculation shows that

$$
\left(\operatorname{div} J_{q}(\nabla \psi)+\psi, \Delta \phi\right)_{\mathbb{R}^{n}}=0 \text { for all } \psi \in L^{q}\left(\mathbb{R}^{n}\right), \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right) .
$$

Since the space $H=\left\{\Delta \phi ; \phi \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)\right\}$ is dense in $L^{q^{\prime}}\left(\mathbb{R}^{n}\right)$, the above identity yields that $\operatorname{div} J_{q}(\nabla \psi)=-\psi$ for all $\psi \in L^{q}\left(\mathbb{R}^{n}\right)$. Then we see that the pair $\{u, p\}$ defined by

$$
u=J_{q} f+J_{q}\left(\nabla \operatorname{div} J_{q}(f+\nabla g)\right), p=-\operatorname{div} J_{q}(f+\nabla g)
$$

has the desired property.
Now it remains to show the uniqueness. Let $\left\{u^{\prime}, p^{\prime}\right\} \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)^{n} \times L^{q}\left(\mathbb{R}^{n}\right)$ satisfy (2.8) in the sense of distributions. Then $\bar{u}=u-u^{\prime}, \bar{p}=p-p^{\prime}$ satisfies (2.8) with $f=0, g=0$. Applying the operator div to both sides of the first equation, we get $\Delta \bar{p}=0$ in the sense of distributions in $\mathbb{R}^{n}$. Since $\bar{p} \in L^{q}\left(\mathbb{R}^{n}\right)$, it follows from the Liouville theorem that $\bar{p} \equiv 0$ in $\mathbb{R}^{n}$. Therefore $\Delta \bar{u} \equiv 0$ in $\mathbb{R}^{n}$. Since $\bar{u} \in \hat{H}_{0}^{1, q}\left(\mathbb{R}^{n}\right)^{n}$, we have by (2.1) that $\bar{u} \equiv 0$ in $\mathbb{R}^{n}$.

Remark. There have been several results related to Lemma 2.9 (KozonoSohr [22, Proposition 2.9], Borchers-Miyakawa [7, Proposition 3.7], GaldiSimader [18, Theorem 3.1]). Our proof seems to be rather simple: we used only the variational inequality (2.1).

## 3. - Stokes equations in the class $\left(D_{q}\right)$

In this section we shall give some results in $\hat{H}_{0}^{1, q}(\Omega)$ analogous to those of subsection 2.3. In exterior domains, because of the boundary condition, we have restrictions on $q$ and $r$.

Theorem 3.1 (Regularity theory in $\Omega$ ). Let $n \geq 2,1<q<\infty$ and $r>n^{\prime}(=n /(n-1))$ and let $f \in \hat{H}^{-1, q}(\Omega)^{n} \cap \hat{H}^{-1, r}(\Omega)^{n}$ and $g \in L^{q}(\Omega) \cap L^{r}(\Omega)$. Suppose that $\{u, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ is a generalized solution of $(S)$. Then
we have $\nabla u \in L^{r}(\Omega)^{n^{2}}$ and $p \in L^{r}(\Omega)$. In case $r \geq n(n \geq 3)$ and $r>2(n=2)$, we have $u \in \hat{H}_{0}^{1, r}(\Omega)^{n}$, and in case $1<q<n$, we have also $u \in \hat{H}_{0}^{1, r}(\Omega)^{n}$.

Proof. We use the cut-off procedure. Take a ball $B_{R} \equiv\left\{x \in \mathbb{R}^{n} ;|x|<R\right\}$ so that $\partial \Omega \subset B_{R}$ and take a function $\psi_{1} \in C_{0}^{\infty}\left(\mathbb{R}^{n}\right)$ satisfying $0 \leq \psi_{1} \leq 1$, $\psi_{1}(x)=1$ for $x \in \mathbb{R}^{n} / \Omega, \psi_{1}(x)=0$ for $|x| \geq R$, and set $\psi_{2}=1-\psi_{1}$. Then from $(S)$ it follows

$$
\begin{equation*}
-\Delta\left(\psi_{i} u\right)+\nabla\left(\psi_{i} p\right)=f_{i}, \operatorname{div}\left(\psi_{i} u\right)=g_{i}(i=1,2) \tag{i}
\end{equation*}
$$

where $f_{i}=\psi_{i} f-2 \nabla \psi_{i} \nabla u-\left(\Delta \psi_{i}\right) u+\left(\nabla \psi_{i}\right) p, g_{i}=\psi_{i} g+\left(\nabla \psi_{i}\right) u(i=1,2)$. We may regard ( $S_{1}$ ) and ( $S_{2}$ ) as equations in $\Omega_{R} \equiv \Omega \cap B_{R}$ and in $\mathbb{R}^{n}$, respectively. Set $\Omega_{1} \equiv \Omega_{R}$ and $\Omega_{2} \equiv \mathbb{R}^{n}$.

Let us first assume that $1 / q-1 / n \leq 1 / r<1 / n^{\prime}=1-1 / n$. Taking $s \in(1, \infty)$ so that $1 / s=1 / r+1 / n$, we have $s \leq q$ and $1 / s^{\prime}=1 / r^{\prime}-1 / n$. Since $q^{\prime} \leq s^{\prime}$, it follows from the Sobolev embedding $H_{0}^{1, r^{\prime}}\left(\Omega_{i}\right) \subset L^{s^{\prime}}\left(\Omega_{i}\right)$ that

$$
\left\|\phi_{i}\right\|_{q^{\prime}, \Omega_{R}} \leq C\left\|\phi_{i}\right\|_{s^{\prime}, \Omega_{R}} \leq C\left\|\nabla \phi_{i}\right\|_{r^{\prime}, \Omega_{i}} \text { for all } \phi_{i} \in C_{0}^{\infty}\left(\Omega_{i}\right)(i=1,2)
$$

Since supp $\nabla \psi_{i}$ and $\operatorname{supp} \Delta \psi_{i}$ are contained in $\Omega_{R}$, we have by assumption and the above inequality that $f_{i} \in \hat{H}^{-1, r}\left(\Omega_{i}\right)^{n}(i=1,2)$. By the Sobolev embedding $H^{1, q}\left(\Omega_{R}\right) \subset L^{r}\left(\Omega_{R}\right)$, we get easily $g_{i} \in L^{r}\left(\Omega_{i}\right)(i=1,2)$, and also

$$
\int_{\Omega_{1}} g_{1} \mathrm{~d} x=\int_{\Omega_{R}} \operatorname{div}\left(\psi_{1} u\right) \mathrm{d} x=\int_{\partial \Omega} u \cdot \nu \mathrm{~d} S=0
$$

Now applying Corollary 2.7 and Lemma 2.8 to $\left\{\psi_{1} u, \psi_{1} p\right\}$ and $\left\{\psi_{2} u, \psi_{2} p\right\}$, respectively, we obtain

$$
\begin{equation*}
\nabla\left(\psi_{i} u\right) \in L^{r}\left(\Omega_{i}\right)^{n^{2}}, \psi_{i} p \in L^{r}\left(\Omega_{i}\right)(i=1,2) \tag{3.1}
\end{equation*}
$$

We next consider the case $1 / q-2 / n \leq 1 / r<1 / q-1 / n$. Taking $\bar{q}=(1 / q-1 / n)^{-1}$, we have by (3.1) that $\nabla u \in L^{\bar{q}}(\Omega)^{n^{2}}$ and $p \in L^{\bar{q}}(\Omega)$. Now, taking $\bar{q}$ instead of $q$ in the above, we get (3.1) also for $r>n^{\prime}$ with $1 / r \geq 1 / q-2 / n$. Proceeding in the same way to the case $1 / r<1 / q-2 / n$, by the bootstrap argument, we get (3.1) for all $r>n^{\prime}$ and hence $\nabla u \in L^{r}(\Omega)^{n^{2}}, p \in L^{r}(\Omega)$ for all $r>n^{\prime}$.

