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## **Hamiltonian Dynamics on Pseudodifferential Symbols**

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# Hamiltonian dynamics on pseudodifferential symbols

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The talk is based on the joint work with Ilya ZAKHAREVICH ( Math. Dept., MIT ) on the Lie-Poisson group of pseudodifferential operators. We present here a short summary of the results and refer to [1,2] for all the details.

The ring  $\mathfrak{G}$  of pseudodifferential symbols on the circle by definition consists of formal series  $A(x, D) = \sum_{-\infty}^n a_j(x)D^j$  with respect to  $D (= d/dx)$  where  $a_j \in C^\infty(S^1, \mathbb{R} \text{ or } \mathbb{C})$ . The multiplication law in  $\mathfrak{G}$  is given by the Leibnitz rule for multiplication of symbols:  $A(x, \xi) \circ B(x, \xi) = \sum_{n \geq 0} \frac{1}{n!} A_\xi^{(n)}(x, \xi) B_x^{(n)}(x, \xi)$  where  $A_\xi^{(n)} = d^n/d\xi^n A(x, \xi)$ ,  $B_x^{(n)} = d^n/dx^n B(x, \xi)$ . The Lie algebra structure on  $\mathfrak{G}$  is natural:  $[A, B] = A \circ B - B \circ A$ , and the operator  $\text{res} : \mathfrak{G} \rightarrow C^\infty(S^1)$  is defined by  $\text{res}(\sum a_i(x)D^i) = a_{-1}(x)$ .

The formal expression  $\log D$  defines an outer derivation of  $\mathfrak{G}$  by:  $[\log D, A] = \log D \circ A - A \circ \log D$  where the r.h.s is understood in the sense of the Leibnitz law above and belongs to  $\mathfrak{G}$  for any  $A \in \mathfrak{G}$ . Then the ( $\mathbb{R}$  or  $\mathbb{C}$ -valued) 2-cocycle

$$c(A, B) = \int \text{res}([\log D, A] \circ B) = \int \text{res} \left( \sum_{k \geq 1} \frac{(-1)^{k+1}}{k} A_x^{(k)} D^{-k} \circ B \right)$$

gives a nontrivial central extension of the Lie algebra  $\mathfrak{G}$  (here  $A$  and  $B$  are arbitrary pseudodifferential symbols on  $S^1$ ) (see [3]). The restrictions of this cocycle to the subalgebra  $\mathfrak{G}_{DO}$  of purely differential operators  $\{\sum_0^n a_j(x)D^j\}$  gives the Kac-Peterson cocycle [4], while restricting it to the Lie algebra of vector fields  $\{f(x)D\}$  one obtains the Gelfand-Fuchs cocycle of the Virasoro algebra.

Let the algebra  $\tilde{\mathfrak{G}} = \left\{ \left( \sum_{j=-\infty}^n a_j(x)D^j + \lambda \log D, c \right) \right\}$  be the extension of  $\mathfrak{G}$  by the

2-cocycle  $c(A, B)$  and by the cocentral element  $\log D$ . The algebra  $\tilde{\mathfrak{G}}$  has a natural *ad*-invariant nondegenerate inner product (“Killing form”) and two remarkable isotropic (with

respect to this inner product) subalgebras: 1)  $\tilde{\mathfrak{G}}_{DO}$  which is the algebra of centrally extended differential operators  $\{(\sum_{j \geq 0} a_j(x)D^j, c)\}$ , and 2)  $\tilde{\mathfrak{G}}_{\text{Int}}$  which is the algebra of integral

symbols together with  $\log D : \{(\sum_{j=-\infty}^{-1} a_j(x)D^j + \lambda \log D)\}$ . The triple  $(\tilde{\mathfrak{G}}, \tilde{\mathfrak{G}}_{DO}, \tilde{\mathfrak{G}}_{\text{Int}})$  is a Manin triple (or, equivalently,  $\tilde{\mathfrak{G}}_{\text{Int}}$  is a Lie bialgebra).

This implies that the Lie group  $\tilde{G}_{\text{Int}} = \{(1 + \sum_{k=-\infty}^{-1} u_k(x)D^k) \circ D^\alpha | \alpha \in \mathbb{R} \text{ ou } \mathbb{C}\}$

corresponding to the Lie bialgebra  $\tilde{\mathfrak{G}}_{\text{Int}}$  has a natural Poisson-Lie structure. This structure generalises the second Adler-Gelfand-Dickey structure [5,6] on fixed order differential operators.

This Lie group is quasi-nilpotent, and its the exponential map  $\exp : \tilde{\mathfrak{G}}_{\text{Int}} \rightarrow \tilde{G}_{\text{Int}}$  is one-to-one. It allows one to define arbitrary complex powers of an operator  $L \in \tilde{G}_{\text{Int}}$ . Then the following infinite sequence of evolution equations on the coefficients of  $L$ :  $\frac{\partial L}{\partial t_m} = [L, (L^{m/\alpha})_+]$ ,  $m = 1, 2, \dots$  is well defined ( $\deg L = \alpha$ ,  $\deg L^{m/\alpha} = m$ , and the operation  $+$ , taking the differential part, makes sense). For any  $\alpha \neq 0$  and any integral  $m$  these equations are Hamiltonian relative to the Poisson structure on  $\tilde{G}_{\text{Int}}$  with Hamiltonian functions  $H_m(L) = \frac{\alpha}{m} \int \text{res}(L^{m/\alpha})$ . The set  $(\{H_m\}, m = 1, 2, \dots)$  is the set of integrals in involution, and this family interpolates between the KP hierarchy ( $\alpha = 1$ ) and n-KdV hierarchies ( $\alpha = n$ ).

The same system of evolution equations on the subspace  $\{L = D + \psi(x)D^{-1}\psi^*(x)\}$  (where  $\psi(x)$  is a complex-valued function on the circle, and Hamiltonians  $\tilde{H}_m = i^m H_m$ ) generates NLS hierarchy. The classical NLS-equation corresponds to the Hamiltonian function  $\tilde{H}_2$ , cf.[7,8].

This Poisson-Lie point of view suggests a geometric interpretation of relation between  $W_n$  and  $W_\infty$  - algebras appeared recently in theoretical physics.

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