It remains to show that $u \in \hat{H}_{0}^{1, r}(\Omega)^{n}$ in case $r \geq n(n \geq 3), r>2(n=2)$, and in case $1<q<n$. To this end, we may show $\psi_{2} u \in \hat{H}_{0}^{1, r}(\Omega)^{n}$ in (3.1). Consider first the case when $r \geq n(n \geq 3)$ and $r>2(n=2)$. Since $\nabla\left(\psi_{2} u\right) \in L^{r}\left(\mathbb{R}^{n}\right)^{n^{2}}$ and since $\psi_{2} u$ vanishes in a neighbourhood of $\partial \Omega$, we get by Lemma 2.2(ii) that $\psi_{2} u \in \hat{H}_{0}^{1, r}(\Omega)^{n}$. We next consider the case when $1<q<n, n^{\prime}<r<n(n \geq 3)$. Since $f_{2} \in \hat{H}^{-1, r}\left(\mathbb{R}^{n}\right)^{n}$ and $g_{2} \in L^{r}\left(\mathbb{R}^{n}\right)$, it follows from Lemma 2.9 that there is a unique pair $\{v, \chi\} \in \hat{H}_{0}^{1, r}\left(\mathbb{R}^{n}\right)^{n} \times L^{r}\left(\mathbb{R}^{n}\right)$ such that $-\Delta v+\nabla \chi=f_{2}$, $\operatorname{div} v=g_{2}$ in the sense of distributions in $\mathbb{R}^{n}$. Taking
$w=v-\psi_{2} u$ and $\eta=\chi-\psi_{2} p$, we see that $\{w, \eta\}$ satisfies (2.8) with $f=0$ and $g=0$; applying div to both sides of the first equation, we obtain that $\eta$ is harmonic in $\mathbb{R}^{n}$. Since $\eta \in L^{r}\left(\mathbb{R}^{n}\right)$, the Liouville theorem yields that $\eta \equiv 0$ in $\mathbb{R}^{n}$; hence $w$ is also harmonic in $\mathbb{R}^{n}$. Moreover, by the Sobolev embedding theorem, we obtain $w \in L^{\bar{q}}\left(\mathbb{R}^{n}\right)+L^{\bar{r}}\left(\mathbb{R}^{n}\right)$, where $1 / \bar{q}=1 / q-1 / n$ and $1 / \bar{r}=1 / r-1 / n$. Using the same argument as in the proof of Lemma 2.4 , we get $w \equiv 0$ in $\mathbb{R}^{n}$, from which $\psi_{2} u \in \hat{H}_{0}^{1, r}\left(\mathbb{R}^{n}\right)^{n}$ follows. Now, again by Lemma 2.2(i), we have $\psi_{2} u \in \hat{H}_{0}^{1, \tau}(\Omega)^{n}$.

REMARK. The restriction $n^{\prime}<r$ is a critical condition; we cannot take $1<r \leq n^{\prime}$ in Theorem 3.1. Indeed, let us assume the main results in Section 1. Taking some $n<q<\infty$ and $A=0, a \neq 0$ in Theorem B (ii), we get such $\{u, p\} \in \mathbb{N}_{q}^{0}$ as $\lim _{x \rightarrow \infty} u(x)=a$ in case $n \geq 3$ and as $\int_{\Omega}|\nabla u(x)-\nabla E(x) a|^{2} \mathrm{~d} x<\infty$ in case $n=2$, and by Lemma 2.2(ii), we have $u \in \hat{H}_{0}^{1, q}(\Omega)^{n}$. Suppose now that Theorem 3.1 is true for $1<r \leq n^{\prime}$. Then it follows that $\nabla u \in L^{r}(\Omega)^{n^{2}}$ for some $1<r \leq n^{\prime}$. Thus by Theorem A, we get $u \equiv 0$ in $\Omega$, which contradicts $a \neq 0$. Note that $\int_{\Omega}|\nabla E(x)|^{2} \mathrm{~d} x=\infty$ in case $n=2$.

We shall next give an a priori estimate in the class $\left(D_{q}\right)$. For this purpose we need:

LEMMA 3.2. Let $n \geq 2$ and $1<q<\infty$ and let $\{f, g\} \in \hat{H}^{-1, q}(\Omega)^{n} \times L^{q}(\Omega)$. Suppose that $\{u, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ is a generalized solution of $(S)$. Then we have

$$
\begin{align*}
& \|\nabla u\|_{q}+\|p\|_{q} \\
& \leq C\left(\|f\|_{-1, q}+\|g\|_{q}+\|u\|_{q, \Omega_{R}}+\|p\|_{-1, q, \Omega_{R}}+\left|\int_{\Omega_{R}} \psi_{1}(x) p(x) \mathrm{d} x\right|\right) \tag{3.2}
\end{align*}
$$

where $\Omega_{R}=\Omega \cap B_{R}$ and $\psi_{1}$ are the same as in the proof of Theorem 3.1 and where $C$ is a constant independent of $u$ and $p$.

Proof. We use again the cut-off method. Recalling the equations $\left(S_{i}\right)(i=$ $1,2)$ in the proof of Theorem 3.1, we first consider $\left(S_{1}\right)$ in $\Omega_{R}$. Since $\operatorname{supp} \nabla \psi_{1}$, $\operatorname{supp} \Delta \psi_{1} \subset \Omega_{R}$, we obtain

$$
\begin{aligned}
& \left\|f_{1}\right\|_{-1, q, \Omega_{R}} \leq C\left(\|f\|_{-1, q}+\|u\|_{q, \Omega_{R}}+\|p\|_{-1, q, \Omega_{R}}\right) \\
& \left\|g_{1}\right\|_{q, \Omega_{R}} \leq C\left(\|g\|_{q}+\|u\|_{q, \Omega_{R}}\right)
\end{aligned}
$$

Applying Theorem 2.6 to $\left\{\psi_{1} u, \psi_{1} p\right\}$ in $\left(S_{1}\right)$ and then using the above inequa-
lities, we get

$$
\begin{align*}
& \left\|\nabla\left(\psi_{1} u\right)\right\|_{q, \Omega_{R}}+\left\|\psi_{1} p\right\|_{q, \Omega_{R}} \\
& \leq C\left(\|f\|_{-1, q}+\|g\|_{q}+\|u\|_{q, \Omega_{R}}+\|p\|_{-1, q, \Omega_{R}}+\left|\int_{\Omega_{R}} \psi_{1}(x) p(x) \mathrm{d} x\right|\right) \tag{3.3}
\end{align*}
$$

with $C$ independent of $u$ and $p$.
We next consider $\left(S_{2}\right)$ in $\mathbb{R}^{n}$. Since supp $f_{2} \subset \operatorname{supp} \psi_{2}$, it follows from Lemma 2.5 that

$$
\begin{equation*}
\left\|f_{2}\right\|_{-1, q, \mathbb{R}^{n}} \leq C\left\|f_{2}\right\|_{-1, q} \text { with } C=C(n, q) \tag{3.4}
\end{equation*}
$$

Since the inequality $\|\phi\|_{q^{\prime}, \Omega_{R}} \leq C\|\nabla \phi\|_{q^{\prime}}$ holds for all $\phi \in C_{0}^{\infty}(\Omega)$ and since $\operatorname{supp} \nabla \phi_{2}$, supp $\Delta \phi_{2} \subset \Omega_{R}$, we obtain

$$
\begin{align*}
& \left\|f_{2}\right\|_{-1, q} \leq C\left(\|f\|_{-1, q}+\|u\|_{q, \Omega_{R}}+\|p\|_{-1, q, \Omega_{R}}\right)  \tag{3.5}\\
& \left\|g_{2}\right\|_{q} \leq C\left(\|g\|_{q}+\|u\|_{q, \Omega_{R}}\right)
\end{align*}
$$

Now applying Lemma 2.9 to $\left\{\psi_{2} u, \psi_{2} p\right\}$ in ( $S_{2}$ ) and then using (3.4)-(3.5), we obtain

$$
\begin{align*}
& \left\|\nabla\left(\psi_{2} u\right)\right\|_{q, \mathbb{R}^{n}}+\left\|\psi_{2} p\right\|_{q, \mathbb{R}^{n}}  \tag{3.6}\\
& \leq C\left(\|f\|_{-1, q}+\|g\|_{q}+\|u\|_{q, \Omega_{R}}+\|p\|_{-1, q, \Omega_{R}}\right)
\end{align*}
$$

Then the desired result follows from (3.3) and (3.6).
Now we introduce the weak Stokes operator $S_{q}$. Let $X_{q} \equiv \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ and $Y_{q} \equiv \hat{H}^{-1, q}(\Omega)^{n} \times L^{q}(\Omega)$. We define two bounded linear operators $S_{q}$ and $T_{q}$ by

$$
\begin{aligned}
& S_{q}:\{u, p\} \in X_{q} \rightarrow\{-\Delta u+\nabla p, \operatorname{div} u\} \in Y_{q} \\
& T_{q}:\{v, \chi\} \in X_{q} \rightarrow\{-\Delta v-\nabla \chi,-\operatorname{div} v\} \in Y_{q},
\end{aligned}
$$

respectively. It is easy to see that

$$
\begin{equation*}
S_{q}^{*}\left(\text { adjoint operator of } S_{q}\right)=T_{q^{\prime}} \text { for all } 1<q<\infty \tag{3.7}
\end{equation*}
$$

Then Lemma 3.2 enables us to apply such a standard argument as Lions-Magenes [24, p. 153, Lemma 5.1], so we see that

Ker $S_{q}$ (the kernel of $S_{q}$ ) is of finite dimension and $R\left(S_{q}\right)$ (the range of $S_{q}$ ) closed in $Y_{q}$.

More precisely we have

THEOREM 3.3 (A priori estimate). Let $n \geq 2,1<q<n$ for $n \geq 3$ and $1<q \leq 2$ for $n=2$, and let $\{f, g\} \in \hat{H}^{-1, q}(\Omega)^{n} \times L^{q}(\Omega)$. Suppose that $\{u, p\} \in \hat{H}_{0}^{1, \bar{q}}(\Omega)^{n} \times L^{q}(\Omega)$ is a generalized solution of $(S)$. Then it holds

$$
\begin{equation*}
\|\nabla u\|_{q}+\|p\|_{q} \leq C\left(\|f\|_{-1, q}+\|g\|_{q}\right) \tag{3.9}
\end{equation*}
$$

where $C=C(\Omega, q, n)$.
Proof. We show first that $\operatorname{Ker} S_{q}=\{0,0\}$ for such $q$ as in the theorem. Let $\{u, p\} \in \operatorname{Ker} S_{q}$. Then it is enough to show that $\{u, p\} \in \hat{H}_{0}^{1,2}(\Omega)^{n} \times L^{2}(\Omega)$, because we can insert $\Phi=u$ as a test function in the definition of the generalized solution and hence $\|\nabla u\|_{2}^{2}=0, \nabla p=0$ follows. Then we get $u \equiv 0, p \equiv 0$. If $n \geq 3$, then we can take $r=2>n^{\prime}$ in Theorem 3.1 and get $\{u, p\} \in \hat{H}_{0}^{1,2}(\Omega)^{n} \times L^{2}(\Omega)$. If $n=2$, then we get by Theorem 3.1 and the interpolation property that $\nabla u \in L^{r}(\Omega)^{n^{2}}, p \in L^{r}(\Omega)$ for all finite $r \geq q$. Since $n=2$, it follows from Lemma 2.2(ii) that $u \in \hat{H}_{0}^{1,2}(\Omega)^{2}$.

Now we prove (3.9) by contradiction. Suppose the contrary. Then there is a sequence $\left\{u_{k}, p_{k}\right\}_{k=1}^{\infty}$ in $\hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ such that $\left\|\nabla u_{k}\right\|_{q}+\left\|p_{k}\right\|_{q}=1$ and that $-\Delta u_{k}+\nabla p_{k} \rightarrow 0$ in $\hat{H}^{-1, q}(\Omega)^{n}$, div $u_{k} \rightarrow 0$ in $L^{q}(\Omega)$ as $k \rightarrow \infty$. A well known compactness argument yields that there is a subsequence, which we denote by $\left\{u_{k}, p_{k}\right\}_{k=1}^{\infty}$ for simplicity, such that $\left\{u_{k}\right\}_{k=1}^{\infty},\left\{p_{k}\right\}_{k=1}^{\infty}$ and $\left\{\int_{\Omega_{R}} \psi_{1}(x) p_{k}(x) \mathrm{d} x\right\}_{k=1}^{\infty}$ converge strongly in $L^{q}\left(\Omega_{R}\right)^{n}, \quad \hat{H}^{-1, q}\left(\Omega_{R}\right)$ and $\mathbb{R}$, respectively. Then, applying Lemma 3.2 to $\left\{u_{k}-u_{k^{\prime}}, p_{k}-p_{k^{\prime}}\right\}_{k, k^{\prime}=1}^{\infty}$, we see that $\left\{u_{k}\right\}_{k=1}^{\infty}$ and $\left\{p_{k}\right\}_{k=1}^{\infty}$ are Cauchy sequences in $\hat{H}_{0}^{1, q}(\Omega)^{n}$ and in $L^{q}(\Omega)$, respectively. Thus, there is a pair $\{u, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ such that $u_{k} \rightarrow u$ in $\hat{H}_{0}^{1, q}(\Omega)^{n}$ and $p_{k} \rightarrow p$ in $L^{q}(\Omega)$. Moreover, we have $\{u, p\} \in \operatorname{Ker} S_{q}$ and that $\|\nabla u\|_{q}+\|p\|_{q}=1$, but this contradicts Ker $S_{q}=\{0,0\}$.

Using (3.7)-(3.8) and a closed range theorem, we have by Theorem 3.3 the following corollary.

COROLLARY 3.4. (i) Let $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Then we have $\operatorname{Ker} S_{q}=\{0,0\}$ and $R\left(S_{q}\right)=\left(\operatorname{Ker} T_{q^{\prime}}\right)^{\perp}$.
(ii) Let $n^{\prime}<q<n$ for $n \geq 3$ and $q=2$ for $n=2$. Then we have Ker $S_{q}=\{0,0\}$ and $R\left(S_{q}\right)=Y_{q}$.
(iii) Let $n \leq q<\infty$ for $n \geq 3$ and $2<q<\infty$ for $n=2$. Then we have Ker $S_{q}=R\left(T_{q^{\prime}}\right)^{\perp}$ and $R\left(S_{q}\right)=Y_{q}$.

Here $W^{\perp}$ denotes the annihilator of the subspace $W$.
REMARK. Theorem 3.3 was first proved by Kozono-Sohr [22] in case $n \geq 3$ and $n^{\prime}<q<n$. Borchers-Miyakawa [7] extended the result to the case when $n \geq 3$ and $1<q \leq n^{\prime}$. Recently Galdi-Simader [18] gave a similar result for $n \geq 2$, but with a different method from ours.

## 4. - Proof of the main results

4.1. Stokes paradox; Proof of Theorems A and $\mathrm{A}^{\prime}$.

Let us first give the following auxiliary lemma due to Bogovskii [5, 6].
Lemma 4.1 (Bogovskii). (i) Let $1<q<\infty$. Suppose that $w \in \hat{H}^{-1, q}(\Omega)^{n}$ satisfies $(w, \Phi)=0$ for all $\Phi \in C_{0}^{\infty}(\Omega)^{n}$ with $\operatorname{div} \Phi=0$. Then there is a unique $p \in L^{q}(\Omega)$ such that $w=\nabla p$, i.e., $(w, \Psi)=-(p$, div $\Psi)$ for all $\Psi \in \hat{H}_{0}^{1, q^{\prime}}(\Omega)^{n}$.
(ii) Let $1<r<n$ and let $u \in L_{\mathrm{loc}}^{1}(\bar{\Omega})$ with $\nabla u \in L^{r}(\Omega)^{n}$. Then there is a constant $C=C(u, n, r)$ such that $u+C \in L^{q}(\Omega)$ with $1 / q=1 / r-1 / n$.

For the proof, see also Giga-Sohr [19, Corollary 2.2] and Borchers-Sohr [9, Lemma 4.1].

In the forthcoming argument, we use the linear extension operator $\Gamma: C^{2}(\partial \Omega)^{n} \rightarrow C_{0}^{2}\left(B_{R}\right)^{n}$ satisfying

$$
\Gamma \phi=\phi \text { on } \partial \Omega,\|\Gamma \phi\|_{H^{m, q}\left(B_{R}\right)} \leq C\|\phi\|_{H^{m-1 / q, q}(\partial \Omega)}, \quad(m=1,2)
$$

for all $\phi \in C^{2}(\partial \Omega)^{n}$ with $C=C(\partial \Omega, R, n, m, q)$. Here $B_{R}=\left\{x \in \mathbb{R}^{n} ;|x|<R\right\}$ is a ball containing $\partial \Omega$.

Proof of Theorem A. If $n=n^{\prime}=q=2$, then the desired result follows from Corollary 3.4(ii) and Lemma 2.2(ii), so we may prove only the case $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Since $\nabla u \in L^{q}(\Omega)^{n^{2}}$, it follows from Lemma 4.1(ii) that there is a constant vector $a=a(u, q) \in \mathbb{R}^{n}$ such that $u-a \in L^{n q /(n-q)}(\Omega)^{n}$. Set $w=\Gamma a \in C_{0}^{2}\left(B_{R}\right)$ and define $\hat{u}=u-a+w$. Then we see by Lemma 2.2(i) that $\hat{u} \in \hat{H}_{0}^{1, q}(\Omega)^{n}$, and by assumption we get

$$
\begin{equation*}
-\Delta \hat{u}+\nabla p=-\Delta w, \operatorname{div} \hat{u}=\operatorname{div} w \text { in } \Omega \tag{4.1}
\end{equation*}
$$

in the sense of distributions. Since $\{-\Delta w$, $\operatorname{div} w\} \in \hat{H}^{-1, \gamma}(\Omega)^{n} \times L^{\gamma}(\Omega)$ for all $\gamma>1$, it follows from Theorem 3.1 that $\{\hat{u}, p\} \in \hat{H}_{0}^{1, r}(\Omega)^{n} \times L^{r}(\Omega)$ for all $r>n^{\prime}$. Moreover, since $q^{\prime} \geq n$ for $n \geq 3$ and $q^{\prime}>2$ for $n=2$, we obtain from Lemma 2.2(ii) that $\{u,-p\} \in \operatorname{Ker} T_{q^{\prime}}$. Therefore it follows from (4.1) and Corollary 3.4(i) that

$$
\begin{aligned}
0 & =(-\Delta w, u)+(\operatorname{div} w,-p) \\
& =-\int_{\Omega}|\nabla u(x)|^{2} \mathrm{~d} x+(\nabla u, \nabla \hat{u})-(p, \operatorname{div} \hat{u}) .
\end{aligned}
$$

Since $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q^{\prime}}(\Omega)^{n} \times L^{q^{\prime}}(\Omega)$ and since $C_{0}^{\infty}(\Omega)^{n} \times C_{0}^{\infty}(\Omega)$ is dense in $\hat{H}_{0}^{1, q^{\prime}}(\Omega)^{n} \times L^{q^{\prime}}(\Omega)$, it follows from the assumption on $\{u, p\}$ that $(\nabla u, \nabla \hat{u})-$ $(p, \operatorname{div} \hat{u})=0$. So we get $u \equiv 0$ in $\Omega$, then $\nabla p=0$ in $\Omega$. Since $p \in L^{q}(\Omega)$, we have also $p \equiv 0$ in $\Omega$.

Proof of Theorem A'. Since $\nabla u \in L^{q}(\Omega)^{n^{2}}$, we have $-\Delta u \in \hat{H}^{-1, q}(\Omega)^{n}$ and, by assumption, $(-\Delta u, \Phi)=0$ for all $\Phi \in C_{0}^{\infty}(\Omega)^{n}$ with $\operatorname{div} \Phi=0$. Then it follows from Lemma 4.1(i) that there is a scalar function $p \in L^{q}(\Omega)$ such that

$$
(\nabla u, \nabla \Psi)-(p, \operatorname{div} \Psi)=0 \text { for all } \Psi \in C_{0}^{\infty}(\Omega)^{n}
$$

Now Theorem A yields that $u \equiv 0$ in $\Omega$.

### 4.2. Characterization of the null space; Proof of Theorem B.

We shall first consider the cases (i) and (ii), i.e., $n \geq 2,1<q<\infty$ except for $n=q=2$. The proof will be done by three lemmas. Let us define the vector spaces $V$ and $\hat{V}_{q}$ for $1<q \leq n^{\prime}(n \geq 3)$ and $1<q<2(n=2)$ as follows:

$$
\begin{aligned}
V= & \left\{\{A, a\} \in \mathbb{R}^{n^{2}} \times \mathbb{R}^{n} ; \operatorname{Tr} A=0\right\} \\
\hat{V}_{q}= & \left\{\{A, a\} \in V ; \int_{\partial \Omega}\left[(A x+a) \cdot \frac{\partial v}{\partial \nu}-\chi(A x+a) \cdot \nu\right] \mathrm{d} S=0\right. \\
& \text { for all } \left.\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}\right\} .
\end{aligned}
$$

Then the existence of a generalized solution with (1.2) and (1.3-3') in Theorem $\mathrm{B}(\mathrm{i})$-(ii) is guaranteed by the following lemma.

LEMMA 4.2 (Existence). (i) Let $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Then there is a linear operator $K_{q}:\{A, a\} \rightarrow\{u, p\}$ from $\hat{V}_{q}$ to $\underset{r \geq q}{\cap} \mathbb{N}_{r}$ such that $\nabla u-A \in L^{r}(\Omega)^{n^{2}}$ for all $r \geq q$ and such that (1.3) holds.
(ii) There is a linear operator $L:\{A, a\} \rightarrow\{u, p\}$ from $V$ to $\bigcap_{r>n^{\prime}} \mathbb{N}_{r}$ such that: $\nabla u-A \in L^{r}(\Omega)^{n^{2}}$ for all $r>n^{\prime}$, (1.3) holds if $n \geq 3$, and

$$
\begin{equation*}
\int_{\Omega}|\nabla[u(x)-A x-E(x) a]|^{s} \mathrm{~d} x<\infty \tag{4.2}
\end{equation*}
$$

holds for all $s \geq 2$ if $n=2$.
Proof. (i) Since $q^{\prime} \geq n$ for $n \geq 3$ and $q^{\prime}>2$ for $n=2$, we have by Lemma 2.2(ii) that $\{v,-\chi\} \in \operatorname{Ker} T_{q^{\prime}}$ for $\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}$. Taking $w=\Gamma(A \cdot+a)$ for $\{A, a\} \in \hat{V}_{q}$, where $\Gamma$ is the extension operator defined above and where $A \cdot+a$ is a function on $\partial \Omega$ defined by $x \in \partial \Omega \rightarrow A x+a \in \mathbb{R}^{n}$, a direct calculation shows that

$$
\begin{aligned}
& (-\Delta w, v)+(\operatorname{div} w,-\chi) \\
& =\int_{\partial \Omega}\left\{(A x+a) \cdot \frac{\partial v}{\partial \nu}-\chi(A x+a) \cdot \nu\right\} \mathrm{d} S=0
\end{aligned}
$$

for all $\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}$. Hence it follows from Corollary 3.4(i) that there is a unique generalized solution $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ of (4.1). Moreover, by Theorem 3.1 and the interpolation inequality, $\nabla \hat{u} \in L^{r}(\Omega)^{n^{2}}$ and $p \in L^{r}(\Omega)$ for all $r \geq q$. If $n \geq 3$, we have again by Theorem 3.1 that $\hat{u} \in \hat{H}_{0}^{1, \gamma}(\Omega)^{n}$ for all $\gamma>n^{\prime}$. For $\gamma \in(n / 2, n), \hat{H}_{0}^{1, \gamma}(\Omega)$ is continuously embedded into $L^{n \gamma /(n-\gamma)}(\Omega)$ (see Lemma 2.2(i)). For such $\gamma$, we have $n \gamma /(n-\gamma)>n$ and hence, in particular, $\hat{u} \in H^{1, s}(\Omega)^{n}$ for $s>n$. By the Sobolev embedding theorem, $\hat{u} \in C^{0}(\bar{\Omega})^{n}$ and $\lim _{x \rightarrow \infty}|\hat{u}(x)|=0$. If $n=2$, we have by Lemma 2.2(i) that $\hat{u} \in H_{0}^{1,2 q /(2-q)}(\Omega)^{2}$. Since $2 q /(2-q)>2, \hat{u}$ has the same properties as above. Now, setting $u=\hat{u}+A x+a-w$ and then defining $K_{q}\{A, a\}=\{u, p\}$, we obtain the operator $K_{q}$.
(ii) We first consider the case $n \geq 3$. Let $\{A, a\} \in V$. We set $w=\Gamma(A \cdot+a)$. Then it follows from Corollary 3.4(ii) and Theorem 3.1 that there is a unique generalized solution $\{\hat{u}, p\}$ of (4.1) such that $\{\hat{u}, p\} \in \hat{H}_{0}^{1, r}(\Omega)^{n} \times L^{r}(\Omega)$ for all $r>n^{\prime}$. Then in the same way as above, we can show that $\hat{u} \in C^{0}(\bar{\Omega})^{n}$, $\lim _{x \rightarrow \infty}|\hat{u}(x)|=0$ and that the map $L:\{A, a\} \rightarrow\{u, p\}$ with $u=\hat{u}+A x+a-w$ satisfies the required conditions.

We next construct $L$ for $n=2$. Without loss of generality, we may assume that $0 \in \mathbb{R}^{n} / \bar{\Omega}$. Set $w=\Gamma(A \cdot+E a)$, where $A \cdot+E a$ is the function on $\partial \Omega$ defined by $x \in \partial \Omega \rightarrow A x+E(x) a \in \mathbb{R}^{2}$. Then it follows from Corollary 3.4(ii) and Theorem 3.1 that there is a unique generalized solution $\{\hat{u}, p\}$ of (4.1) belonging to $\hat{H}_{0}^{1, r}(\Omega)^{2} \times L^{r}(\Omega)$ for all $r \geq 2$. Setting $u=\hat{u}+A x+E a-w$, we see that the map $L:\{A, a\} \rightarrow\{u, p\}$ enjoys the desired properties.

We next show the uniqueness of generalized solutions.
LEMMA 4.3 (Uniqueness). (i) Let $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. Then for every $\{A, a\} \in \hat{V}_{q}$, there is a unique $\{u, p\} \in \mathbb{N}_{q}$ with properties (1.2) and (1.3).
(ii) Let $n^{\prime}<q<\infty, n \geq 2$. Then for every $\{A, a\} \in V$, there exists $a$ unique $\{u, p\} \in \mathbb{N}_{q}$ with properties (1.2)-(1.3) if $n \geq 3$, and (1.2)-(1.3') if $n=2$.

Proof. The proof of existence is contained in Lemma 4.2 so we may only prove uniqueness.
(i) Suppose that $\{\tilde{u}, \tilde{p}\} \in \mathbb{N}_{q}$ satisfies (1.2) and (1.3) with $u$ replaced by $\tilde{u}$. Set $\bar{u}=u \tilde{u}$ and $\bar{p}=p-\tilde{p}$ (note that $p, \tilde{p}$ and $\bar{p}$ do not denote integral exponents but functions of the pressure). Then we have $\nabla \bar{u} \in L^{q}(\Omega)^{n^{2}}, \bar{u} \in C^{0}(\bar{\Omega})^{n}$ and from Theorem 3.1 that there is a unique generalized solution $\lim _{x \rightarrow \infty}|\bar{u}(x)|=0$. On the other hand, by Lemma 4.1 (ii), there is a constant vector $C \in \mathbb{R}^{n}$ such that $\bar{u}+C \in L^{n q /(n-q)}(\Omega)^{n}$. Since $\lim _{x \rightarrow \infty}|\bar{u}(x)|=0$, we have $C=0$ and hence $\bar{u} \in L^{n q /(n-q)}(\Omega)^{n}$. Then from Lemma 2.2(i) we obtain $\bar{u} \in \hat{H}_{0}^{1, q}(\Omega)^{n}$ and so $\{\bar{u}, \bar{p}\} \in \operatorname{Ker} S_{q}$. Now applying Corollary 3.4(i), we get $\bar{u} \equiv 0, \bar{p} \equiv 0$ and the assertion on uniqueness follows.
(ii) Let us first assume $n \geq 3$. Let $\{\tilde{u}, \tilde{p}\} \in \mathbb{N}_{q}$ and $\{\bar{u}, \bar{p}\}$ as above. If $n^{\prime}<q<n$, we can argue in the same way as above and get $\{\bar{u}, \bar{p}\} \in \operatorname{Ker} S_{q}$.

Then it follows from Corollary 3.4(ii) that $\bar{u} \equiv 0, \bar{p} \equiv 0$. If $n \leq q<\infty$, we see by Lemma 2.2(ii) that $\bar{u} \in \hat{H}_{0}^{1, q}(\Omega)^{n}$. From Theorem 3.1, we obtain $\nabla \bar{u} \in L^{r}(\Omega)^{n^{2}}$, $\bar{p} \in L^{r}(\Omega)$ for all $r>n^{\prime}$. In the same way as in (i), we get $\{\bar{u}, \bar{p}\} \in \operatorname{Ker} S_{\gamma}$ for some $\gamma$ with $n^{\prime}<\gamma<n$ and Corollary 3.4(ii) yields $\bar{u} \equiv 0, \bar{p} \equiv 0$. In case $n=2$, we see by Lemma 2.2(ii) and (1.3') that $\{\bar{u}, \bar{p}\} \in \operatorname{Ker} S_{2}=\{0,0\}$.

Now it remains to give the dimensions of $\mathbb{N}_{q}$ and $\mathbb{N}_{q}^{0}$. To this end, we shall make use of the operators $K_{q}$ and $L$ constructed in Lemma 4.2.

LEMMA 4.4. (i) For each $q$ with $1<q \leq n^{\prime}(n \geq 3)$ and with $1<q<2(n=2), K_{q}$ defines a bijection from $\hat{V}_{q}$ onto $\mathbb{N}_{q}$.
(ii) For each $q$ with $n^{\prime}<q<\infty(n \geq 2)$, $L$ defines a bijection from $V$ onto $\mathbb{N}_{q}$.

Proof. (i) Injectivity. Let $K_{q}\{A, a\}=\{0,0\}$ for $\{A, a\} \in \hat{V}_{q}$. Then by (1.3), $|A x+a| \rightarrow 0$ as $x \rightarrow \infty$; hence we get $A=0$ and $a=0$.

Surjectivity. Suppose that $\{u, p\} \in \mathbb{N}_{q}$. Then $\nabla u-A \in L^{q}(\Omega)^{n^{2}}$ for some $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$. By Lemma 4.1(ii), there is a constant vector $a \in \mathbb{R}^{n}$ such that $u-A x-a \in L^{n q /(n-q)}(\Omega)^{n}$. Introducing $w=\Gamma(A \cdot+a)$ as in the proof of Lemma 4.2 and then defining $\hat{u}=u-A x-a+w$, we see by Lemma 2.2(i) that $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ and that $\{\hat{u}, p\}$ satisfies (4.1) in the sense of distributions. Moreover it follows from Theorem 3.1 that $\hat{u} \in \hat{H}_{0}^{1, r}(\Omega)^{n}$ for all $r>n^{\prime}$. Now using the same argument as in the proof of Lemma 4.2, we get $\hat{u} \in C^{0}(\bar{\Omega})^{n}, \lim _{x \rightarrow \infty}|\hat{u}(x)|=0$ and hence $u$ satisfies (1.3). Then the uniqueness stated in Lemma 4.3(i) yields that $\{u, p\}=K_{q}\{A, a\}$.
(ii) Injectivity. Let $L\{A, a\}=\{0,0\}$ for $\{A, a\} \in V$. If $n \geq 3$, we get in the same way as above $A=0, a=0$. If $n=2$, we obtain by (4.2) that $A-\nabla E a \in L^{s}(\Omega)^{2^{2}}$ for all $s \geq 2$. The explicit expression of $E$ shows that $\nabla E$ is not in $L^{2}(\Omega)^{2^{3}}$, but in $L^{r}(\Omega)^{2^{3}}$ for all $r>2$. Hence $A=0, a=0$.

Surjectivity. Let us first assume that $n \geq 3$. The proof for $q$ with $n^{\prime}<q<n$ is parallel to that of case (i), so we may only show it for $n \leq q<\infty$. Suppose that $\{u, p\} \in \mathbb{N}_{q}(q \geq n)$. Then $\nabla u-A \in L^{q}(\Omega)^{n^{2}}$ for some $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A=0$, and taking $w=\Gamma(A \cdot)$ and $\hat{u}=u-A x+w$, we see by Lemma 2.2(ii) that $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ and that $\{\hat{u}, p\}$ is a generalized solution of (4.1). Moreover, by Theorem 3.1, $\nabla \hat{u} \in L^{r}(\Omega)^{n^{2}}$ for all $r>n^{\prime}$, and in particular, we have $\nabla \hat{u} \in L^{\gamma}(\Omega)^{n^{2}}$ for $\gamma$ with $n / 2<\gamma<n$. By Lemma 4.1(ii), there is a constant vector $a \in \mathbb{R}^{n}$ such that $\hat{u}-a \in L^{\sigma}(\Omega)^{n}$ with $1 / \sigma=1 / \gamma-1 / n$. Since $\sigma>n$, we have $\hat{u}-a \in H^{1, \sigma}(\Omega)^{n}$ and hence by the Sobolev embedding theorem $\hat{u}-a \in C^{0}(\bar{\Omega})^{n}$ and $\lim _{x \rightarrow \infty}|\hat{u}(x)-a|=0$, from which (1.3) follows. Now the uniqueness result of Lemma 4.3(ii) yields that $\{u, p\}=L\{A, a\}$.

We next consider the case $n=2$ and $q>2$. Let $\{u, p\} \in \mathbb{N}_{q}$. Then $\nabla u-A \in L^{q}(\Omega)^{2^{2}}$ for some $A \in \mathbb{R}^{2^{2}}$ with $\operatorname{Tr} A=0$. Since $u-A x-\Gamma(A \cdot) \in$ $\hat{H}_{0}^{1, q}(\Omega)^{2}$, we obtain by Lemma 2.2(ii) that $u(x)-A x=\mathrm{O}\left(|x|^{1-2 / q}\right)$ as $|x| \rightarrow \infty$.

Applying the regularity theorem of Finn-Smith [13, Theorem 5.11-12], we see $u-A x \in C^{2}(\Omega) \cap C^{1}(\bar{\Omega})$ and $p \in C^{1}(\Omega) \cap C^{0}(\bar{\Omega})$. Then it follows from the representation formula of Chang-Finn [11, Theorem 1] that

$$
\begin{aligned}
& u(x)-A x-E(x) a=u^{\infty}-\int_{\partial \Omega} A \xi \cdot T E(x-\xi) \nu_{\xi} \mathrm{d} S_{\xi} \\
& -\int_{\partial \Omega}(E(x-\xi)-E(x)) T(u(\xi) A \xi) \nu_{\xi} \mathrm{d} S_{\xi}
\end{aligned}
$$

where $a=-\int_{\partial \Omega} T(u(\xi)-A \xi) \nu_{\xi} \mathrm{d} S_{\xi}, u^{\infty} \in \mathbb{R}^{2}$ and $T$ denotes the stress tensor. Using the explicit expression of $E$, we see that

$$
\sup _{\xi \in \partial \Omega}\left|\nabla_{x} T E(x-\xi)\right|=\mathrm{O}\left(|x|^{-2}\right), \sup _{\xi \in \partial \Omega}\left|\nabla_{x} E(x-\xi)-\nabla_{x} E(x)\right|=\mathrm{O}\left(|x|^{-2}\right)
$$

as $x \rightarrow \infty$. Therefore $\nabla(u-A x-E a) \in L^{2}(\Omega)^{2^{2}}$ and it follows from the uniqueness proved in Lemma 4.2(ii) that $\{u, p\}=L\{A, a\}$.

Properties of $\operatorname{dim} \mathbb{N}_{q}$ and $\operatorname{dim} \mathbb{N}_{q}^{0}$.
Let us first consider the case $n^{\prime}<q<\infty$ ( $n \geq 2$ ). Then by Lemma 4.4(ii) and the definition of $V$, we obtain

$$
\begin{equation*}
\operatorname{dim} \mathbb{N}_{q}=\operatorname{dim} V=n^{2}+n-1, \operatorname{dim} \mathbb{N}_{q}^{0}=n, n^{\prime}<q<\infty(n \geq 2) \tag{4.3}
\end{equation*}
$$

Hence Theorem B(ii) follows from (4.3) and Lemmas 4.2-3(ii).
We next consider the case $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. By Theorem A and the definition of $\mathbb{N}_{q}^{0}$, we have

$$
\begin{equation*}
\operatorname{dim} \mathbb{N}_{q}^{0}=0,1<q \leq n^{\prime}(n \geq 3), 1<q<2(n=2) \tag{4.4}
\end{equation*}
$$

Moreover, it follows from Lemma 4.4(i) that $\mathbb{N}_{q}^{0}$ is isometric to the subspace $W_{q}$ of $\hat{V}_{q}$ :

$$
W_{q} \equiv\left\{a \in \mathbb{R}^{n} ; \int_{\partial \Omega}\left\{a \cdot \frac{\partial v}{\partial \nu}-\chi a \cdot \nu\right\} \mathrm{d} S=0 \text { for all }\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}\right\}
$$

Hence $W_{q}=\{0\}$. On the other hand, by (4.3), we see $\operatorname{dim} \mathbb{N}_{q^{\prime}}^{0}=n$. Therefore it follows that $\operatorname{dim} \hat{V}_{q}=n^{2}+n-1-\operatorname{dim} \mathbb{N}_{q^{\prime}}^{0}=n^{2}-1$. Now, Lemma 4.4(i) yields

$$
\begin{equation*}
\operatorname{dim} \mathbb{N}_{q}=\operatorname{dim} \hat{V}_{q}=n^{2}-1,1<q \leq n^{\prime}(n \geq 3), 1<q<2(n=2) . \tag{4.5}
\end{equation*}
$$

Hence Theorem $\mathrm{B}(\mathrm{i})$ follows from (4.4-5) and Lemmas 4.2-3(i).
(iii) Case $n=q=2$. In the same way as in Lemmas 4.2-4.3, we can construct a bijective operator $L^{\prime}: A \rightarrow\{u, p\}$ from $V^{\prime} \equiv\left\{A \in \mathbb{R}^{2^{2}} ; \operatorname{Tr} A=0\right\}$ onto $\mathbb{N}_{2}$ such that $u$ satisfies (4.2) with $a=0$. Hence we get $\operatorname{dim} \mathbb{N}_{2}=3$. By Lemma 2.2(ii) and Corollary 3.4(ii), we have $\mathbb{N}_{2}^{0}=\operatorname{Ker} S_{2}=\{0,0\}$. And therefore existence and uniqueness derive from the same argument as before, so we may omit the details.

### 4.3. Inhomogeneous equations; Proof of Theorem C.

Recall the function $A \cdot+a: x \in \partial \Omega \rightarrow A x+a \in \mathbb{R}^{n}$ and set $w=\Gamma(A \cdot+a)$. Taking $\hat{u}=u-A x-a+w$, we get from ( $S$ )

$$
\begin{align*}
& -\Delta \hat{u}+\nabla p=f-\Delta w, \operatorname{div} \hat{u}=g \operatorname{Tr} A+\operatorname{div} w \text { in } \Omega, \\
& \hat{u}=0 \text { on } \partial \Omega \tag{4.6}
\end{align*}
$$

In order to solve ( $S$ ), we shall make use of (4.6).
Proof of Theorem C. (i) Case $1<q \leq n^{\prime}$ for $n \geq 3$ and $1<q<2$ for $n=2$. As we have seen in the proof of Theorem B , (1.5) is equivalent to the identity

$$
(f-\Delta w, v)+(g-\operatorname{Tr} A+\operatorname{div} w,-\chi)=0 \text { for all }\{v, \chi\} \in \mathbb{N}_{q^{\prime}}^{0}
$$

This implies that $\{f-\Delta w, g-\operatorname{Tr} A+\operatorname{div} w\} \in\left(\operatorname{Ker} T_{q^{\prime}}\right)^{\perp}$. Hence by Corollary 3.4(i), there is a unique generalized solution $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ of (4.6). By Lemma 2.2(i) we have also $\hat{u} \in L^{n q /(n-q)}(\Omega)^{n}$. Moreover, from Theorem 3.3 and the continuity of the extension operator $\Gamma$, we obtain

$$
\begin{aligned}
\|\nabla \hat{u}\|_{q}+\|p\|_{q} & \leq C\left(\|f-\Delta w\|_{-1, q}+\|g-\operatorname{Tr} A+\operatorname{div} w\|_{q}\right) \\
& \leq C\left(\|f\|_{-1, q}+\|g-\operatorname{Tr} A\|_{q}+|A|+|a|\right),
\end{aligned}
$$

where $C=C(\Omega, n, q)$. Taking $u=\hat{u}+A x+a-w$, we see that $\{u, p\}$ is the desired generalized solution of $(S)$. The uniqueness follows from the fact that $\mathbb{N}_{q}^{0}=\{0,0\}$. Conversely, suppose that $\{u, p\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ is a generalized solution of ( $S$ ) satisfying (1.2) and (1.4) for some $A \in \mathbb{R}^{n^{2}}$ with $\operatorname{Tr} A-g \in L^{q}(\Omega)$ and $a \in \mathbb{R}^{n}$. Taking $w=\Gamma(A \cdot+a)$, we see by Lemma 2.2(i) that $\hat{u}=u-A x-a+w \in \hat{H}_{0}^{1, q}(\Omega)^{n}$ and that $\{\hat{u}, p\}$ is a generalized solution of (4.6). Hence it follows from Corollary 3.4(i) that

$$
\{f-\Delta w, g-\operatorname{Tr} A+\operatorname{div} w\} \in R\left(S_{q}\right)=\left(\operatorname{Ker} T_{q^{\prime}}\right)^{\perp}
$$

from which we get (1.5).
(ii) Case $n^{\prime}<q<n$ for $n \geq 3$. Take $w=\Gamma(A \cdot+a)$ and consider (4.6) for $\{\hat{u}, p\}$. Then we have by Corollary $3.4(\mathrm{ii})$ that (4.6) is uniquely solvable in $\hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ for all $f, g, A$ and $a$ as given in the assumptions. Then the
proof of existence and uniqueness is quite the same as in the case (i) above. Suppose in addition that $f \in \hat{H}^{-1, r}(\Omega)^{n}$ and $g-\operatorname{Tr} A \in L^{r}(\Omega)$ for some $r>n$. Since $1<r^{\prime}<q^{\prime}<n$, it follows from (2.2) and an interpolation argument (see, e.g., Triebel [32, 1.11.2]) that $f \in \hat{H}^{-1, \gamma}(\Omega)^{n}$ and $g-\operatorname{Tr} A \in L^{\gamma}(\Omega)$ for all $q \leq \gamma \leq r$. Hence we have by Theorem 3.1 that $\{\hat{u}, p\} \in \hat{H}_{0}^{1, \gamma}(\Omega)^{n} \times L^{\gamma}(\Omega)$ for all $q \leq \gamma \leq r$ and that, in particular, $\hat{u} \in H^{1, s}(\Omega)^{n}$ for some $s>n$. By the Sobolev embedding theorem, we obtain $\hat{u} \in C^{0}(\bar{\Omega})^{n}$ and $\lim _{x \rightarrow \infty}|\hat{u}(x)|=0$. Now it is easy to see that $u=\hat{u}+A x+a-w$ satisfies (1.3).
(iii) Case $n \leq q<\infty$ for $n \geq 3$ and $2<q<\infty$ for $n=2$. Taking $w=\Gamma(A \cdot)$ in (4.6), we have by Corollary 3.4(iii) that there is at least one generalized solution $\{\hat{u}, p\} \in \hat{H}_{0}^{1, q}(\Omega)^{n} \times L^{q}(\Omega)$ of (4.6). On the other hand, we have by Lemma 2.2 (ii) that $\mathbb{N}_{q}^{0}=\operatorname{Ker} S_{q}$ and that $R\left(S_{q}\right)$ is isometric to the quotient space $X_{q} / \mathbb{N}_{q}^{0}$. Therefore it follows that

$$
\begin{aligned}
& \inf \left\{\|\nabla \hat{u}-\nabla v\|_{q}+\|p-\chi\|_{q} ;\{v, \chi\} \in \mathbb{N}_{q}^{0}\right\} \\
& \leq C\left(\|f-\Delta w\|_{-1, q}+\|g-\operatorname{Tr} A+\operatorname{div} w\|_{q}\right) \\
& \leq C\left(\|f\|_{-1, q}+\|g-\operatorname{Tr} A\|_{q}+|A|\right)
\end{aligned}
$$

where $C=C(\Omega, n, q)$. Taking $u=\hat{u}+A x-w$, we see that $\{u, p\} \in$ $H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ has the desired property.

Suppose that $\{\tilde{u}, \tilde{p}\} \in H_{\mathrm{loc}}^{1, q}(\bar{\Omega})^{n} \times L^{q}(\Omega)$ is another generalized solution of (S) with (1.2). Set $u^{\prime}=u-\tilde{u}$ and $p^{\prime}=p-\tilde{p}$. Then we get $\left\{u^{\prime}, p^{\prime}\right\} \in \mathbb{N}_{q}^{0}$. Hence

$$
\begin{aligned}
& \inf \left\{\|\nabla \tilde{u}-A-\nabla v\|_{q}+\|\tilde{p}-\chi\|_{q} ; \quad\{v, \chi\} \in \mathbb{N}_{q}^{0}\right\} \\
& =\inf \left\{\|\nabla u-A-\nabla v\|_{q}+\|p-\chi\|_{q} ; \quad\{v, \chi\} \in \mathbb{N}_{q}^{0}\right\}
\end{aligned}
$$

so uniqueness and (1.7) follow.
(iv) Case $n=q=2$. By Theorem $\mathrm{B}(\mathrm{iii})$, we see that $\mathbb{N}_{2}^{0}=\{0,0\}$; so the proof is quite similar to that of the case (ii) above.

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